



Trans-Planckian Physics in Quantum Cosmology

Christian Durán Romero, Mercedes Martín
Benito, Luis J. Garay and Rita B. Neves

 $\langle G | \hbar \rangle$

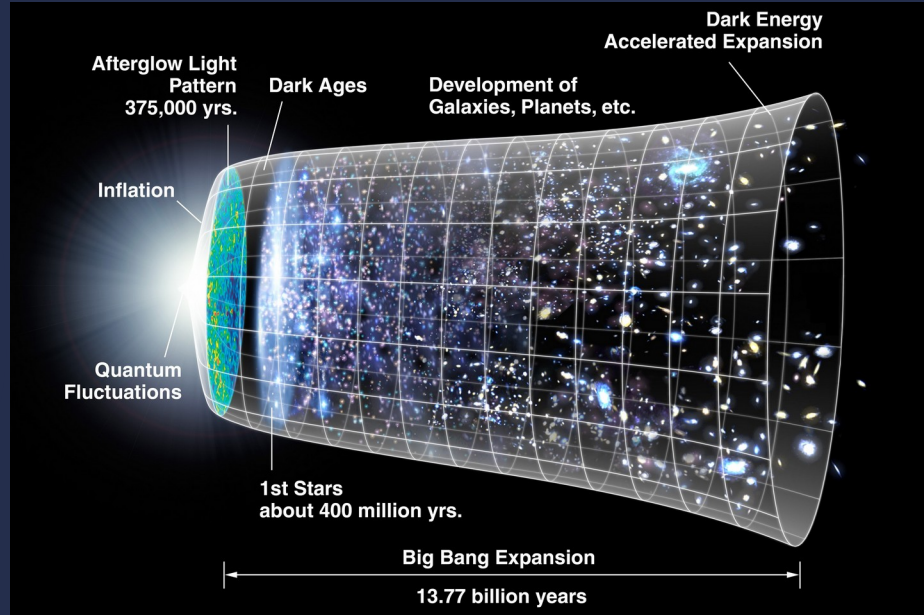
Quantum Fields & Gravity Group

The background features a dark blue gradient with several glowing, interconnected nodes and faint orbital paths, resembling a network or a celestial map. A prominent bright purple and white cluster of points is located in the upper center, while a large, bright cyan and white circular structure with internal connections is visible in the lower right. The overall aesthetic is futuristic and scientific.

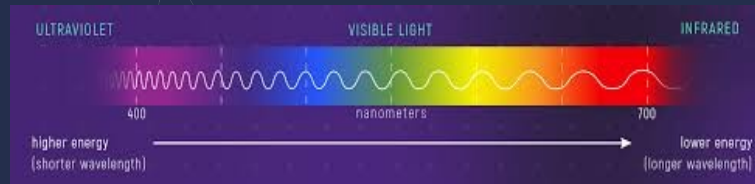
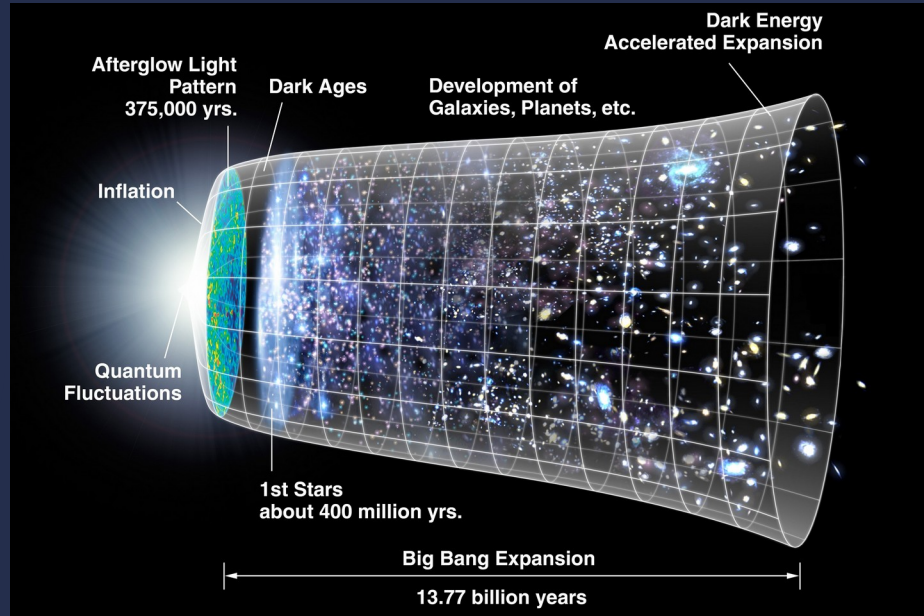
01

The Trans-Planckian problem in Cosmology

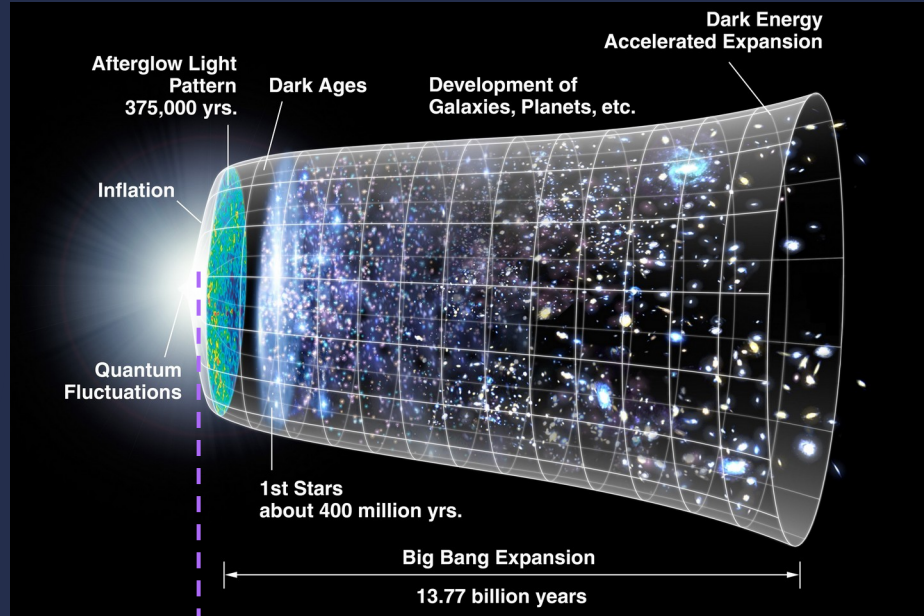
The Universe is expanding!



The Universe is expanding!

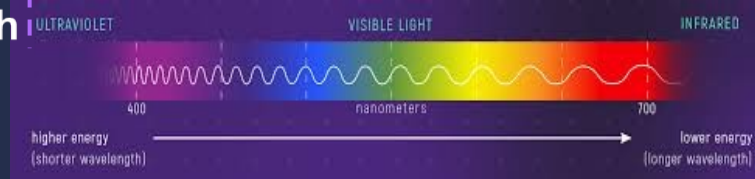


The Universe is expanding!



Planck's length

L_p



At Planck's length we are entering in the framework of Quantum Gravity. Then... what about cosmological predictions?



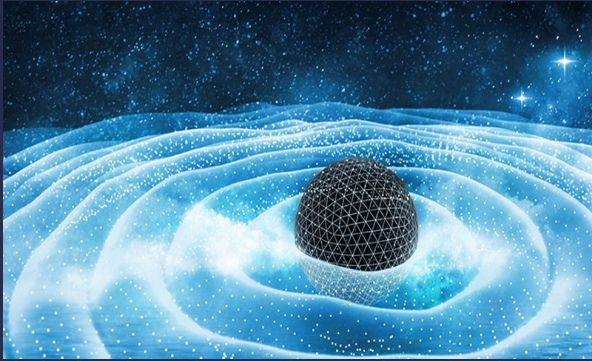
The background features a teal sphere on the left with a bright white light source. A purple starburst is centered behind the text. The text '02 Semiclassical Gravity' is displayed in white.

02 Semiclassical Gravity

Quantum fields in curved spacetimes

Semiclassical Einstein's equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G (T_{\mu\nu}^{cla} + \langle T_{\mu\nu} \rangle)$$



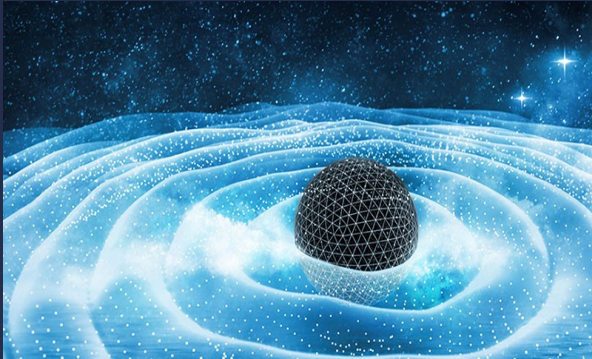
Quantum fields in curved spacetimes

Semiclassical Einstein's equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G (T_{\mu\nu}^{cla} + \langle T_{\mu\nu} \rangle)$$



Einstein's tensor. Responsible for the curvature of the spacetime. Purely classical.



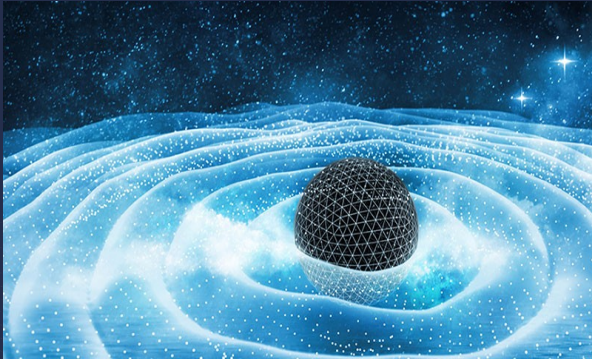
Quantum fields in curved spacetimes

Semiclassical Einstein's equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G (T_{\mu\nu}^{cla} + \langle T_{\mu\nu} \rangle)$$



Vacuum expectation value of the stress-energy tensor. It contains information about the QFT. Purely quantum.



Quantum fields in curved spacetimes

Semiclassical Einstein's equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G (T_{\mu\nu}^{cla} + \langle T_{\mu\nu} \rangle)$$

1. Energy of the particles $E \ll \frac{1}{L_P}$ (**no quantum gravity**)

2. Negligible *backreactions* $|\langle T_{\mu\nu} \rangle| \ll \frac{1}{L_P^2}$ (**no changes in the geometry** from QF)

3. Low curvature $R \ll \frac{1}{L_P^2}$ (the background $g_{\mu\nu}$ is fixed, i.e., **no quantum fluctuations**)



Quantum fields in curved spacetimes

1. Energy of the particles $E \ll \frac{1}{L_P}$ (**no quantum gravity**)
2. Negligible *backreactions* $|\langle T_{\mu\nu} \rangle| \ll \frac{1}{L_P^2}$ (**no changes in the geometry** from QF)
3. Low curvature $R \ll \frac{1}{L_P^2}$ (the background $g_{\mu\nu}$ is fixed, i.e., **no quantum fluctuations**)

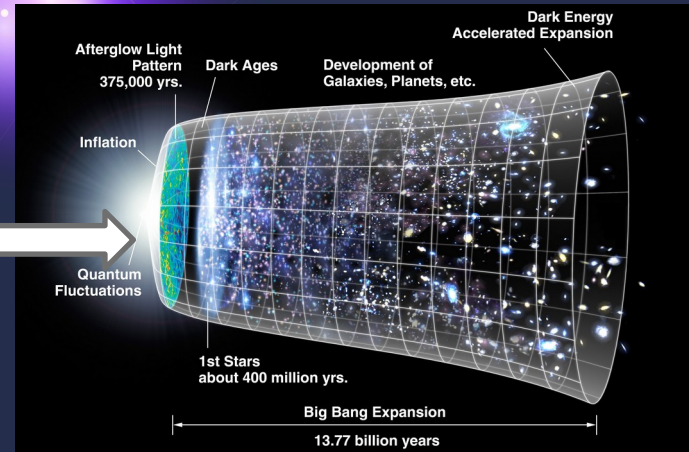
Does this always hold?

Quantum fields in curved spacetimes

1. Energy of the particles $E \ll \frac{1}{L_P}$ (**no quantum gravity**)
2. Negligible *backreactions* $|\langle T_{\mu\nu} \rangle| \ll \frac{1}{L_P^2}$ (**no changes in the geometry from QF**)
3. Low curvature $R \ll \frac{1}{L_P^2}$ (the background $g_{\mu\nu}$ is fixed, i.e., **no quantum fluctuations**)

Does this always hold? **NO**

It breaks down here

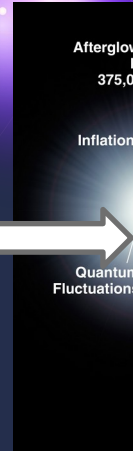


Quantum fields in curved spacetimes

1. Energy of the particles $E \ll \frac{1}{L_P}$ (**no quantum gravity**)
2. Negligible *backreactions* $|\langle T_{\mu\nu} \rangle| \ll \frac{1}{L_P^2}$ (**no changes in the geometry** from QF)
3. Low curvature $R \ll \frac{1}{L_P^2}$ (the background $g_{\mu\nu}$ is fixed, i.e., **no quantum fluctuations**)

Does this always hold? **NO**

It breaks down here



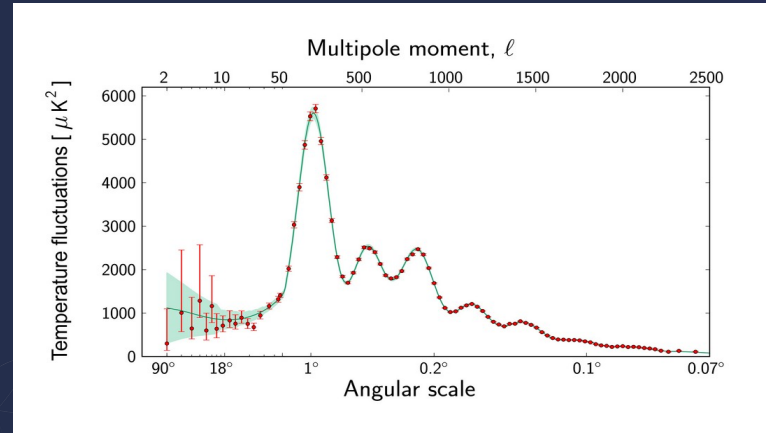
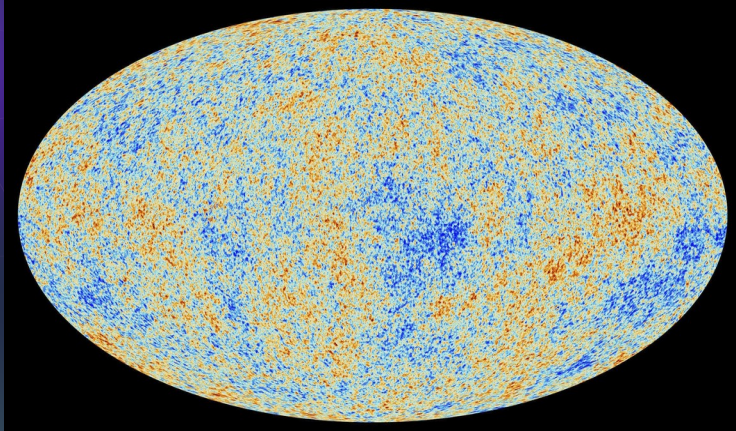
TRANS-
PLANCKIAN
PROBLEM IN
COSMOLOGY



03 Cosmology

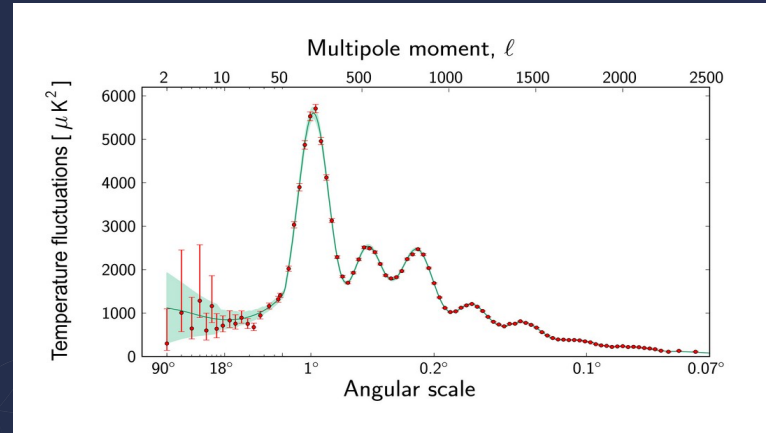
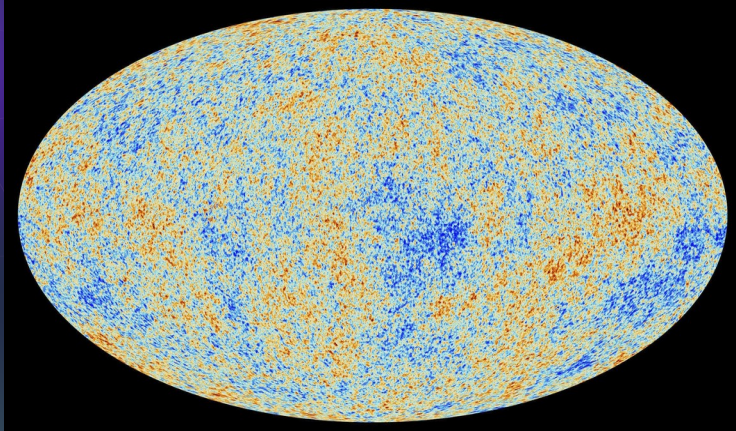
A quick overview of the problem

The standard study of the Cosmos using QFTCS give rise to one of the most impressive predictions in Cosmology, the Cosmological Microwave Background (CMB) and its power spectrum



A quick overview of the problem

The standard study of the Cosmos using QFTCS give rise to one of the most impressive predictions in Cosmology, the Cosmological Microwave Background (CMB) and its power spectrum



And what about the trans-Planckian problem?

04

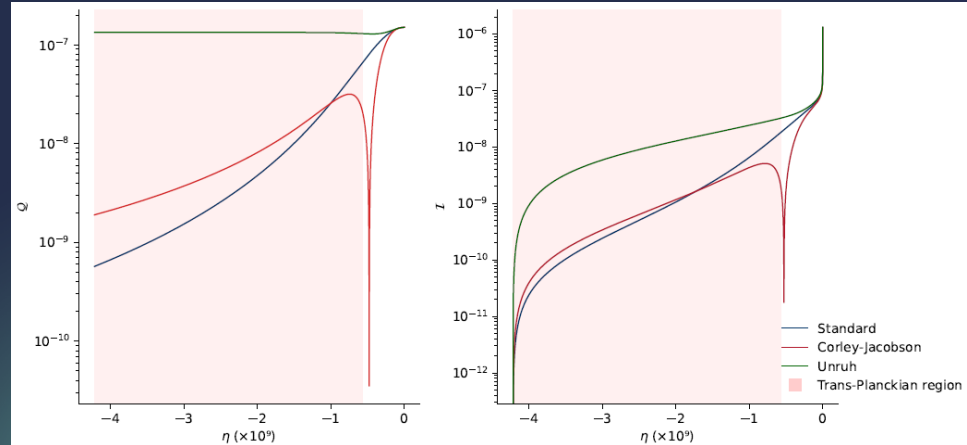
In this thesis

1. Energy of the particles $E \ll \frac{1}{L_P}$ (no quantum gravity)

i) Near L_P quantum gravity effects are expected

ii) The expression for “the energy of the particles”, what is called standard *dispersion relation* should be changed. In the literature this is done by introducing *ad hoc* the so called *modified dispersion relations (MDRs)*.

iii) Is the CMB compatible with this change?



2. Negligible *backreactions* $|\langle T_{\mu\nu} \rangle| \ll \frac{1}{L_P^2}$ (no changes in the geometry from QF)

- i) Under what circumstances the backreactions are negligible?
- ii) Physical quantities (like $\langle T_{\mu\nu} \rangle$ and $\langle \phi^2 \rangle$) are divergent, therefore, must be regularized and renormalized. We have developed a new regularization method and we are collaborating with other members of IPARCOS
- iii) Semiclassical gravity is not a perturbatively renormalizable theory with the standard dispersion relation. What happens in the presence of MDRs?

3. Low curvature $R \ll \frac{1}{L_P^2}$ (the background $g_{\mu\nu}$ is fixed, i.e., no quantum fluctuations)

Since quantum gravity effects are expected to appear at scales near L_P , our porpouse for future work is to include them. How? We plan to develop a MDR coming from a quantum theory approach, namely, from *Loop Quantum Gravity*.

In this approach, spacetime itself becomes quantum and, therefore, we expect that a natural dispersion relation arises from the theory. In this scenario, we could test the robustness of the Power Spectrum and the renormalizability of the resulting semiclassical limit



Conclusions

01

The trans-Planckian
problem in Cosmology
appears as a
consequence of QFTCS



Conclusions

01

The trans-Planckian problem in Cosmology appears as a consequence of QFTCS

02

Since new phenomena are expected, we introduce new physics with MDRs



Conclusions

01

The trans-Planckian problem in Cosmology appears as a consequence of QFTCS

02

Since new phenomena is expected, we introduce new physics with MDRs

03

Cosmological observations must be respected in the new theories



Conclusions

01

The trans-Planckian problem in Cosmology appears as a consequence of QFTCS

04

When quantum meets gravity, new divergences arise and we must be very careful to work in the correct regime

02

Since new phenomena is expected, we introduce new physics with MDRs

03

Cosmological observations must be respected in the new theories



Conclusions

01

The trans-Planckian problem in Cosmology appears as a consequence of QFTCS

04

When quantum meets gravity, new divergences arise and we must be very careful to work in the correct regime

02

Since new phenomena are expected, we introduce new physics with MDRs

05

As future work, we want to derive a MDR based on LQG

03

Cosmological observations must be respected in the new theories



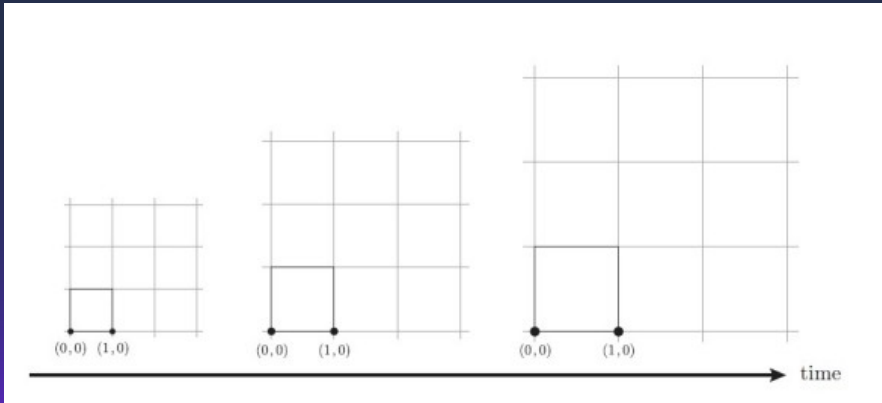
“That’s all Folks!”

Back up slides

Cosmological spacetime FLRW

$$ds^2 = -dt^2 + a^2(t)d\vec{x}^2$$

Cosmological redshift



$$L_{\text{fis}}(t) = a(t)L_{\text{com}}$$

$$k_{\text{fis}}(t) = \frac{k_{\text{com}}}{a(t)}$$

Canonical quantization: Real scalar field

$$(\square - m^2)\phi = 0$$

After the expansion in modes, the temporal satisfies

$$\ddot{F}_k + \omega_k^2 F_k = 0$$

with

$$\omega_k^2 \equiv \left(\frac{k}{a}\right)^2 + m^2 + \sigma, \quad \sigma \equiv -\frac{3}{2} \frac{\ddot{a}}{a} - \frac{3}{4} \frac{\dot{a}^2}{a^2}$$

Back up slides

MDRs:

$$\omega_k^2 = \mathcal{K}^2 + m^2 + \sigma, \quad \mathcal{K} \xrightarrow[k \rightarrow 0]{} k/a$$

Standard

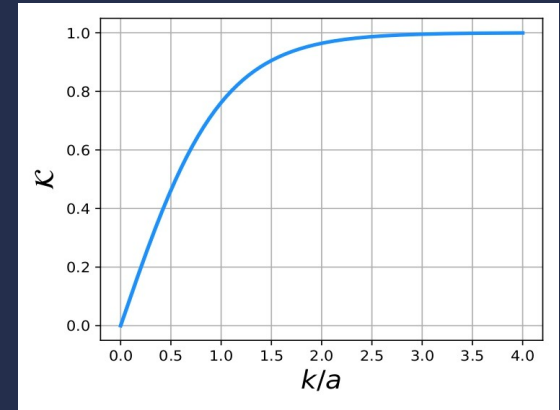
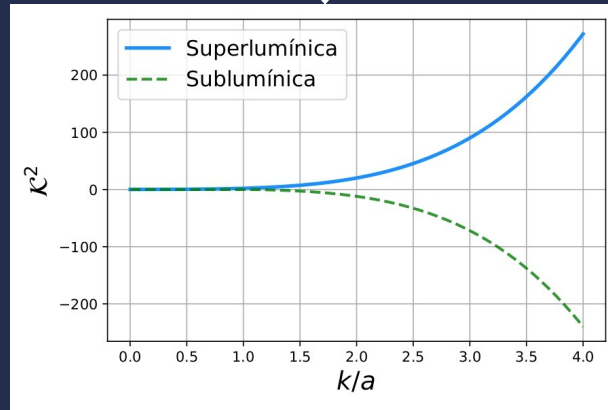
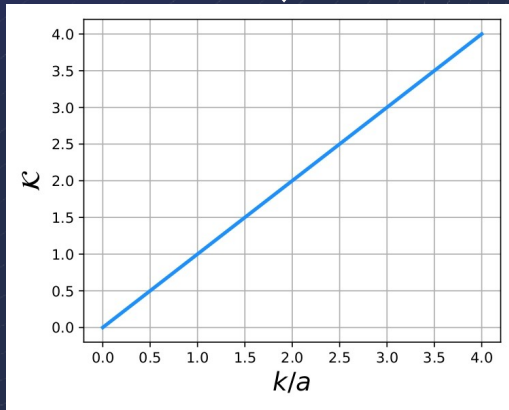
$$\mathcal{K} = k/a$$

Corley-Jacobson

$$\mathcal{K}^2 = \frac{k^2}{a^2} \left(1 \pm \frac{k^2}{a^2 \kappa_c^2} \right)$$

Unruh

$$\mathcal{K} = \kappa_c \tanh \left(\frac{k}{a \kappa_c} \right)$$



Back up slides

Density of states

$$\rho_k \equiv \left| \frac{dk}{d\omega_k} \right|$$

Standard

$$\rho_k \sim \text{const.}$$

Superluminal

$$\rho_k \sim \frac{a\kappa_c}{k}$$

Unruh

$$\rho_k \sim \cosh^2 \frac{k}{a\kappa_c}$$

