

# QCD evolution 2026

# Simultaneous analysis of TMDs and PDFs

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# BASED ON

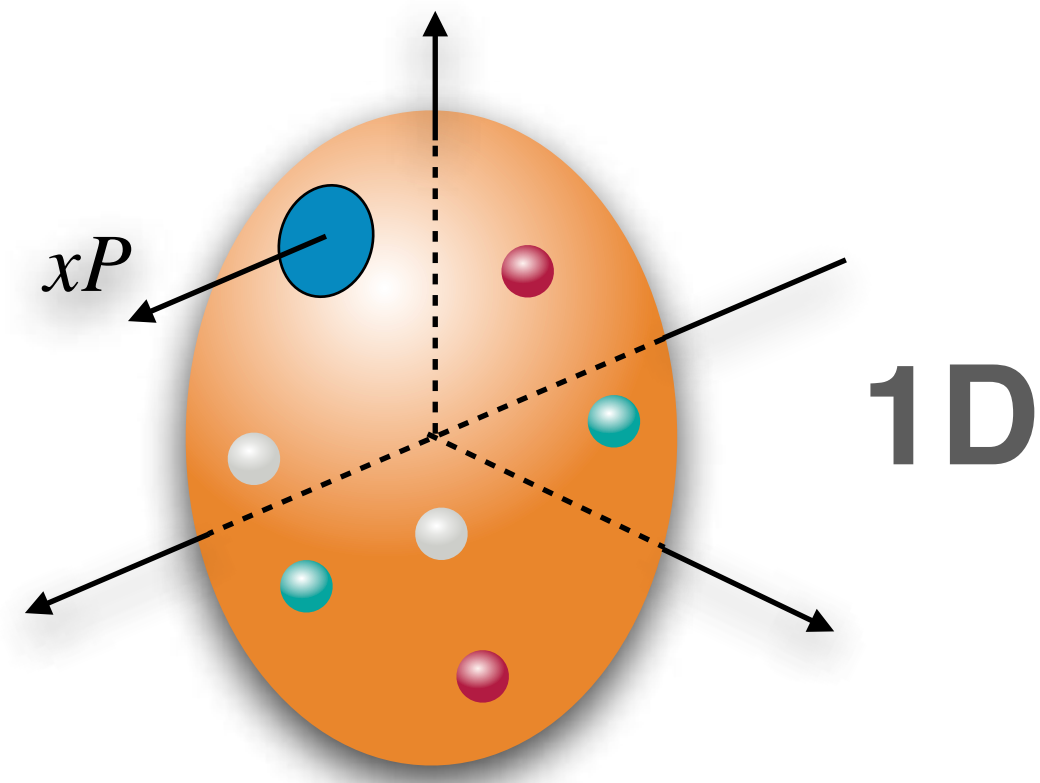


## TMD WORKING GROUP

P. C. Barry, A. Prokudin, T. Anderson, C. Cocuzza, L. Gamberg, W. Melnitchouk, E. Moffat, D. Pitonyak, J.-W. Qiu, N. Sato, A. Vladimirov, and R. Whitehill *2510.13771*

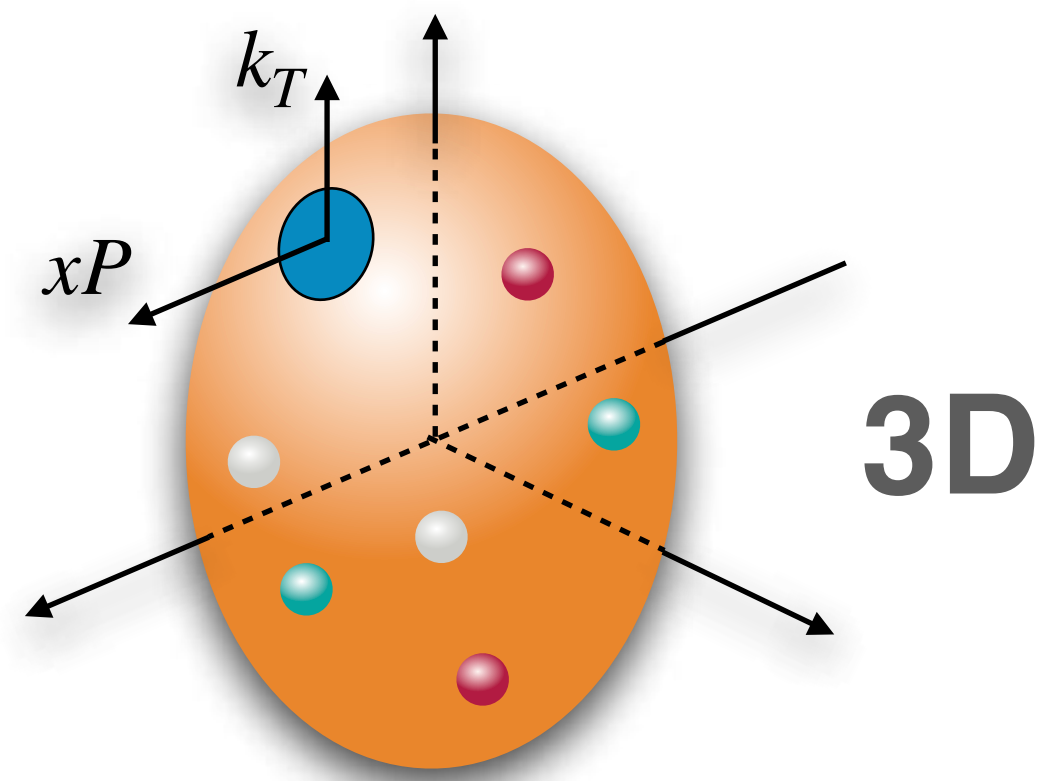


## Collinear Parton Distribution Functions (PDFs)



One large scale ( $Q$ ) sensitive to particle nature of quark and gluons.  
PDFs encode the 1D structure of the nucleon, the distribution in  $x$ .

## Transverse Momentum Dependent Distributions (TMDs)



One large scale ( $Q$ ) sensitive to particle nature of quark and gluons.  
One small scale ( $k_T$ ) sensitive to *how QCD bounds partons* and to the detailed structure at  $\sim$ fm distances.  
TMDs encode the 3D structure of the nucleon, the distribution in  $x$  and  $k_T$  (or Fourier conjugate variable  $b_T$ ).

# TRANSVERSE MOMENTUM DEPENDENT FACTORIZATION

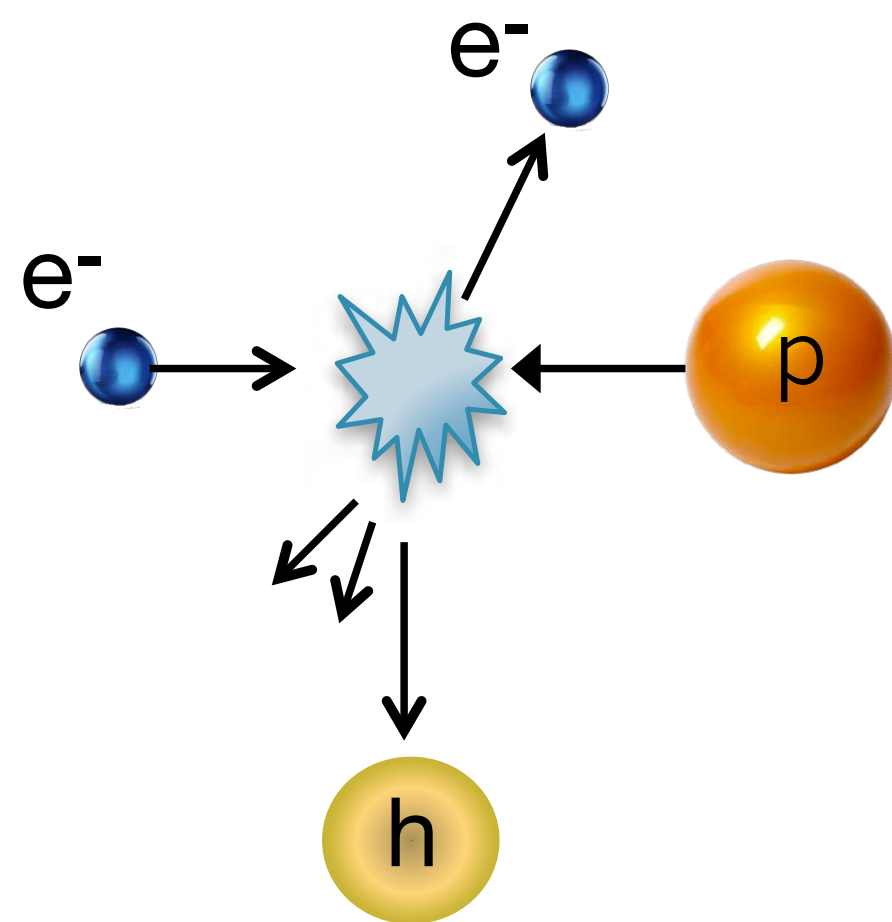
TMD handbook: <https://inspirehep.net/literature/2650019>

Small scale  $\longrightarrow q_T \ll Q \longleftarrow$  Large scale

The confined motion ( $k_T$  dependence) is encoded in TMDs

## Semi-Inclusive DIS

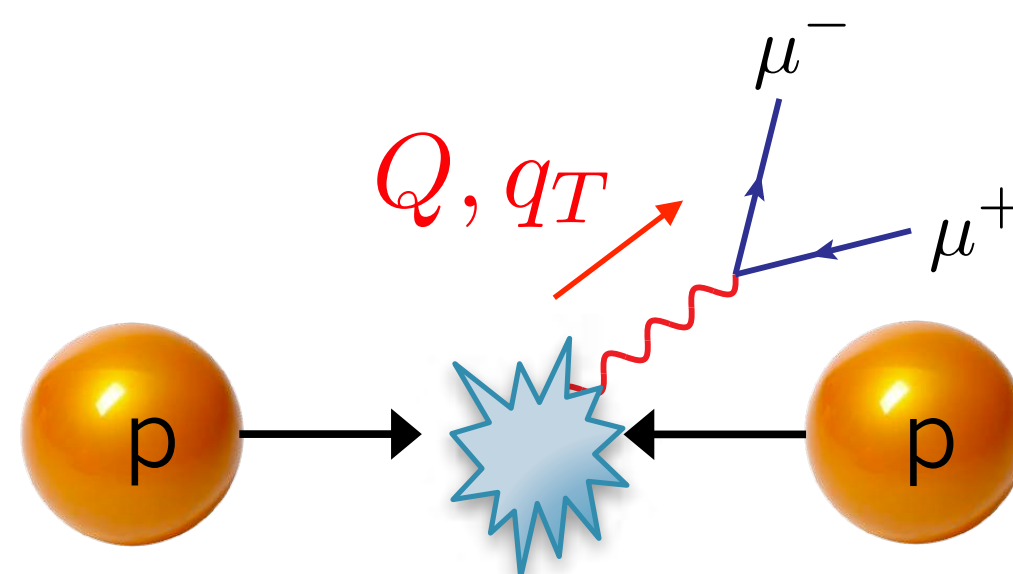
$$\sigma \sim f_{q/P}(x, k_T) D_{h/q}(z, k_T)$$



Meng, Olness, Soper (1992)  
Ji, Ma, Yuan (2005)  
Idilbi, Ji, Ma, Yuan (2004)  
Collins (2011)

## Drell-Yan

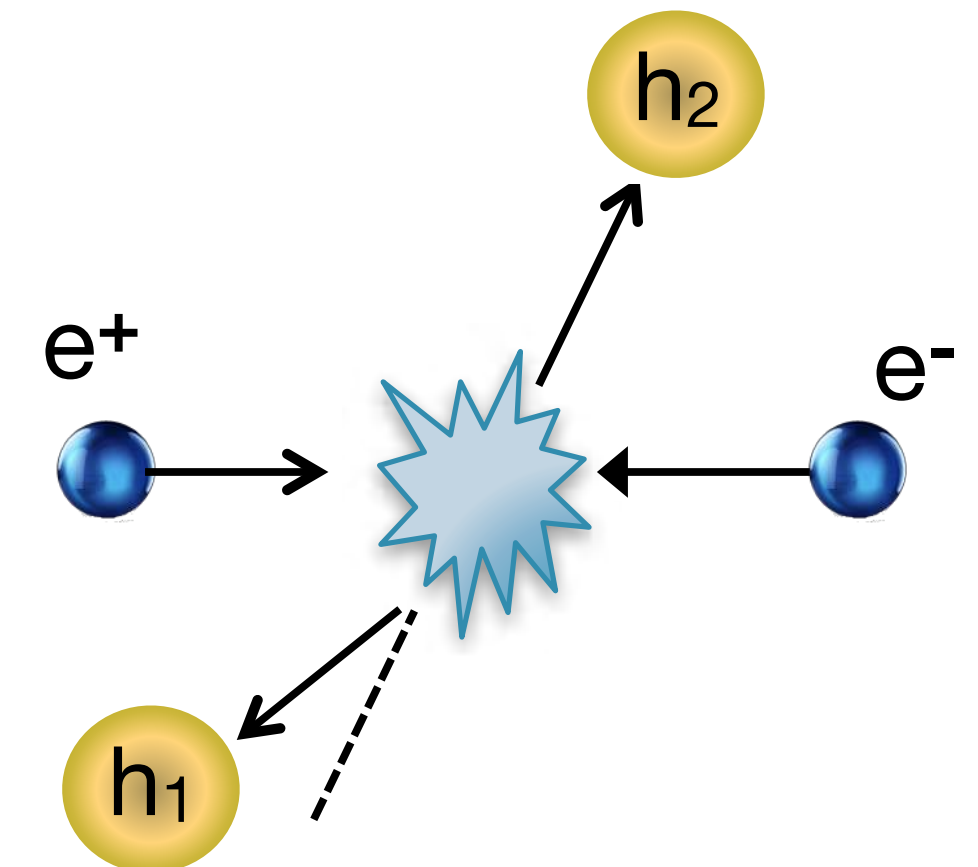
$$\sigma \sim f_{q/P}(x, k_T) f_{q/P}(x, k_T)$$



Collins, Soper, Sterman (1985)  
Ji, Ma, Yuan (2004)  
Collins (2011)

## Dihadron in $e^+e^-$

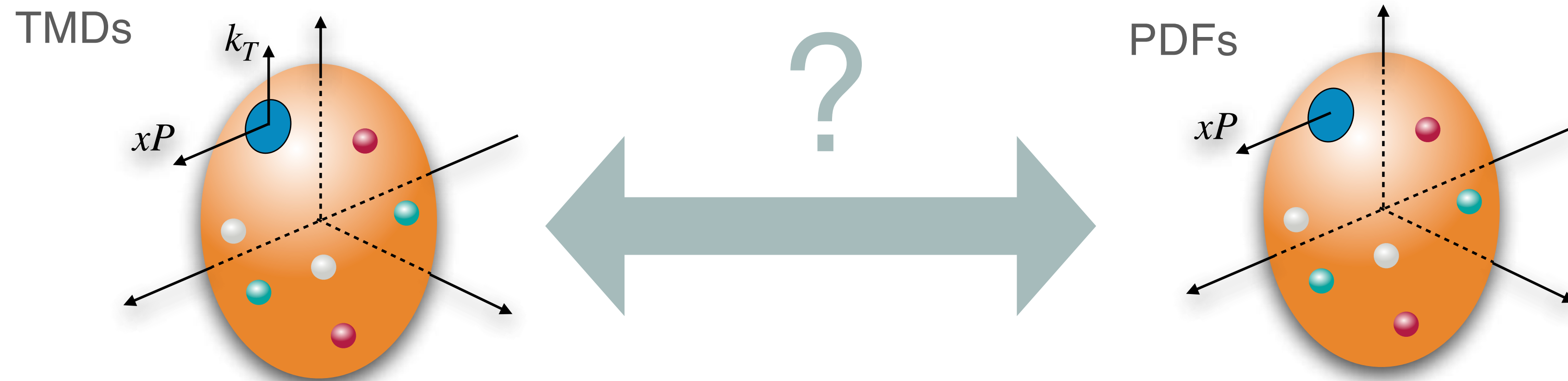
$$\sigma \sim D_{h_1/P}(z, k_T) D_{h_2/q}(z, k_T)$$



Collins, Soper (1983)  
Collins (2011)

# WHAT IS THE RELATIONSHIP?

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TMDs are related to PDFs, but how?

# WHAT IS RELATIONSHIP?

TMD handbook: <https://inspirehep.net/literature/2650019>

Both PDFs and TMDs are extracted from the data of the same type of inclusive processes (DY, SIDIS, etc) in the corresponding regions of validity of the factorization theorems.

TMDs can be related to PDFs (or other collinear functions for polarized TMDs) via Operator Product Expansion at small values\* of  $b_T$

**Known up to N3LO**

*M.A. Ebert, B. Mistlberger and G. Vita, JHEP 09 (2020) 146*

*M.-x. Luo, T.-Z. Yang, H.X. Zhu and Y.J. Zhu, JHEP 06 (2021) 115*

$$\tilde{f}_f(x, b; \mu, \zeta) = \sum_{f'} \int_x^1 \frac{dy}{y} \tilde{C}_{f \rightarrow f'}(y, b, \mu, \zeta) f_{f'}\left(\frac{x}{y}, \mu\right) + O(b^2)$$

TMDs can be related to PDFs via integral relations (TMMs)

**Studied up to N3LO**

$$\int^{|\mu|} d^2 k_T f(x, k_T; \mu, \zeta) = f(x; \mu)$$

*O. del Rio, AP, I. Scimemi, A.Vladimirov Phys.Rev.D 110 (2024)*

*A. Bacchetta, AP, Nucl.Phys.B 875 (2013) 536-551*

*M. A. Ebert, J. K. L. Michel, I. W. Stewart and Z. Sun, JHEP 07 (2022) 129*

*J. O. Gonzalez-Hernandez, T. Rainaldi, T. C. Rogers Phys.Rev.D 107 (2023) 9, 094029*

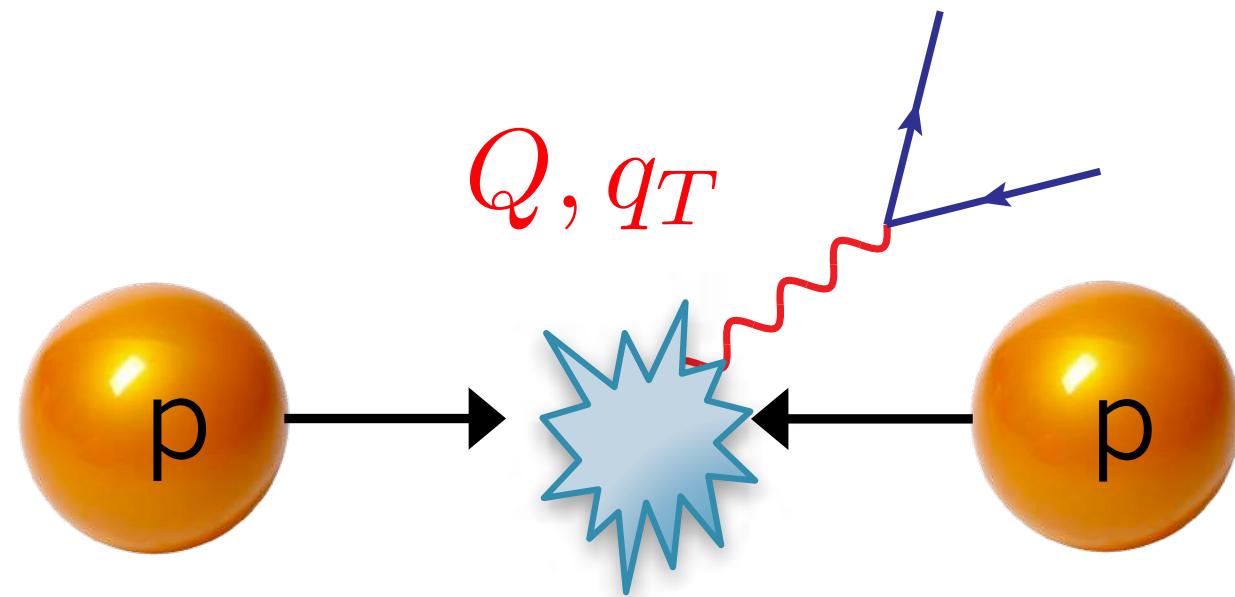
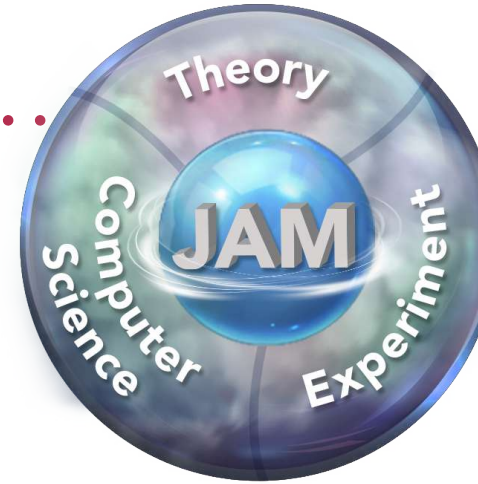
\* I will use  $b$  or  $b_T$  in this presentation as in the literature they are used to denote the same variable for TMDs

# TMD FITS OF UNPOLARIZED DATA

TMD handbook: <https://inspirehep.net/literature/2650019>

	Framework	W+Y	HERMES	COMPASS	DY	Z boson	W boson	N of points
KN 2006 <a href="#">hep-ph/0506225</a>	LO-NLL	W	✗	✗	✓	✓	✗	98
QZ 2001 <a href="#">hep-ph/0506225</a>	NLO-NLL	W+Y	✗	✗	✓	✓	✗	28 (?)
RESBOS <a href="#">resbos@msu</a>	NLO-NNLL	W+Y	✗	✗	✓	✓	✗	>100 (?)
Pavia 2013 <a href="#">arXiv:1309.3507</a>	LO-PM	W	✓	✗	✗	✗	✗	1538
Torino 2014 <a href="#">arXiv:1312.6261</a>	LO-PM	W	✓ (separately)	✓ (separately)	✗	✗	✗	576 (H) 6284 (C)
DEMS 2014 <a href="#">arXiv:1407.3311</a>	NLO-NNLL	W	✗	✗	✓	✓	✗	223
EIKV 2014 <a href="#">arXiv:1401.5078</a>	LO-NLL	W	1 (x,Q <sup>2</sup> ) bin	1 (x,Q <sup>2</sup> ) bin	✓	✓	✗	500 (?)
SIYY 2014 <a href="#">arXiv:1406.3073</a>	NLO-NLL	W+Y	✗	✓	✓	✓	✗	200 (?)
Pavia 2017 <a href="#">arXiv:1703.10157</a>	LO-NLL	W	✓	✓	✓	✓	✗	8059
SV 2017 <a href="#">arXiv:1706.01473</a>	NNLO-NNLL	W	✗	✗	✓	✓	✗	309
BSV 2019 <a href="#">arXiv:1902.08474</a>	NNLO-NNLL	W	✗	✗	✓	✓	✗	457
Pavia 2019 <a href="#">arXiv:1912.07550</a>	NNLO-N3LL	W	✗	✗	✓	✓	✗	353
SV 2019 <a href="#">arXiv:1912.06532</a>	NNLO-N3LL	W	✓	✓	✓	✓	✗	1039
MAP pion 2022 <a href="#">arXiv:2210.01733</a>	NLO-N3LL	W	✗	✗	✓	✗	✗	138
MAP 2022 <a href="#">arXiv:2206.07598</a>	NNLO-N3LL-	W	✓	✓	✓	✓	✗	2031
JAM 2023 <a href="#">arXiv:2302.01192</a>	NLO-NNLL	W	✗	✗	✓	✗	✗	608
ART 2023 <a href="#">arXiv:2305.07473</a>	N3LO-N4LL	W	✗	✗	✓	✓	✓	627
RESBOS <a href="#">arXiv:2311.09916</a>	NNLO-N3LL	W+Y	✗	✗	✓	✓	✗	384
Aslan et al 2024 <a href="#">arXiv:2401.14266</a>	NLO-NLL	W	✗	✗	✓	✗	✗	130
MAP 2025 <a href="#">arXiv:2502.04166</a>	NNLO-N3LL	W	✗	✗	✓	✓	✗	482
ART 2025 <a href="#">arXiv:2503.11201</a>	N3LO-N4LL	W	✓	✓	✓	✓	✓	1269
JAM 2025 <a href="#">arXiv:2510.13771</a>	NLO-NNLL	W	✗	✗	✓	✓	✗	436+4279
KSZZ 2026 <a href="#">arXiv:2604.14133</a>	N3LO-N4LL	W	✗	✗	✓	✓	✗	465

Is it possible to combine the full body of data used in the extractions of TMDs and PDFs, respectively, and to perform a global QCD analysis of both distributions simultaneously?



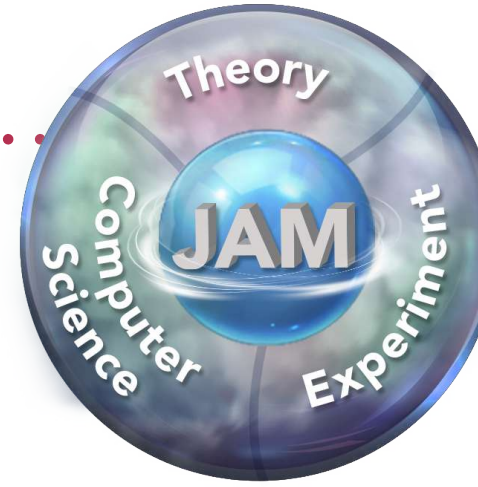
At small values of  $q_T$  Drell-Yan cross section reads:

$$\frac{d^3\sigma}{dQ^2 dy dq_T^2} = \frac{4\pi^2 \alpha_{\text{em}}^2}{9Q^2 s} \mathcal{P} \sum_j c_j^2(Q) H_{j\bar{j}}^{\text{DY}}(\mu, Q) \int \frac{d^2\mathbf{b}_T}{(2\pi)^2} e^{i\mathbf{q}_T \cdot \mathbf{b}_T} \tilde{f}_{q_j/\mathcal{N}}(x_1, b_T; \mu, \zeta) \tilde{f}_{\bar{q}_j/\mathcal{N}}(x_2, b_T; \mu, \zeta),$$

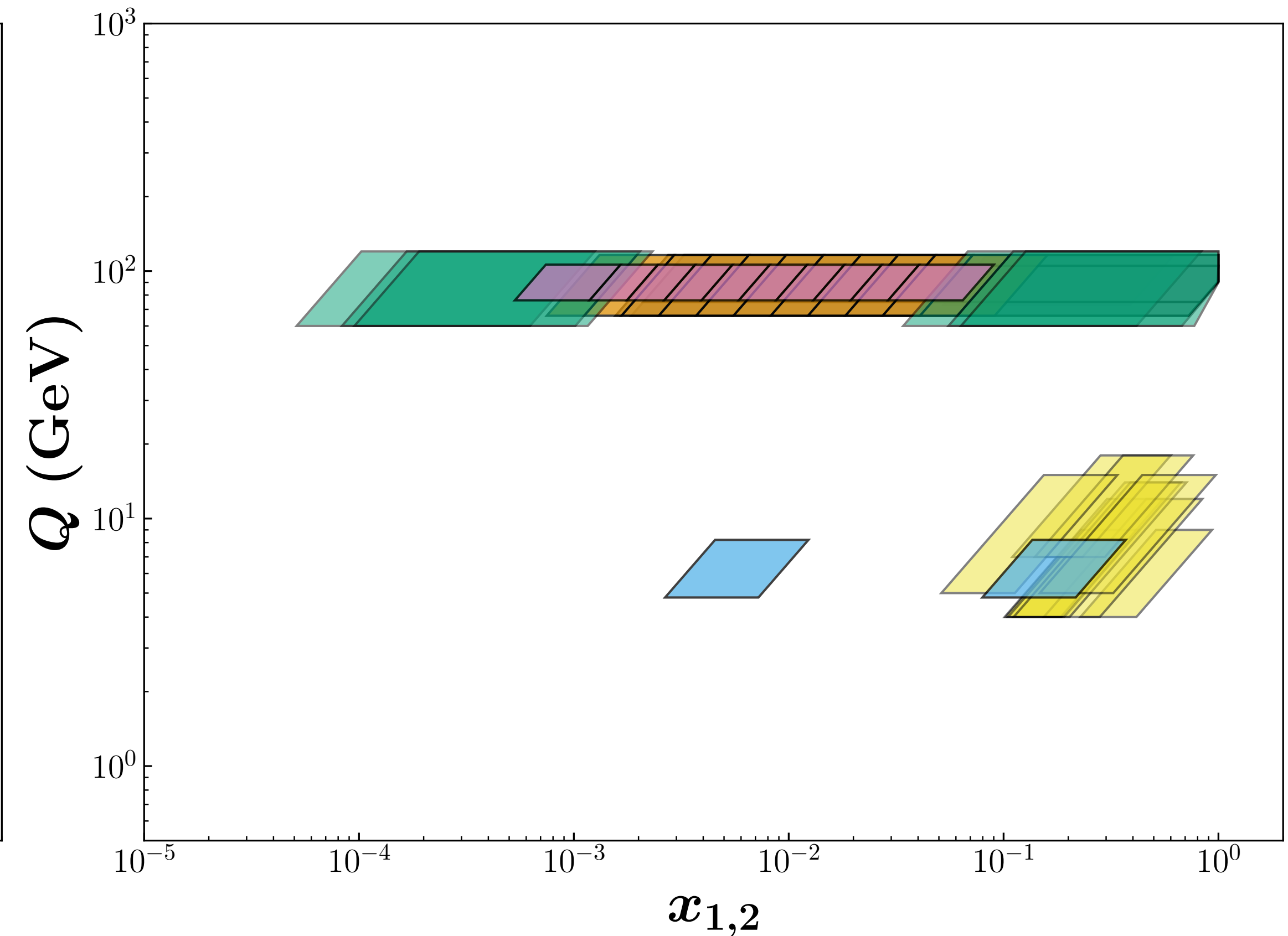
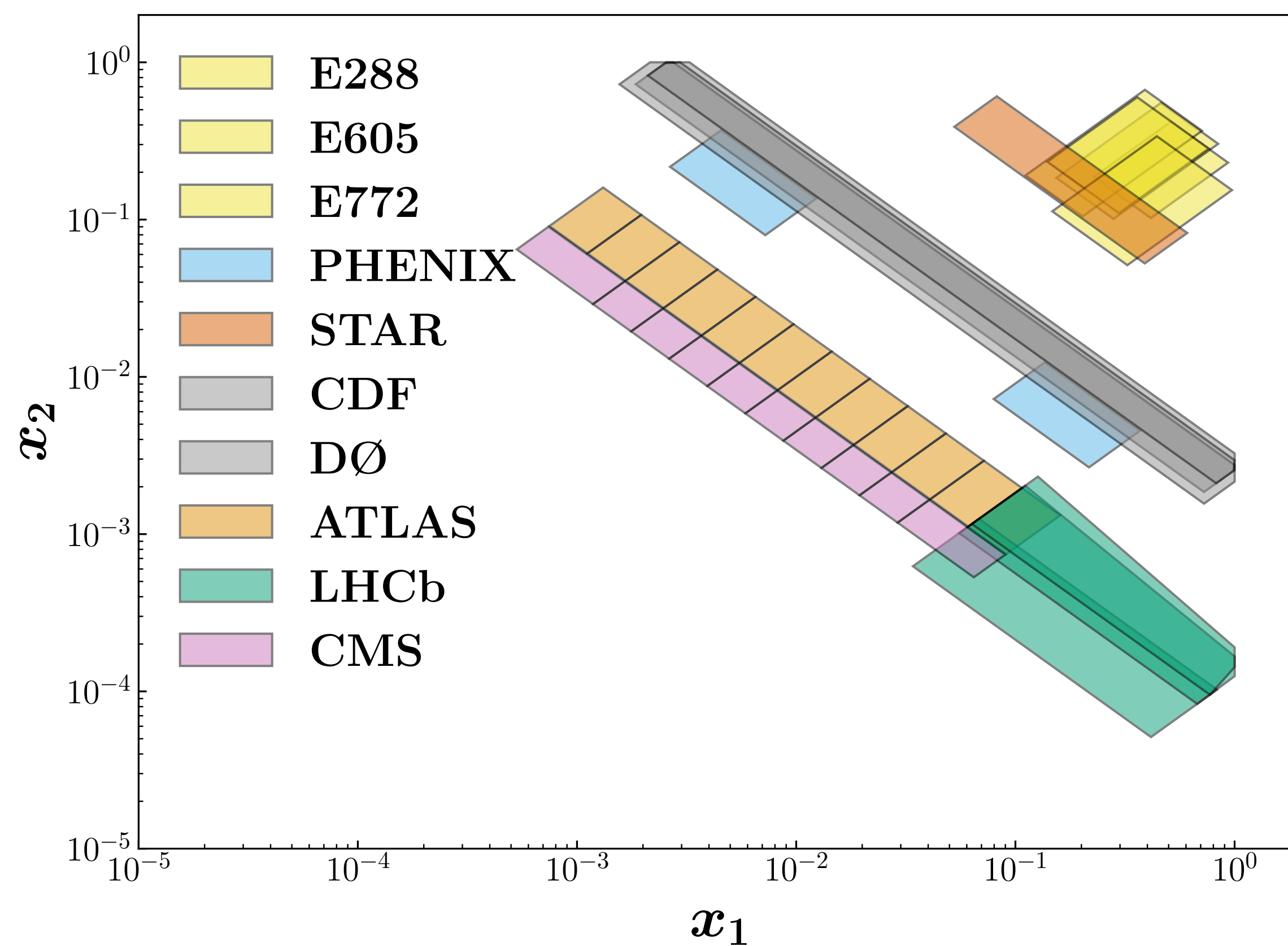
Fiducial factor
Electroweak couplings
Hard factor
TMD for the beam
TMD for the target

To match the NLO precision of collinear PDFs by JAM Collaboration we use **NLO + N2LL** accuracy in our analysis: PDFs are at NLO  $\mathcal{O}(\alpha_S^2)$ , hard factor at  $\mathcal{O}(\alpha_S)$ , Operator Product Expansion at  $\mathcal{O}(\alpha_S)$ , anomalous dimension of TMD at  $\mathcal{O}(\alpha_S^2)$  and, Collins Soper-Kernel at  $\mathcal{O}(\alpha_S^2)$ , and Cusp anomalous dimension at  $\mathcal{O}(\alpha_S^3)$ . We use  $\zeta$  prescription for TMDs

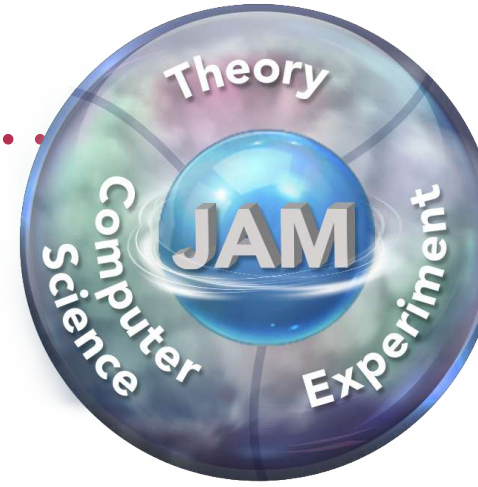
# DRELL-YAN DATASETS



Data from low energy fixed target DY (20-40 GeV). Collider data from RHIC, PHENIX and STAR at 200 and 510 GeV, Fermilab, CDF and D0 at 1.8 TeV, 1.96 TeV, and LHC data from CMS, LHCb, ATLAS at 7, 8, 13 TeV

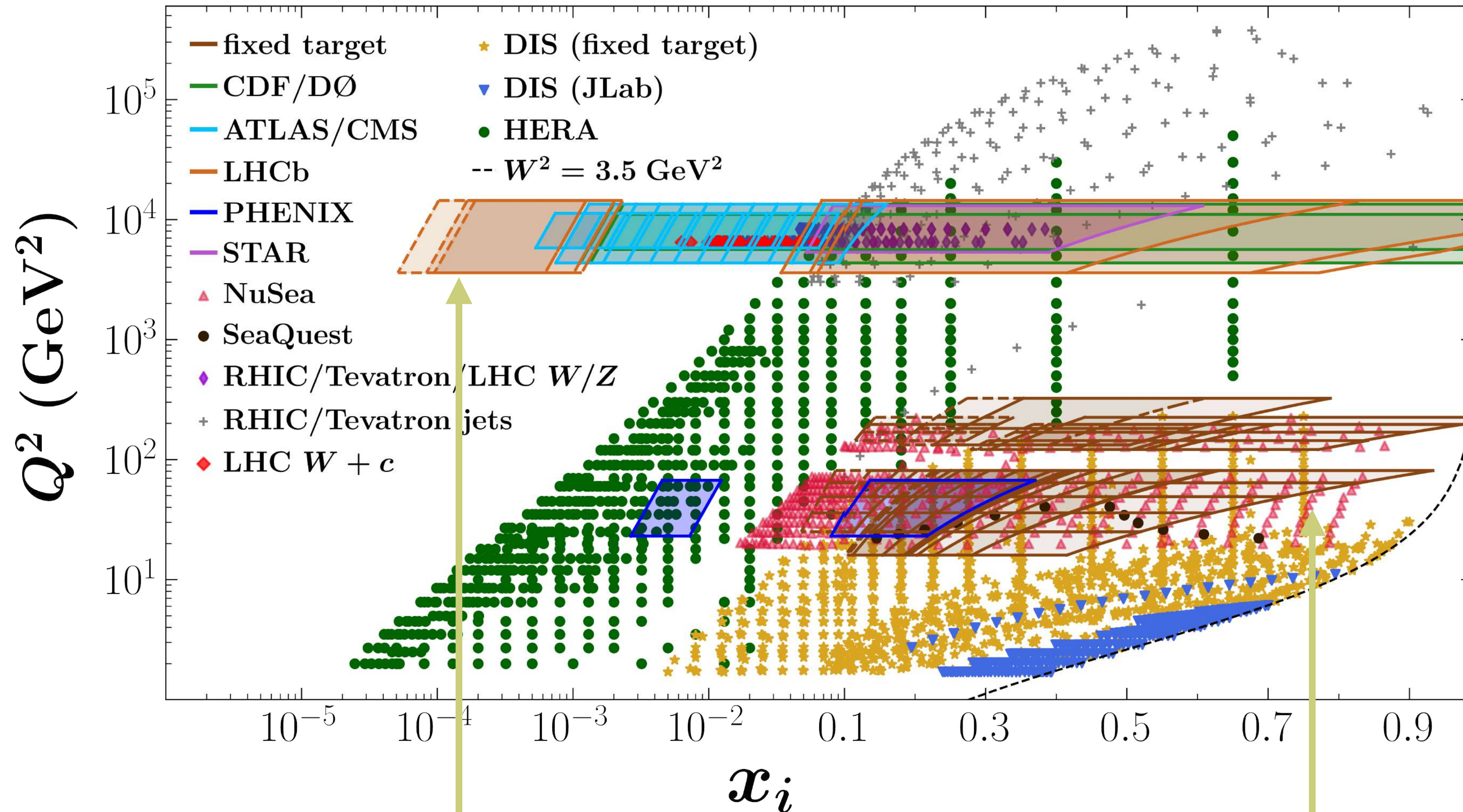


# DATASETS AND KINEMATICS



Symbols are for the data relevant to collinear extraction, regions are for TMD related data

JAM Collaboration (PDF Analysis Group) Trey Anderson et al. *Phys.Rev.D* 112 (2025) 9, 094011  
Jefferson Lab Angular Momentum (JAM) C. Cocuzza et al. *Phys.Rev.Lett.* 127 (2021) 24, 242001



Precise data at small and intermediate x

Data at large-x

# NON PERTURBATIVE MODELS FOR TMDS AND PDFS

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Fit  $\lambda_1$  and  $\lambda_2$  in this functional form for these flavors:  $u, d, \bar{u}, \bar{d}$  and  $sea = s, \bar{s}, c, \bar{c}, b, \bar{b}$ . Total of 10 parameters

$$F^f(x, b) = \frac{1}{\cosh\left(\lambda_1^f(1-x) + \lambda_2^f x\right)b}$$

*Moos, Scimemi, Vladimirov, Zurita: JHEP 05 (2024)*

Collinear PDFs are parametrized and the input scale to the charm quark mass,  $m_c = 1.28$  GeV for  $u_v, d_v, \bar{u}, \bar{d}, s, \bar{s}, g$

$$f(x; m_c) = Nx^\alpha(1-x)^\beta(1 + \gamma\sqrt{x} + \delta x)$$

*JAM Collaboration (PDF Analysis Group), T. Anderson et al. Phys.Rev.D 112 (2025) 9, 09401*

heavy flavors are generated via evolution. We have 35 total parameters for collinear PDFs

# NON PERTURBATIVE MODEL FOR COLLINS-SOPER KERNEL

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Collins-Soper kernel can be expressed as

$$\mathcal{D}(b_T; \mu) = \mathcal{D}_{\text{pert}}(b^*; \mu^*) + \int_{\mu^*}^{\mu} \frac{d\mu'}{\mu'} \Gamma(\mu') + \mathcal{D}_{\text{NP}}(b_T)$$

*Moos, Ignazio Scimemi, Alexey Vladimirov, Zurita: JHEP 05 (2024)*

For the non perturbative Collins-Soper kernel we use

$$\mathcal{D}_{\text{NP}} = bb^* \left( c_0 + c_1 \ln \frac{b^*}{b_{\text{NP}}} \right), \text{ where } b^*(b_T) = \frac{b_T}{\sqrt{1 + b_T^2/B_{\text{NP}}^2}}, \mu^* = \frac{2e^{-\gamma_E}}{b^*(b_T)} \text{ and}$$

$$B_{\text{NP}} = 1.5 \text{ GeV}^{-1}$$

The perturbative part of the kernel  $\mathcal{D}_{\text{pert}}(b^*; \mu^*)$  is calculated at  $\mathcal{O}(\alpha_S^2)$

This parametrization ensures evolution and perturbative convergence,  $c_0, c_1$  are free parameters.

We perform the analysis in four steps:

- **baseline PDF** — an analysis of PDFs only
- **baseline TMD** — an analysis of TMDs with fixed PDFs. Low energy DY, RHIC, and Tevatron data
- **baseline TMD+LHC** — an analysis of TMDs with fixed PDFs, **baseline TMD** data plus LHC data
- **JAM25<sub>TMD+PDF</sub>** — combined analysis of TMDs and PDFs

In each step we generate around 1000 replicas and perform Bayesian inference

The following cuts are used:

$$\frac{q_T}{Q} < 0.2, W^2 > 3.5 \text{ (GeV}^2\text{)}, Q > m_c = 1.28 \text{ (GeV)}$$

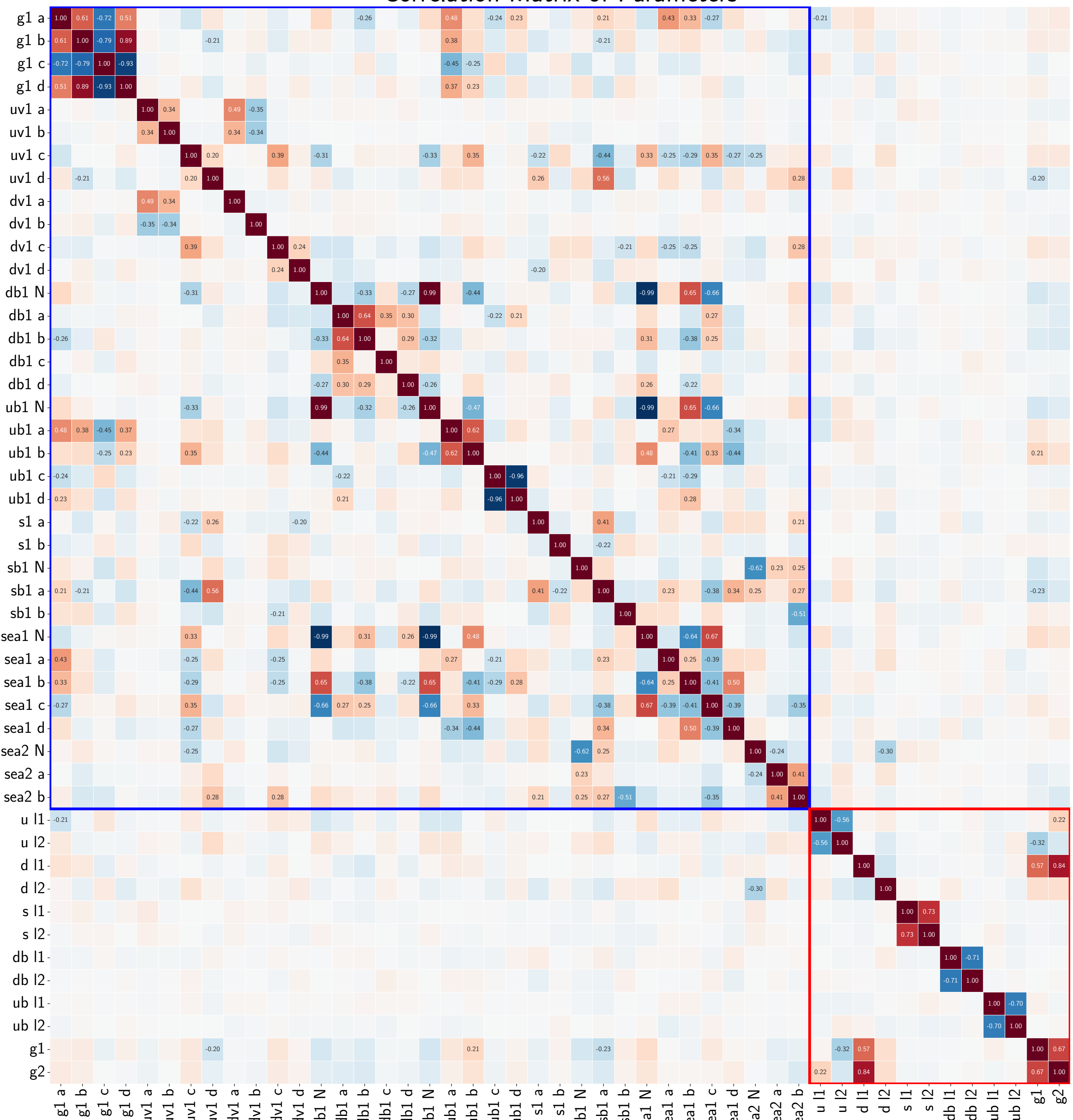
# RESULTS



Collinear				
Process	Experiment	$N_{pts}$	$\chi^2/N_{pts}$ (Z-score)	
			PDF only	<b>TMD+PDF</b>
DIS	fixed target	2501	1.02 (0.82)	<b>1.02 (0.59)</b>
	HERA	1185	1.27 (6.01)	<b>1.27 (6.11)</b>
DY	E866, 906	205	1.27 (2.54)	<b>1.26 (2.50)</b>
$W \ell$ a.	LHC, RHIC	70	0.80 (-1.20)	<b>0.78 (-1.35)</b>
$W$ ch. a.	CDF, D0	27	1.12 (0.51)	<b>1.13 (0.53)</b>
$Z$ rap.	CDF, D0	56	1.09 (0.55)	<b>1.11 (0.60)</b>
Inc. jets	CDF, D0, STAR	198	1.00 (0.00)	<b>0.98 (-0.13)</b>
$W + c$	ATLAS, CMS	37	0.65 (-1.66)	<b>0.62 (-1.84)</b>
Total		4279	1.10 (4.31)	<b>1.09 (4.14)</b>
TMD				
Process	Experiment	$N_{pts}$	TMD only	<b>TMD+PDF</b>
DY	E288, 605, 772	224	1.36 (3.44)	<b>1.31 (3.02)</b>
	CDF, D0	80	1.06 (0.45)	<b>1.11 (0.71)</b>
	STAR, PHENIX	12	1.15 (0.47)	<b>1.20 (0.60)</b>
	ATLAS	30	2.07 (3.29)	<b>1.78 (2.55)</b>
	CMS	64	1.18 (1.03)	<b>0.92 (-0.39)</b>
	LHCb	26	0.53 (-1.97)	<b>0.50 (-2.12)</b>
Total		436	1.27 (3.72)	<b>1.20 (2.76)</b>
<b>Total</b>		<b>4715</b>		<b>1.10 (4.79)</b>

Combined TMD and PDF extraction improves on the quality of description of the data while keeping collinear datasets described well.

# Correlation Matrix of Parameters

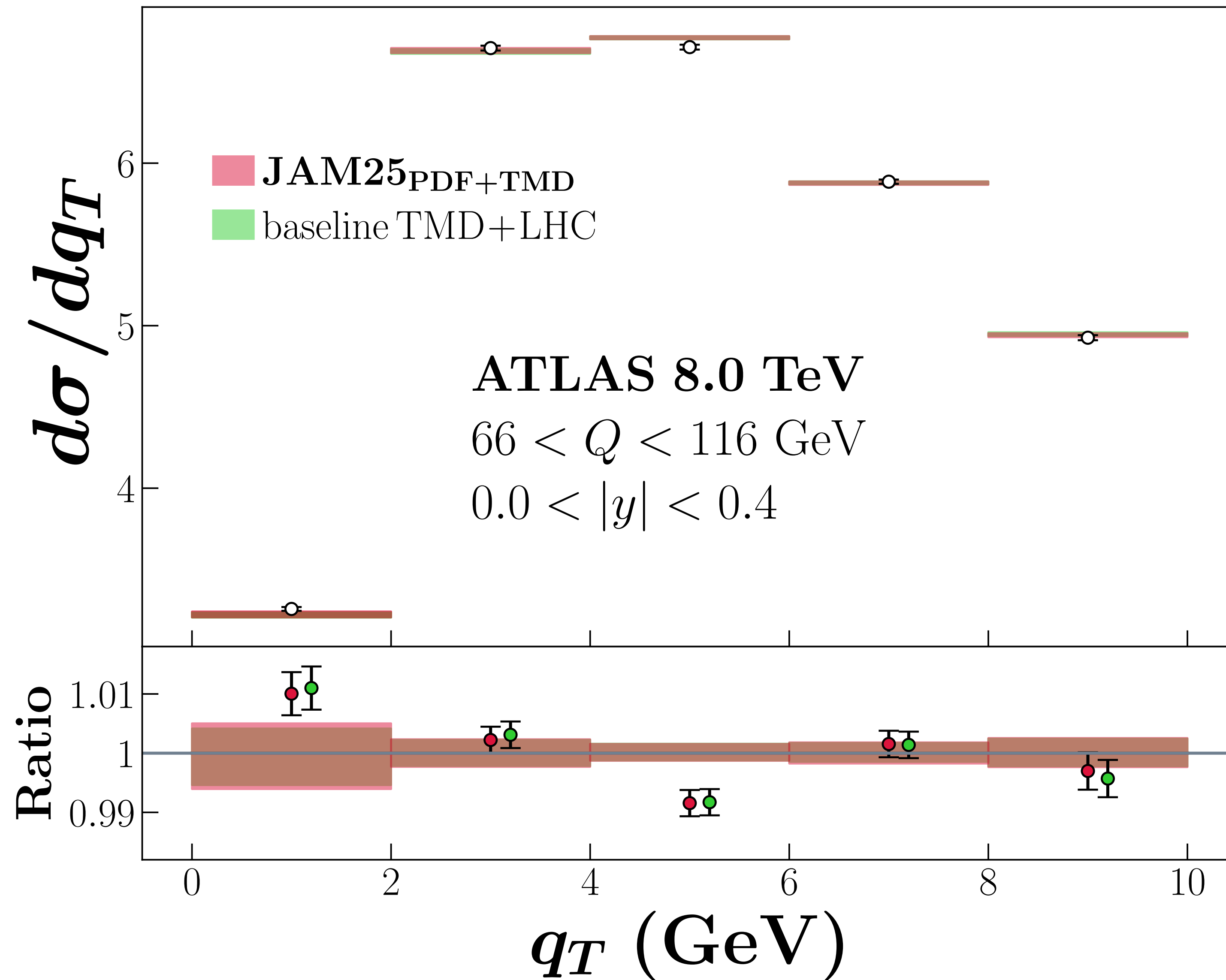
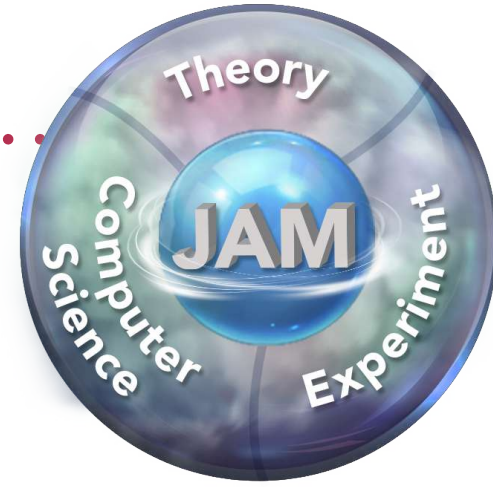


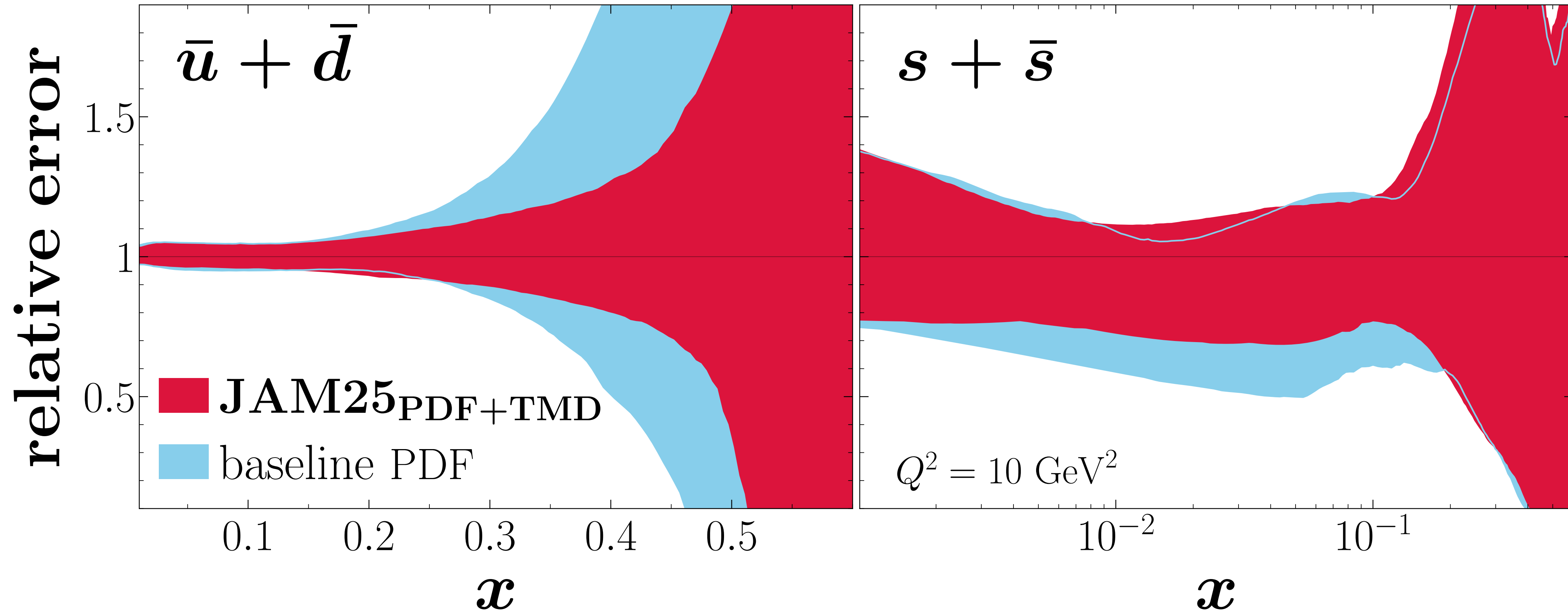
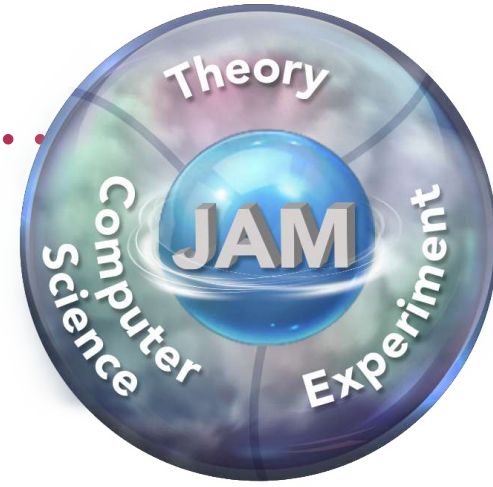
The analysis does not find strong correlation of PDF (blue rectangle) and TMD (red rectangle) parameters

Pearson coefficients are shown

$$r = \frac{\text{Cov}(X, Y)}{\sigma_X \sigma_Y} \in [-1, 1]$$

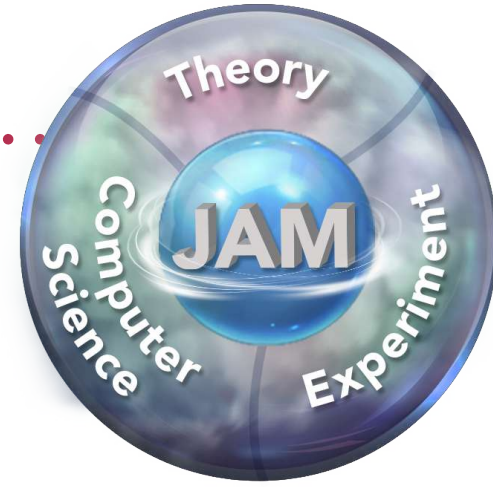




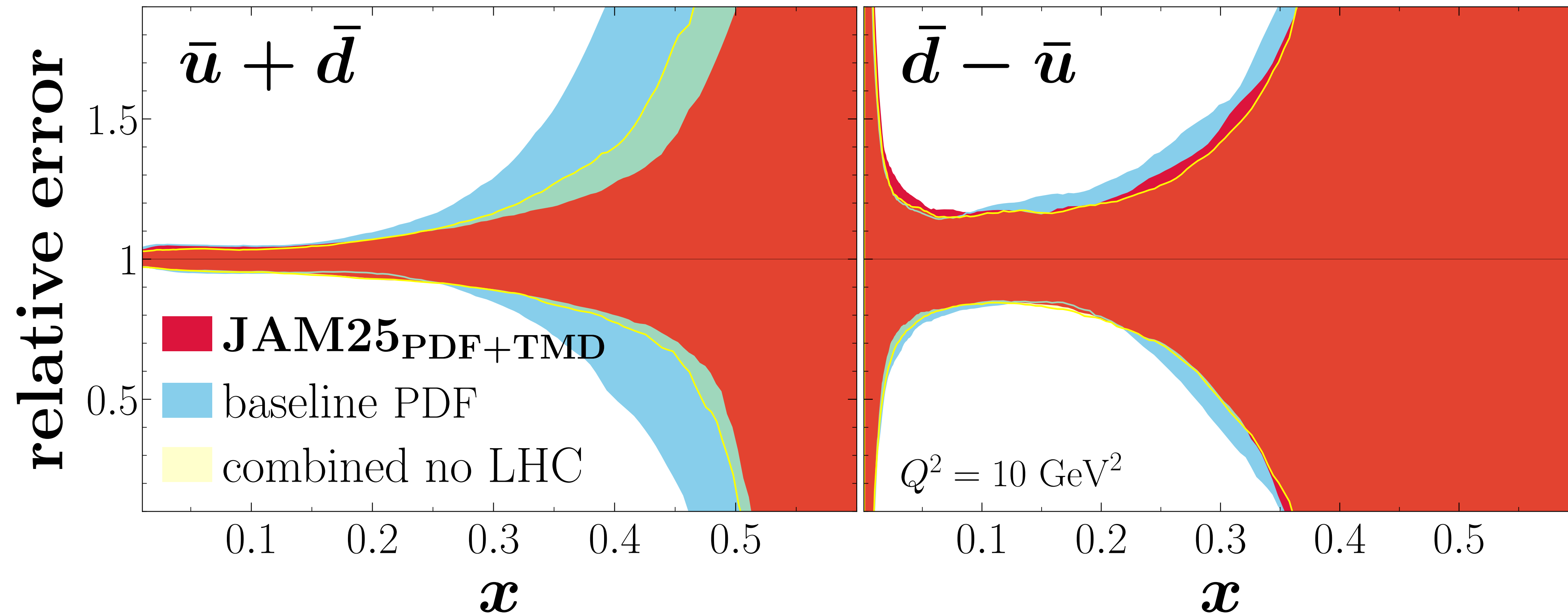


The combined analysis results in the improvement for  $\bar{d} + \bar{u}$  in large  $x$  region due to the inclusion of fixed target DY data and improvement of  $s + \bar{s}$  due to the inclusion of the LHC data.

Shown here is the 95% confidence interval.



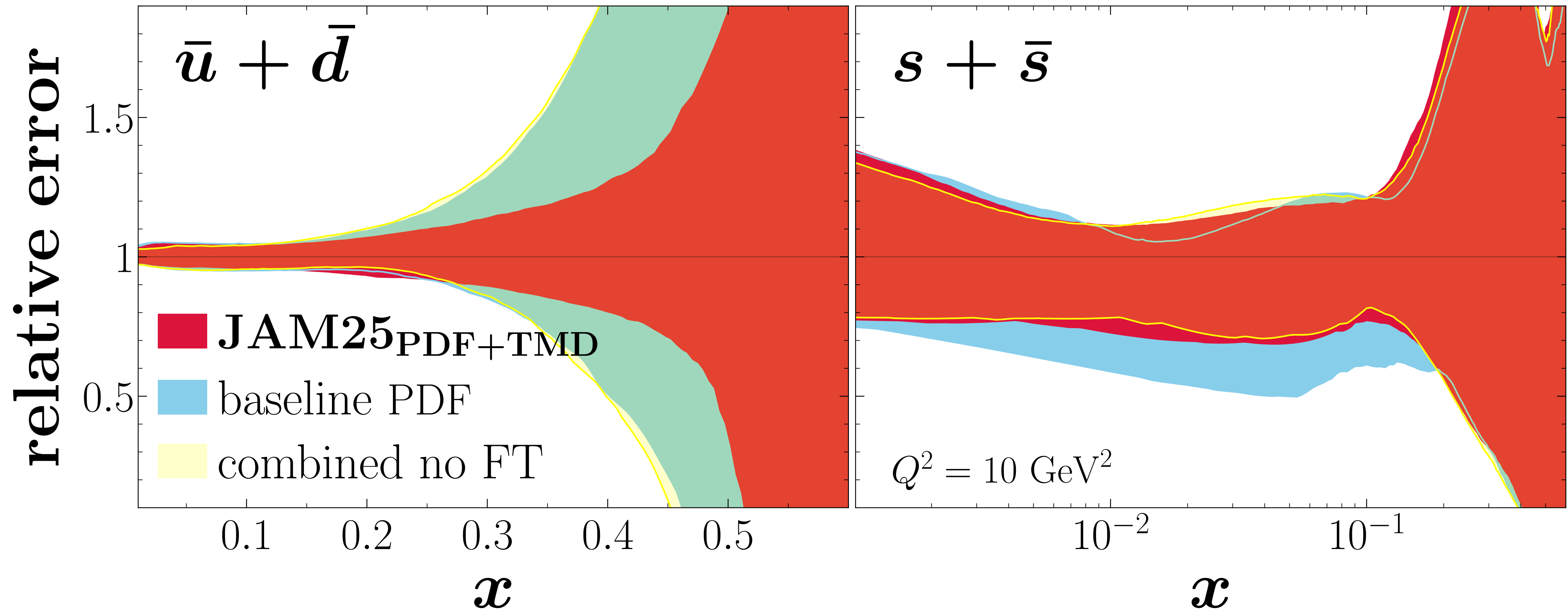
## Exclude LHC data to estimate the impact of fixed target DY



The combined analysis results in the improvement for  $\bar{d} + \bar{u}$  in large  $x$  region due to the inclusion of fixed target DY consistent with the isoscalar nature of the target. Notice that  $\bar{d} - \bar{u}$  is not impacted.

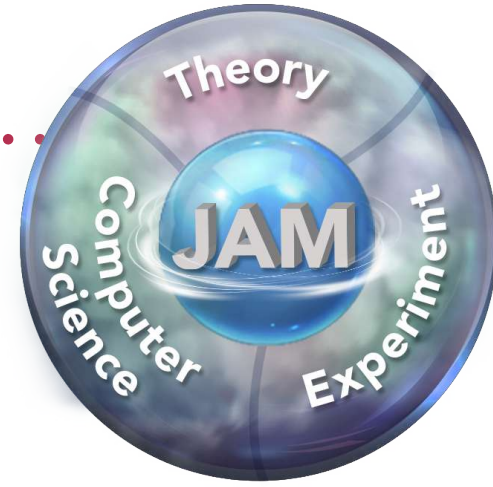


- Exclude fixed target DY data to estimate the impact of LHC data



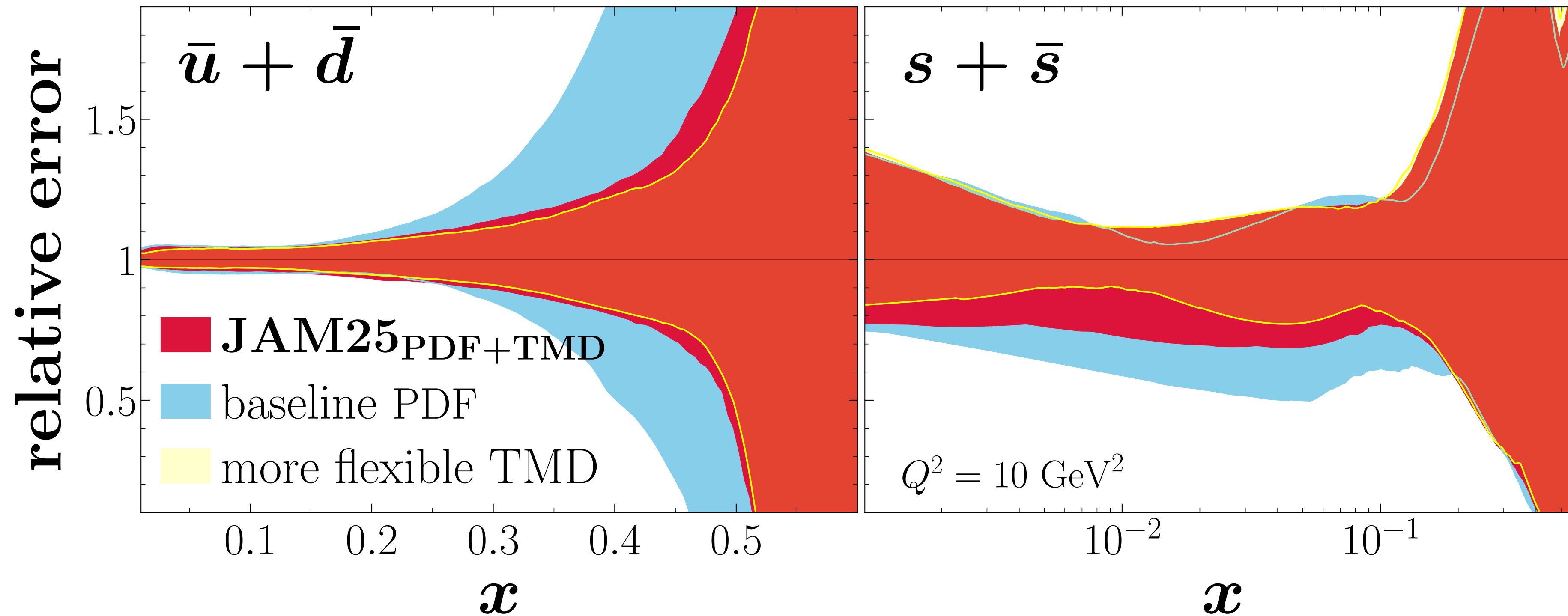
The improvement of  $s + \bar{s}$  due to the inclusion of the LHC data.

# RESULTS COLLINEAR PDFS



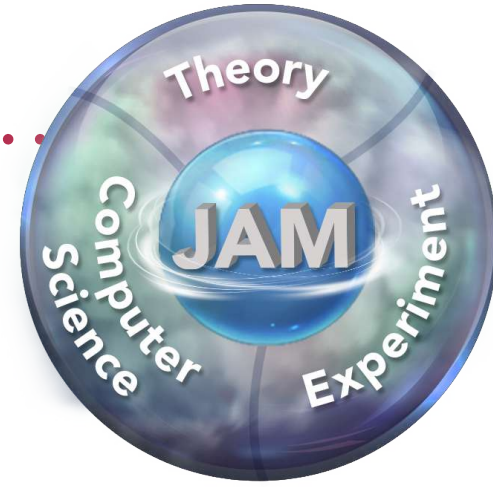
Maybe the rigidity of TMD parametrizations make this improvement?

$$F^f(x, b) = \frac{1}{\cosh\left(\lambda_1^f(1-x)^{\lambda_f^3} + \lambda_2^f x\right)b}$$

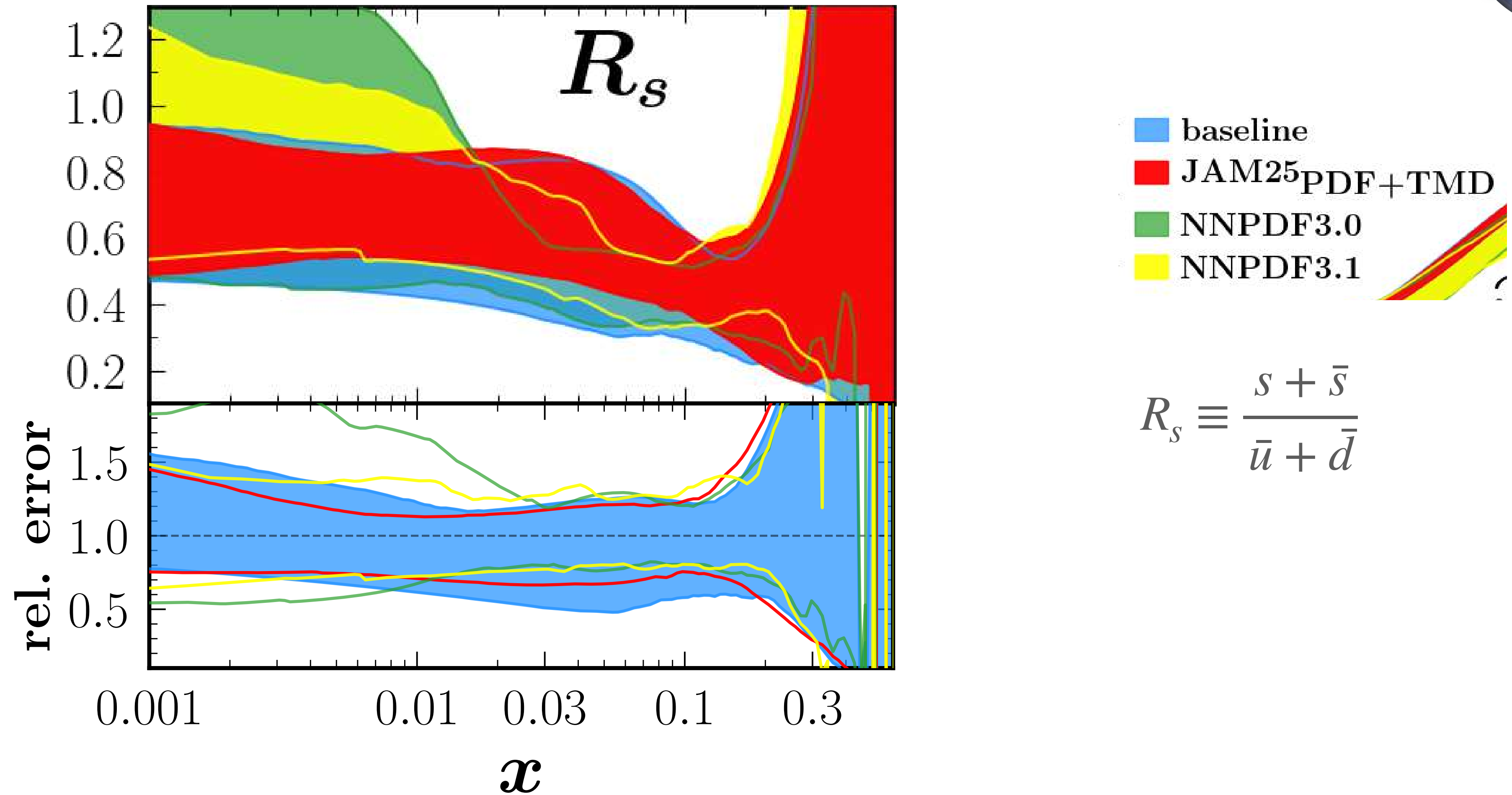


No strong influence of results on TMD parametrization.

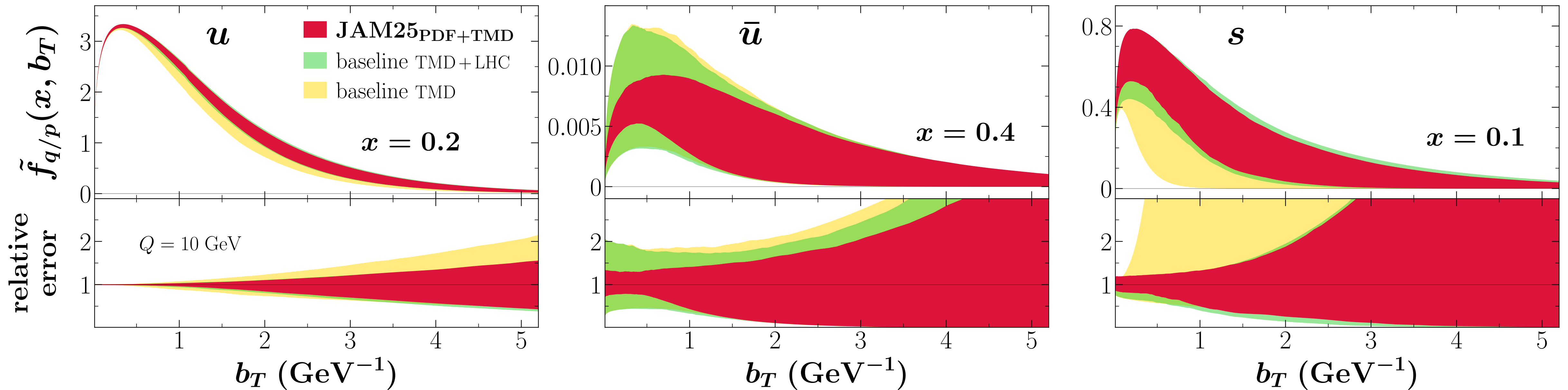
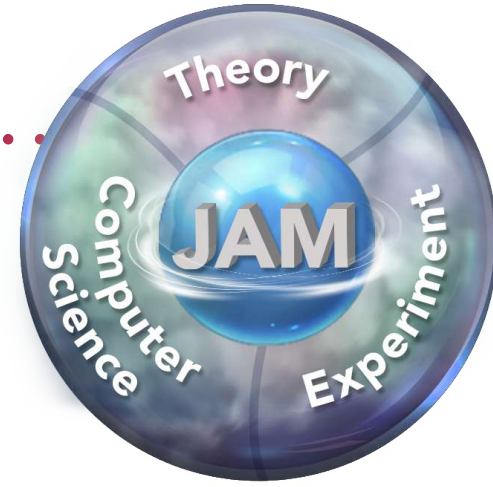
# RESULTS COLLINEAR PDFS



Is the improvement consistent with findings of the analyses that include LHC integrated data?

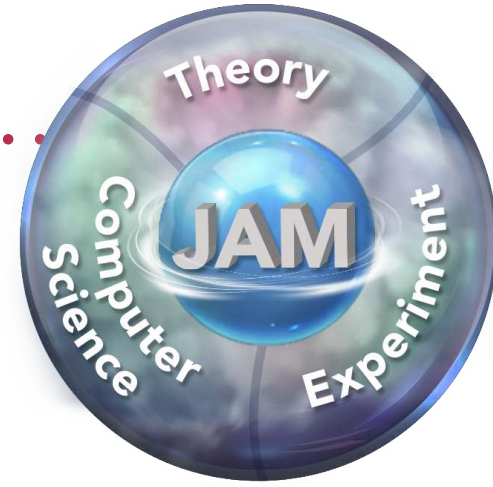


The reduction in uncertainty is consistent with the inclusion of the collinear LHC data as observed by NNPDF. This speaks to the consistency of the two independent theoretical frameworks.

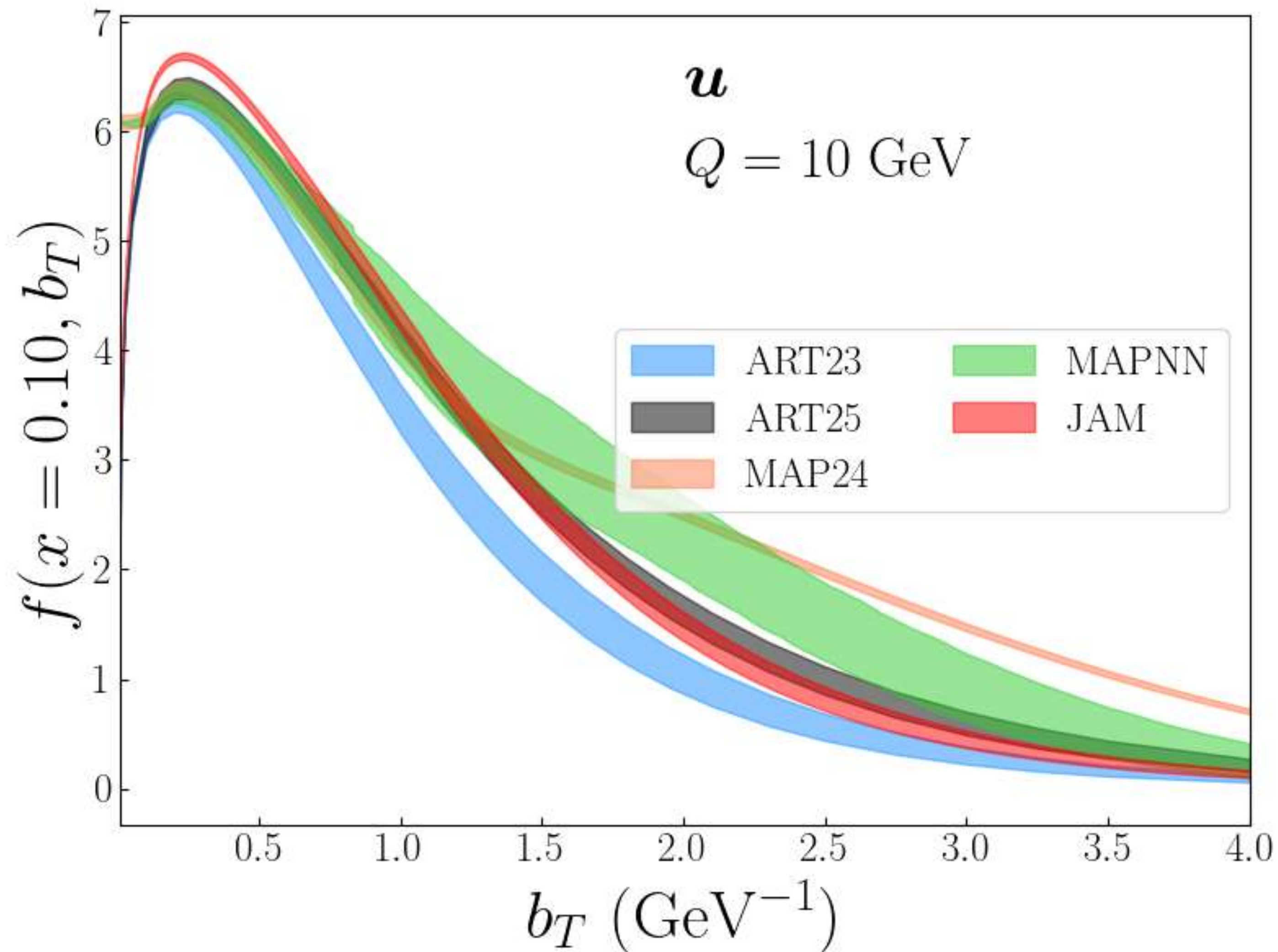


Combined fit results in the improvement of the extraction of TMDs. The LHC data are important for  $u, d, s$  quarks. Fixed target DY data improves  $\bar{u}, \bar{d}$  quarks at large  $x$

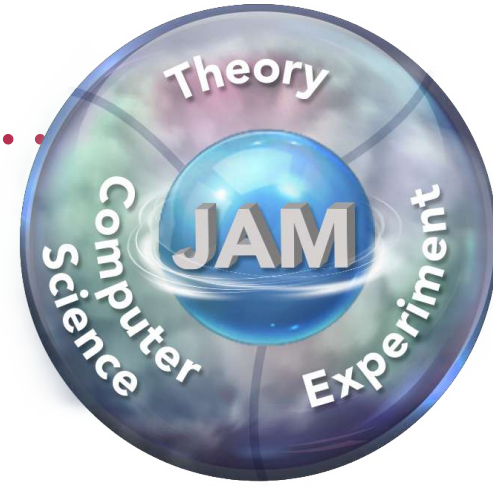
# RESULTS TMDS



Our results are in agreement with results from other groups.

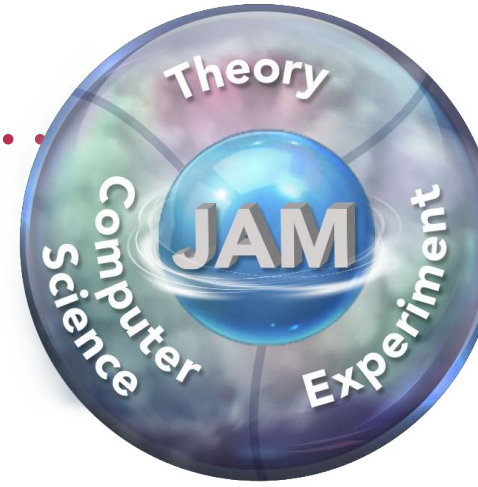


ART23: V. Moos, I. Scimemi, A. Vladimirov, P. Zurita *JHEP* 05 (2024) 036  
ART25: V. Moos, I. Scimemi, A. Vladimirov, P. Zurita *e-Print*: 2503.11201  
MAP24: MAP Collaboration, A. Bacchetta, *JHEP* 08 (2024) 232  
MAPNN: MAP Collaboration, A. Bacchetta, *Phys.Rev.Lett.* 135 (2025) 2, 021904



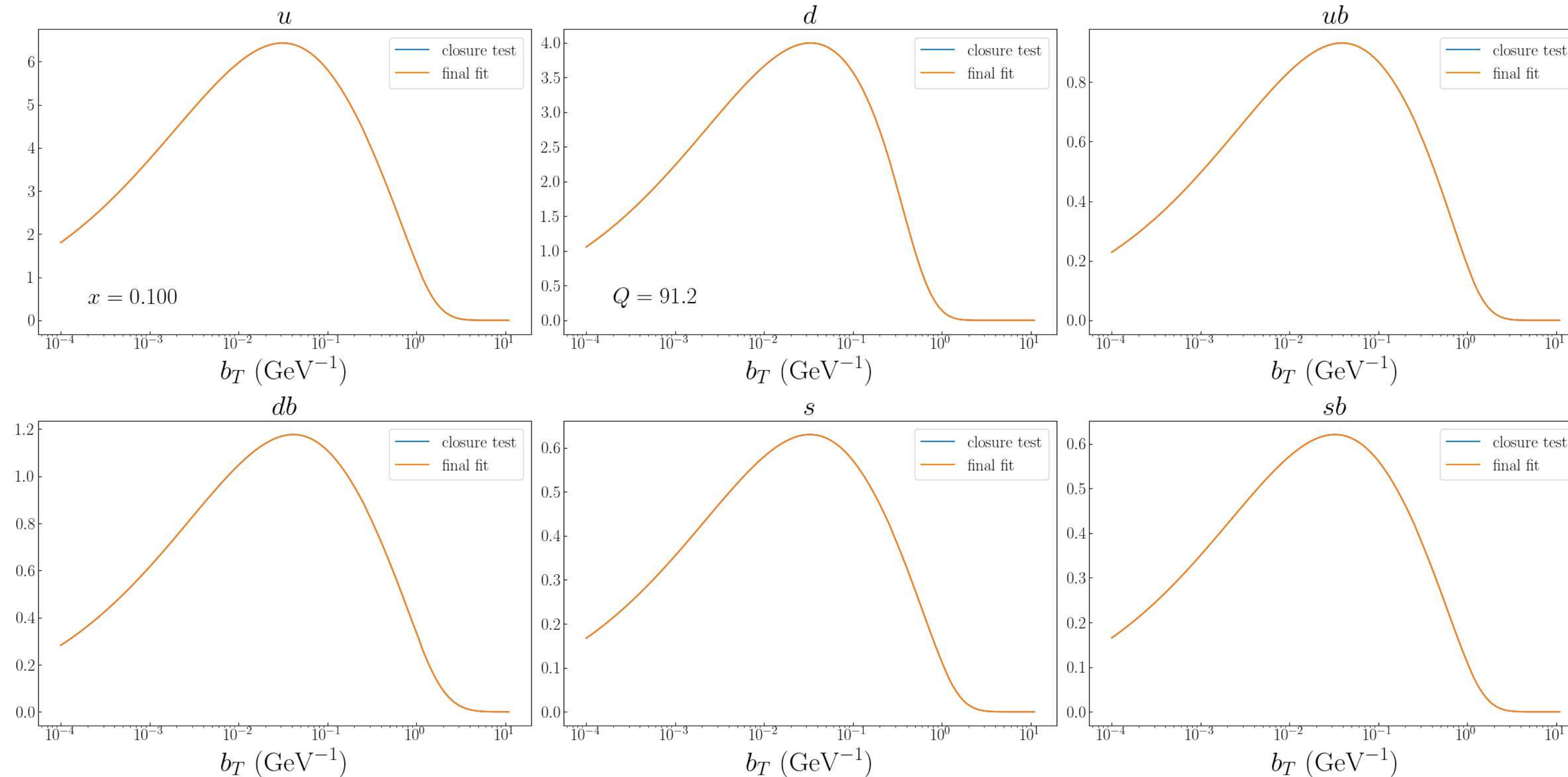
We perform *closure* tests to assess robustness of the framework:

- Fix PDFs and recover TMDs using TMD pseudo data only
- Fix TMDs and recover PDFs using pseudo PDF and TMD data
- Open all and recover TMDs and PDFs using pseudo PDF and TMD data

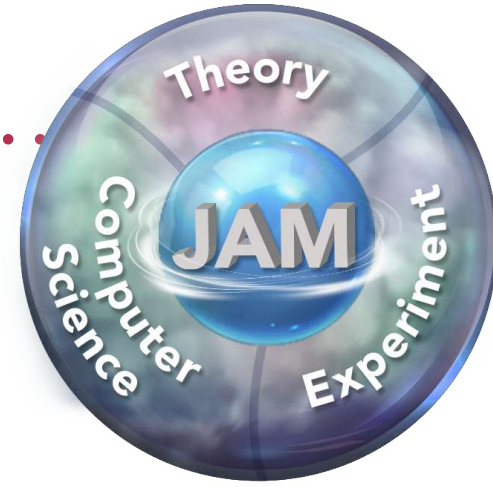


We perform closure tests to assess robustness of the framework:

- Fix PDFs and recover TMDs using TMD pseudo data only

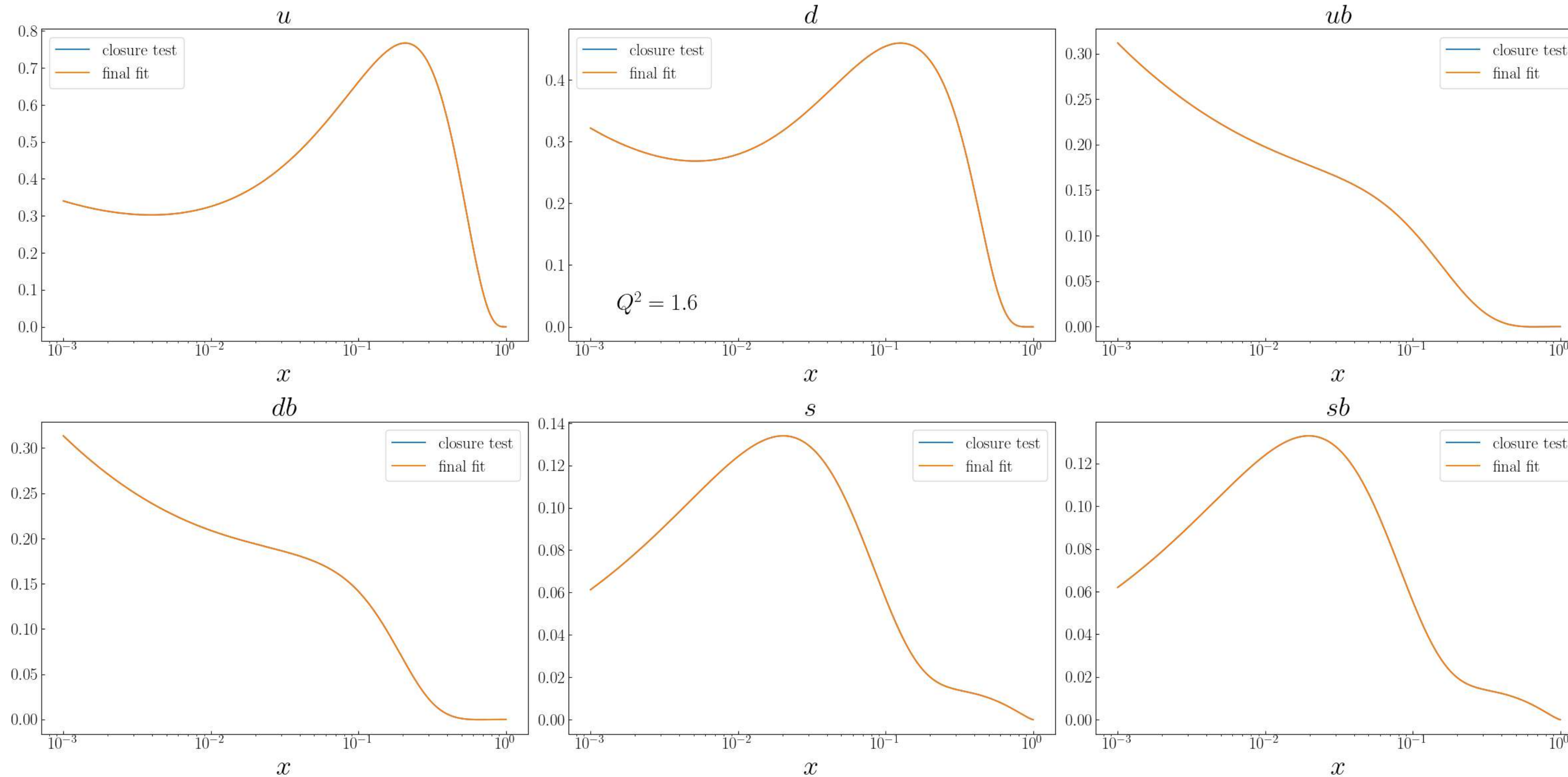


Perfect consistency in TMD sector

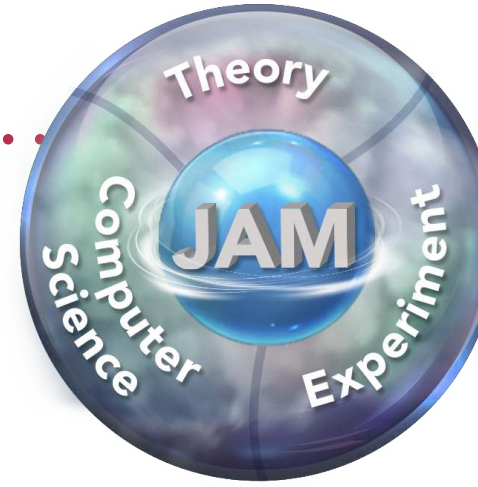


We perform closure tests to assess robustness of the framework:

- Fix TMDs and recover PDFs using pseudo PDF and TMD data

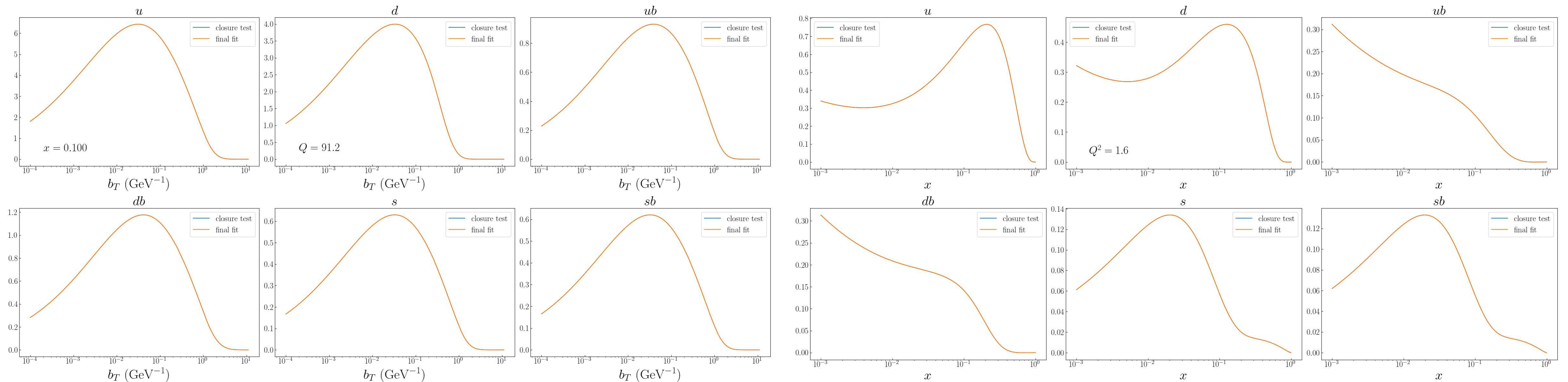


Perfect consistency in PDF sector



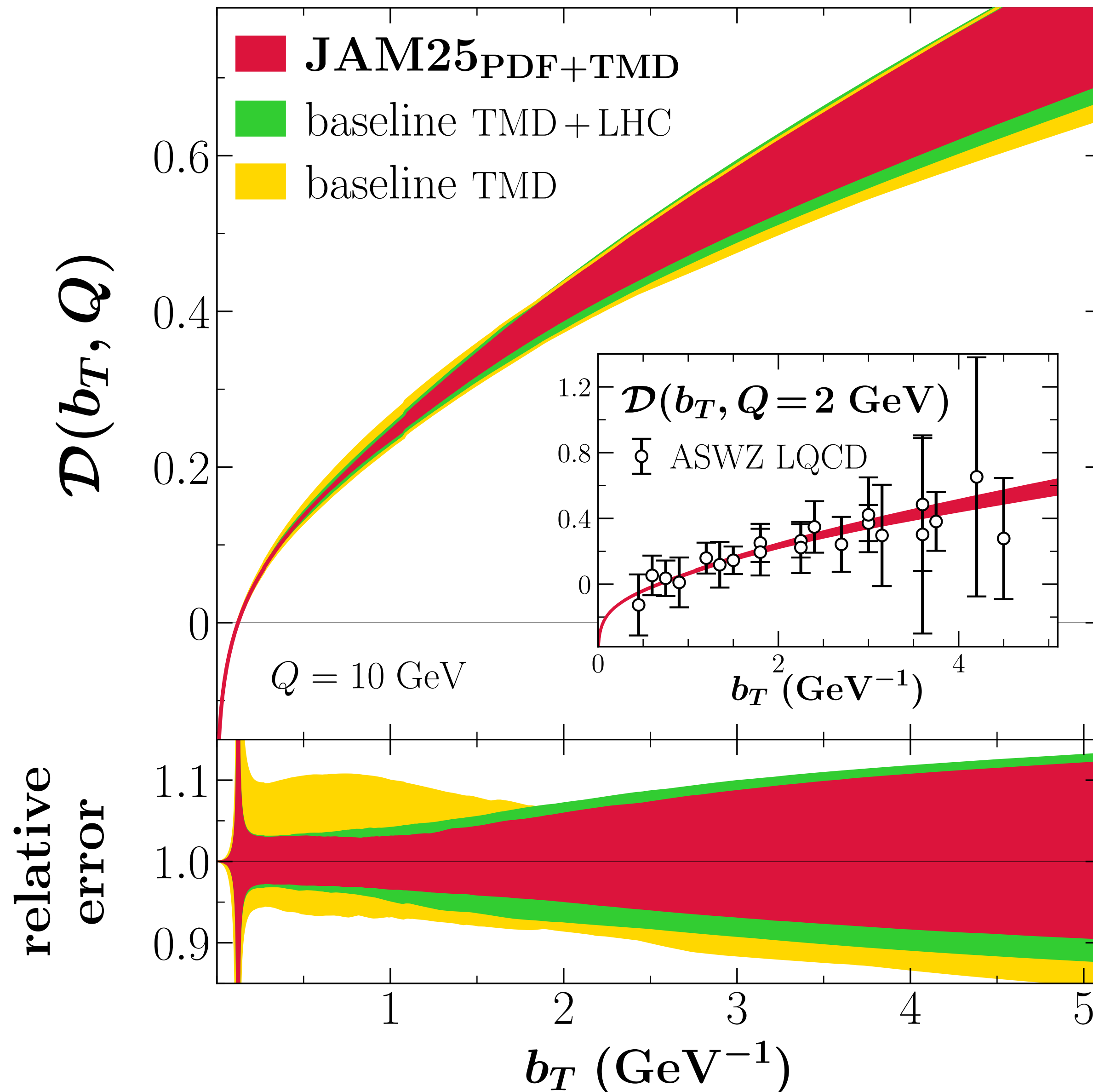
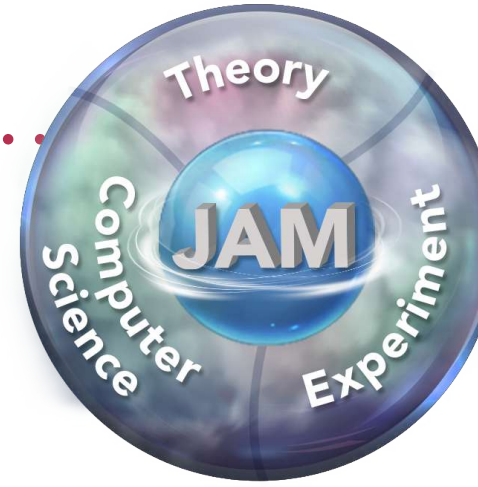
We perform closure tests to assess robustness of the framework:

- Open all and recover TMDs and PDFs using pseudo PDF and TMD data



We demonstrated consistency across these three complementary closure tests and absence of biases hidden in the restricted tests therefore strengthening the reliability of our framework.

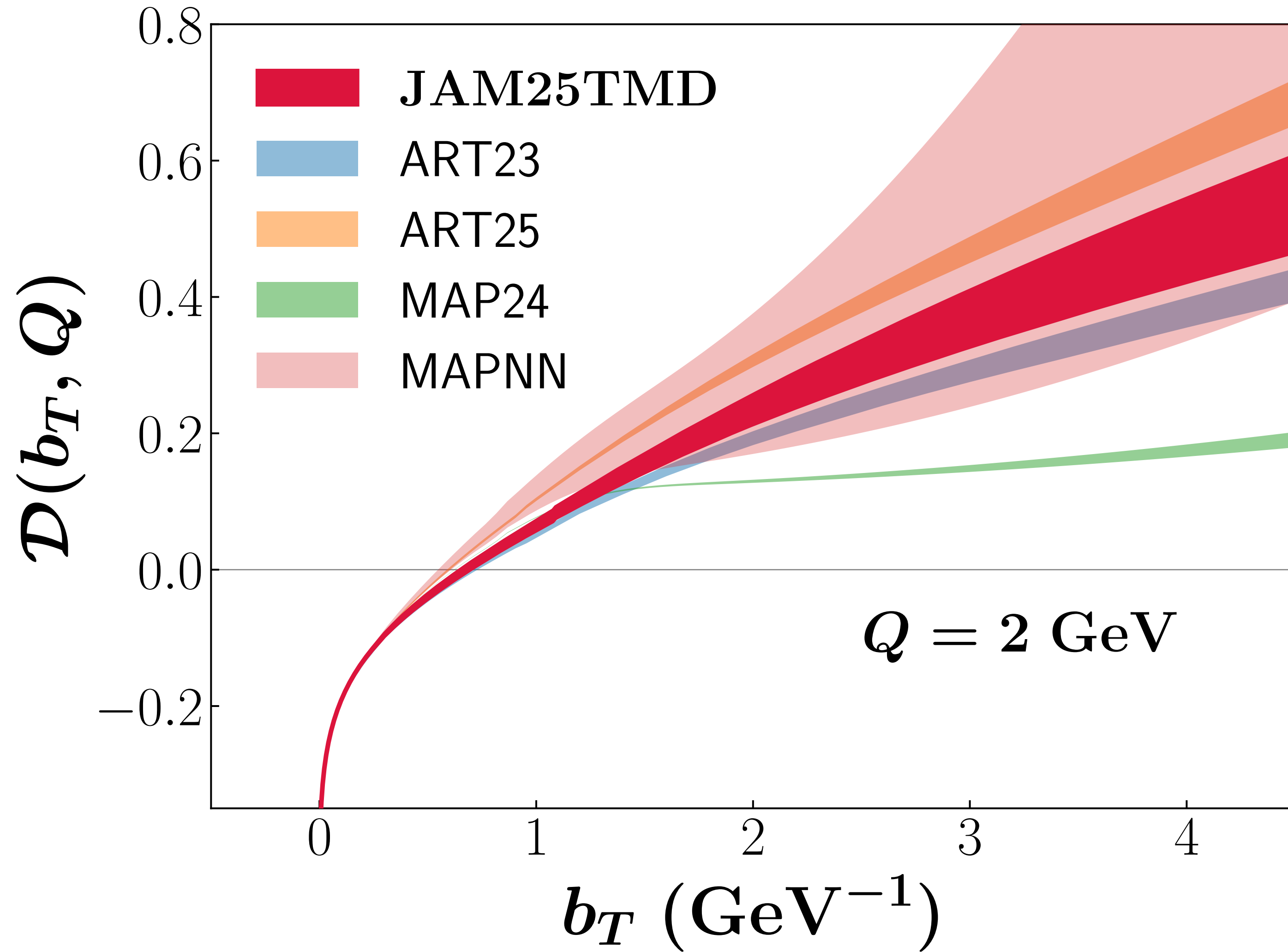
# RESULTS COLLINS-SOPER KERNEL



Collins Soper kernel is an important ingredient of the analysis. The LHC data is particularly impactful on its extraction.

Compared to lattice results from

*A. Avkhadiev, P. E. Shanahan, M. L. Wagman, and Y. Zhao, Phys. Rev. Lett. 132, 231901 (2024)*



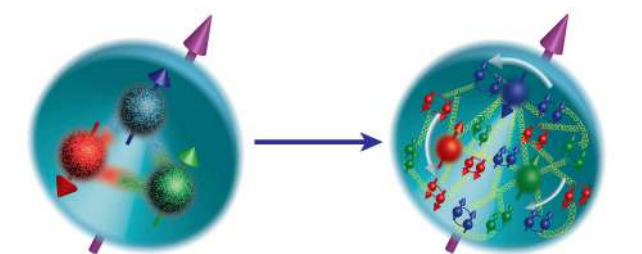
Collins Soper kernel compares well to existing extractions.

ART23: V. Moos, I. Scimemi, A. Vladimirov, P. Zurita *JHEP* 05 (2024) 036  
ART25: V. Moos, I. Scimemi, A. Vladimirov, P. Zurita e-Print: 2503.11201  
MAP24: MAP Collaboration, A. Bacchetta, *JHEP* 08 (2024) 232  
MAPNN: MAP Collaboration, A. Bacchetta, *Phys.Rev.Lett.* 135 (2025) 2, 021904

# CONCLUSIONS

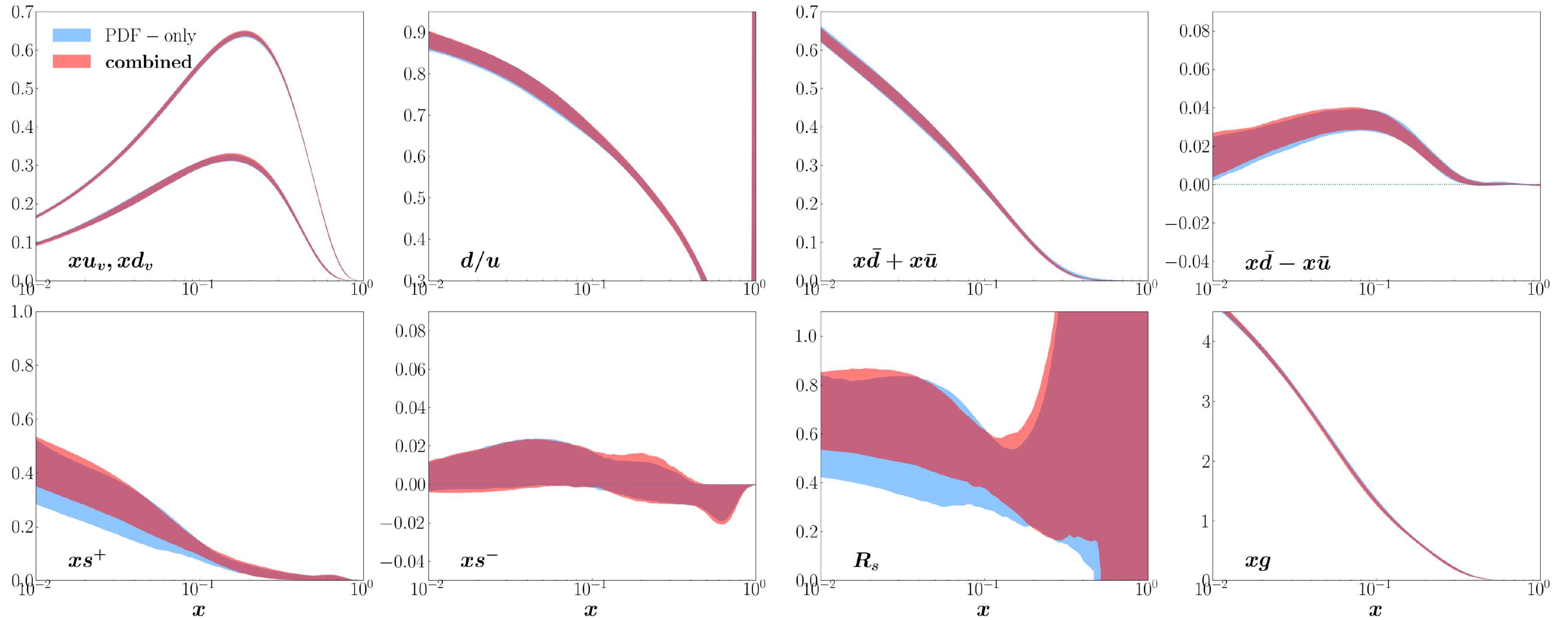
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- A combined analysis of TMDs and PDFs is possible and leads to refinements in both TMDs and PDFs
- We are planning on performing a comparative study of prescriptions for the solution of TMD evolution equations: CSS, Qiu-Zhang,  $\zeta$  prescription
- We are planning on studying collinear and TMD fragmentation functions and on inclusion of Semi-Inclusive Deep Inelastic Scattering data

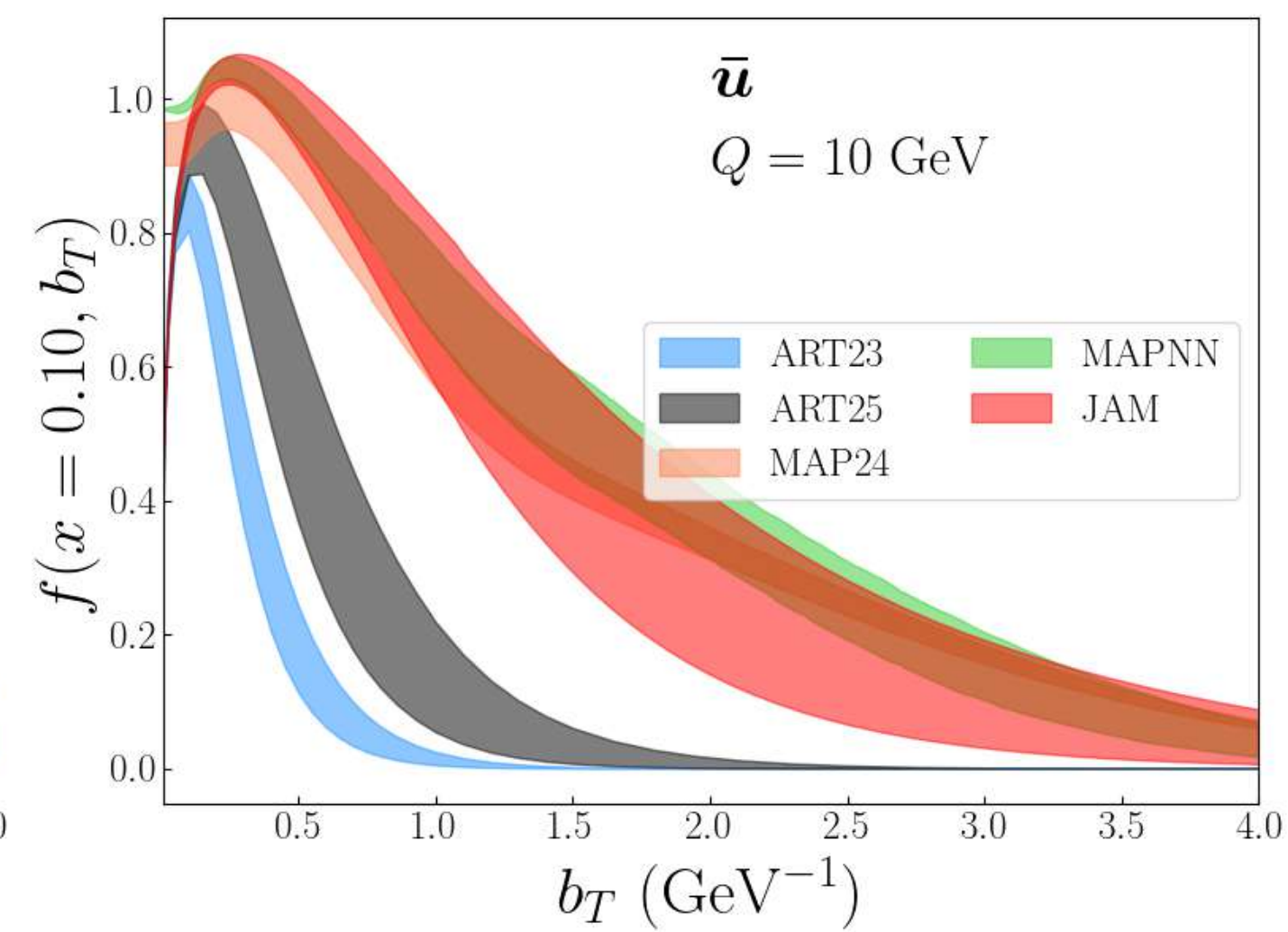
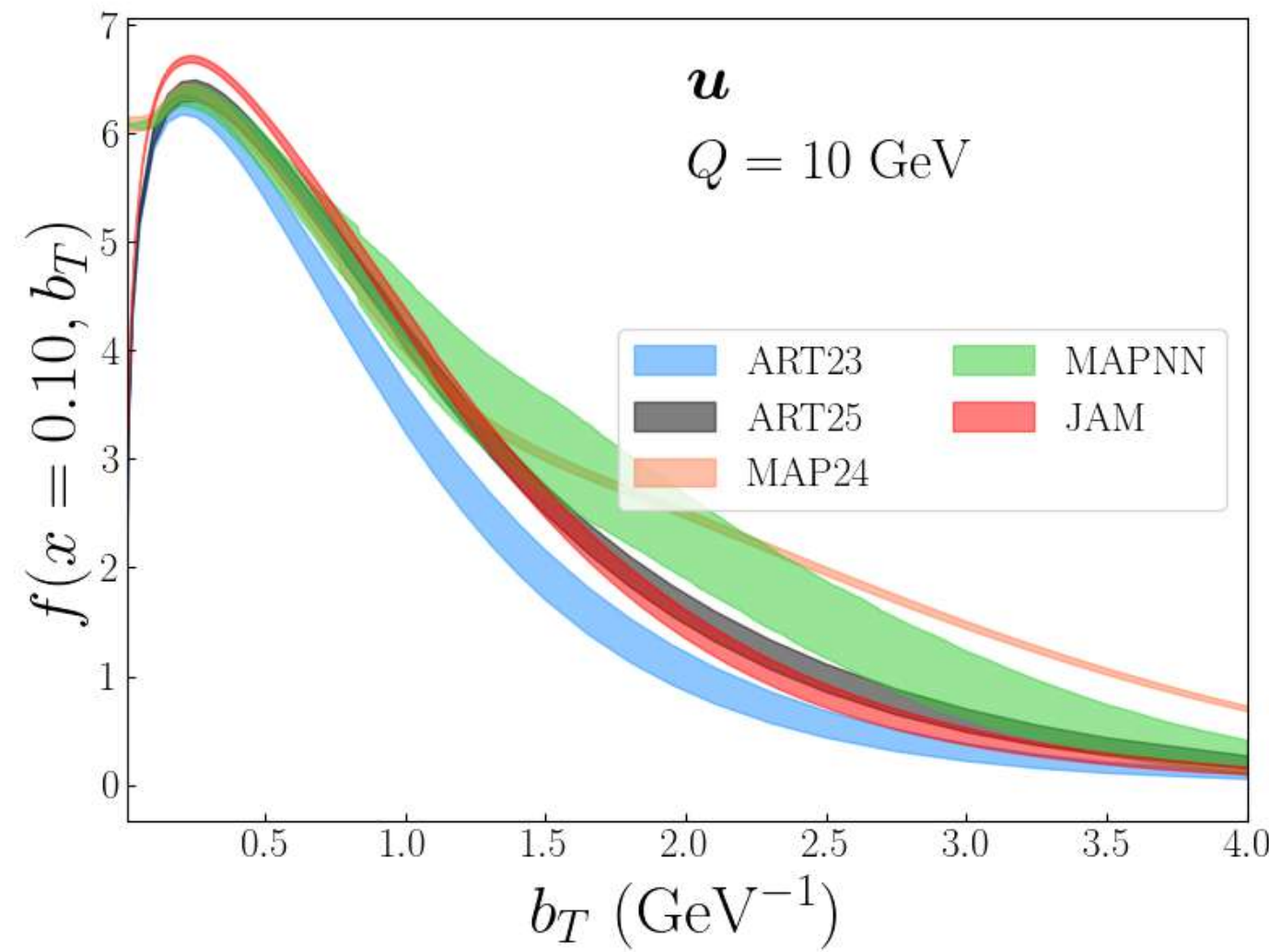
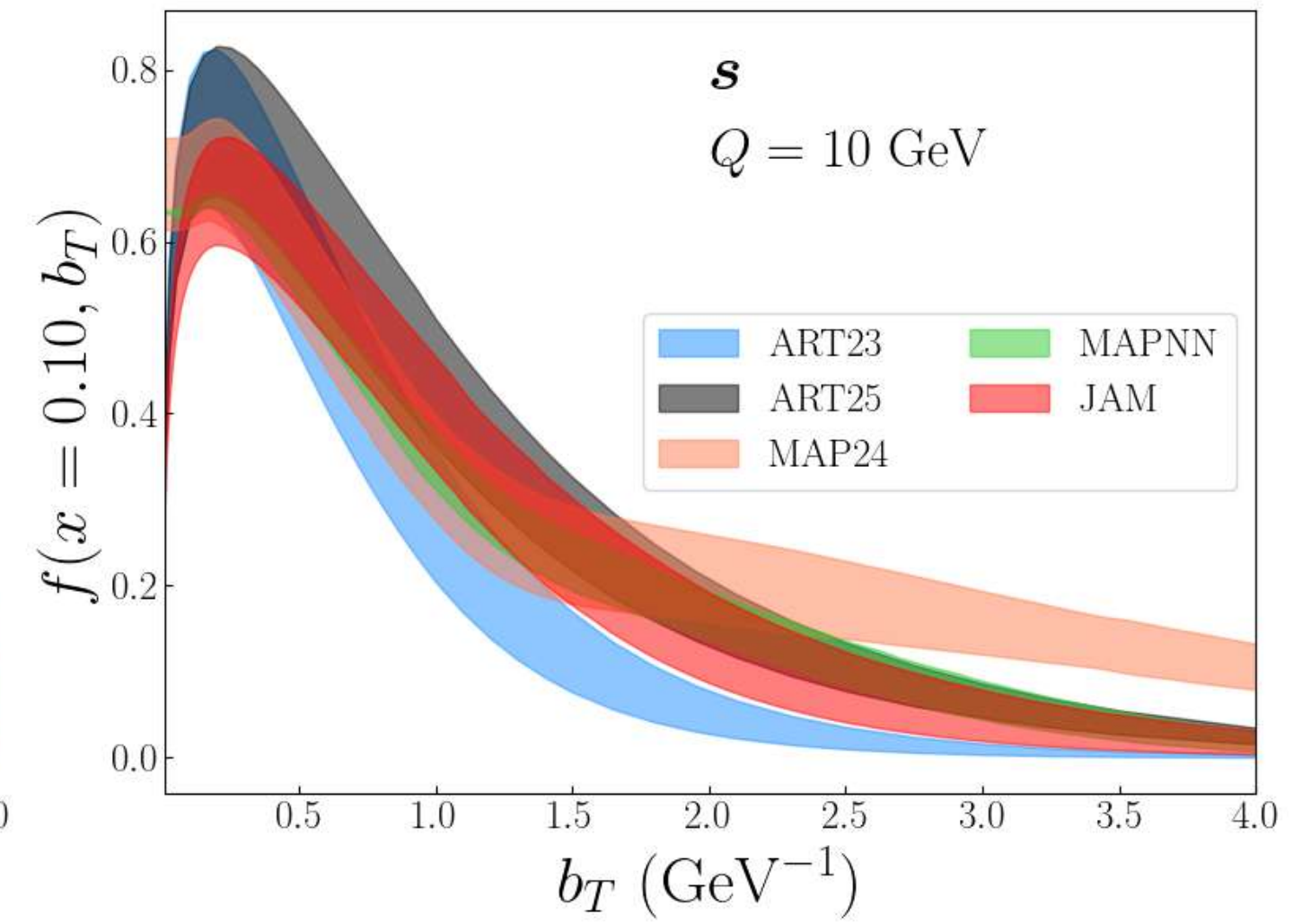
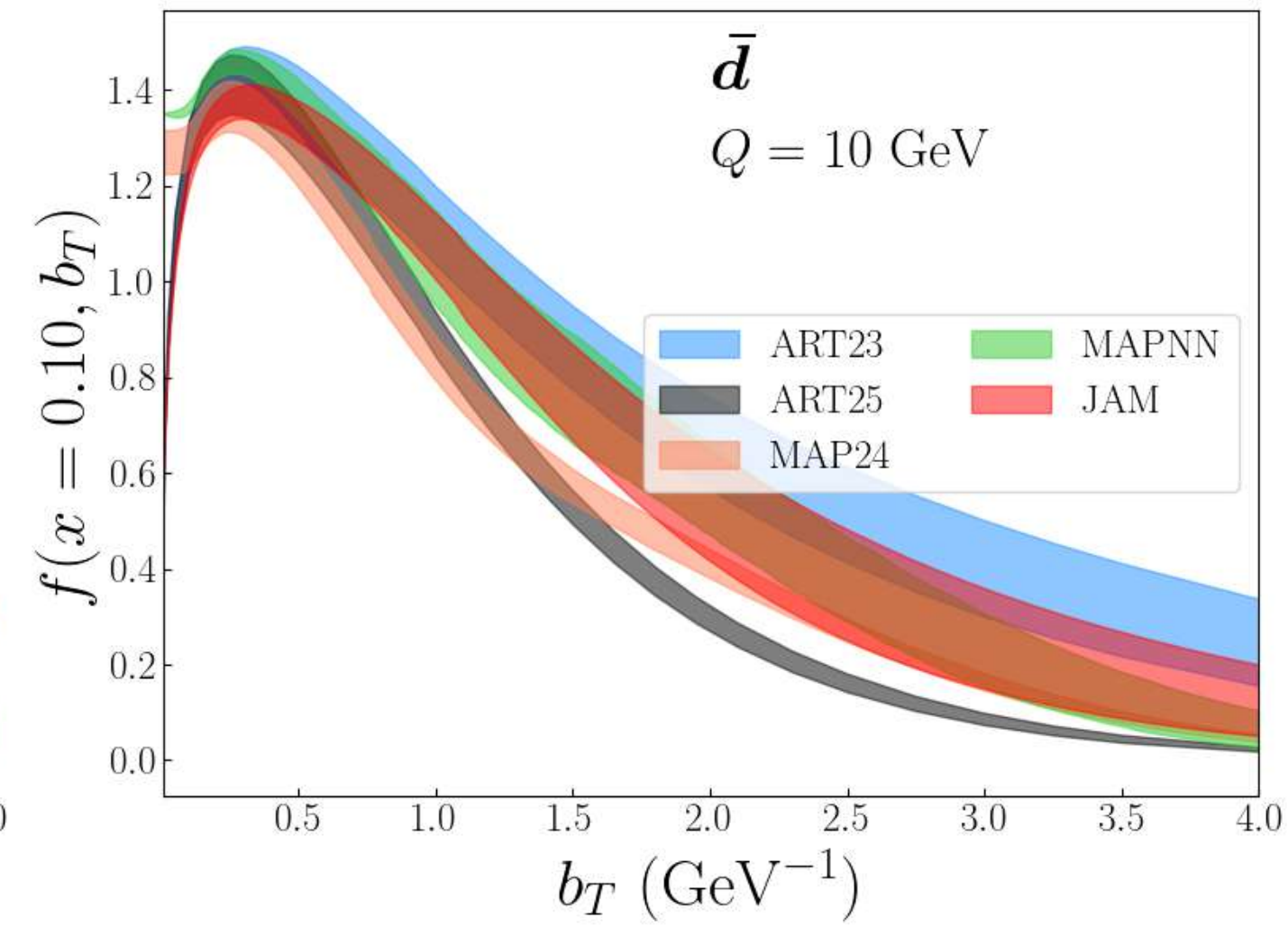
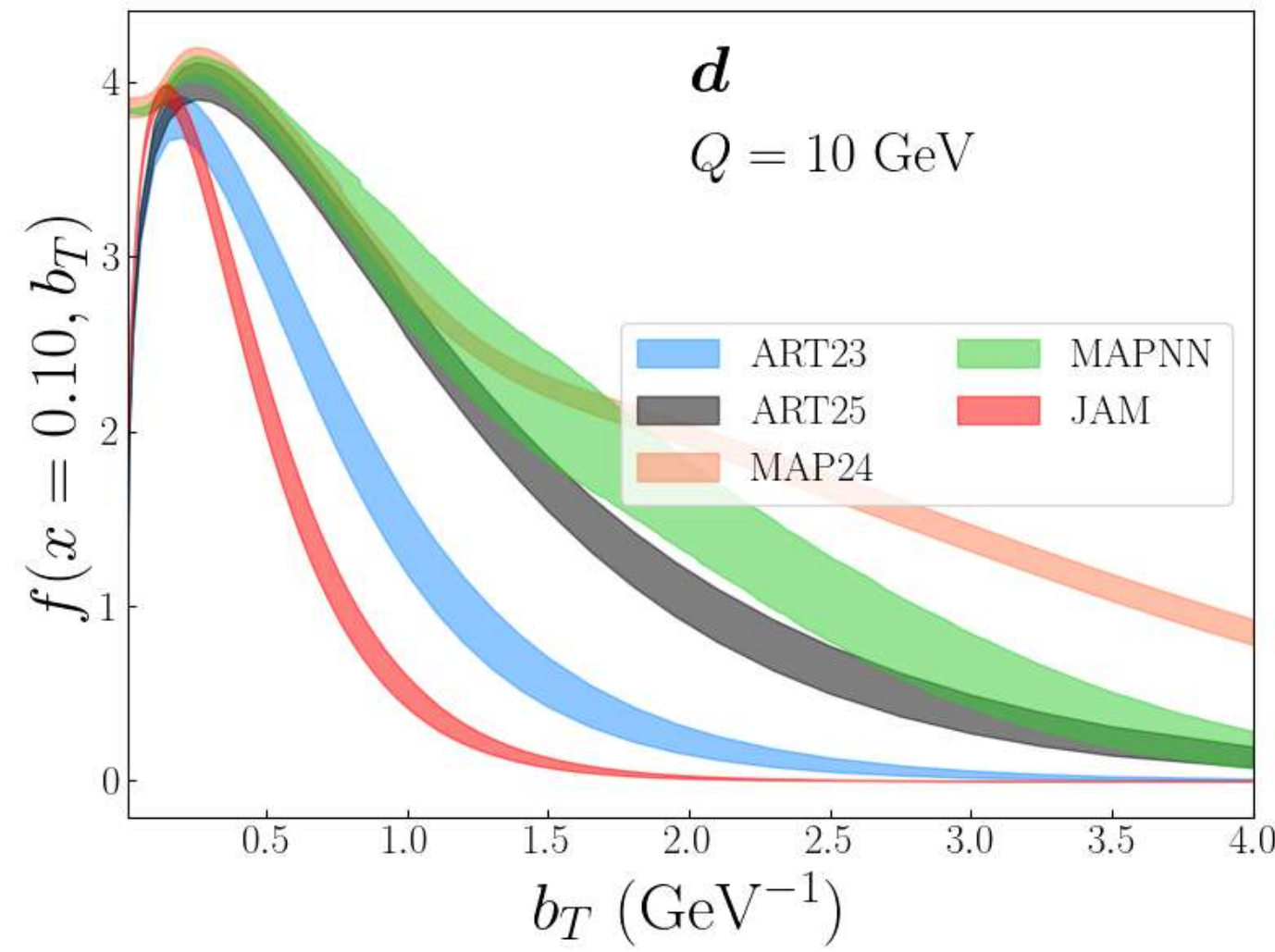


# BACKUP

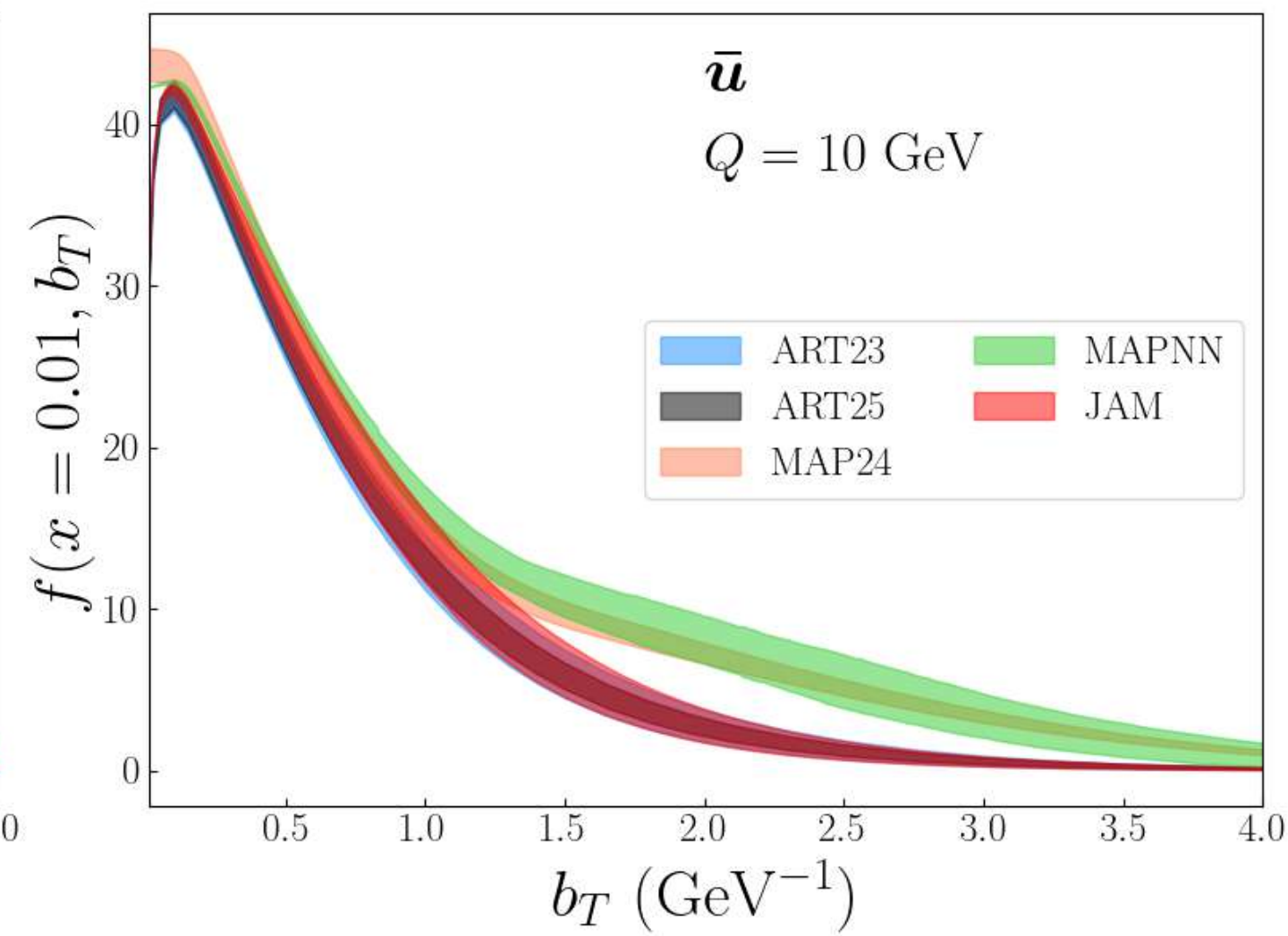
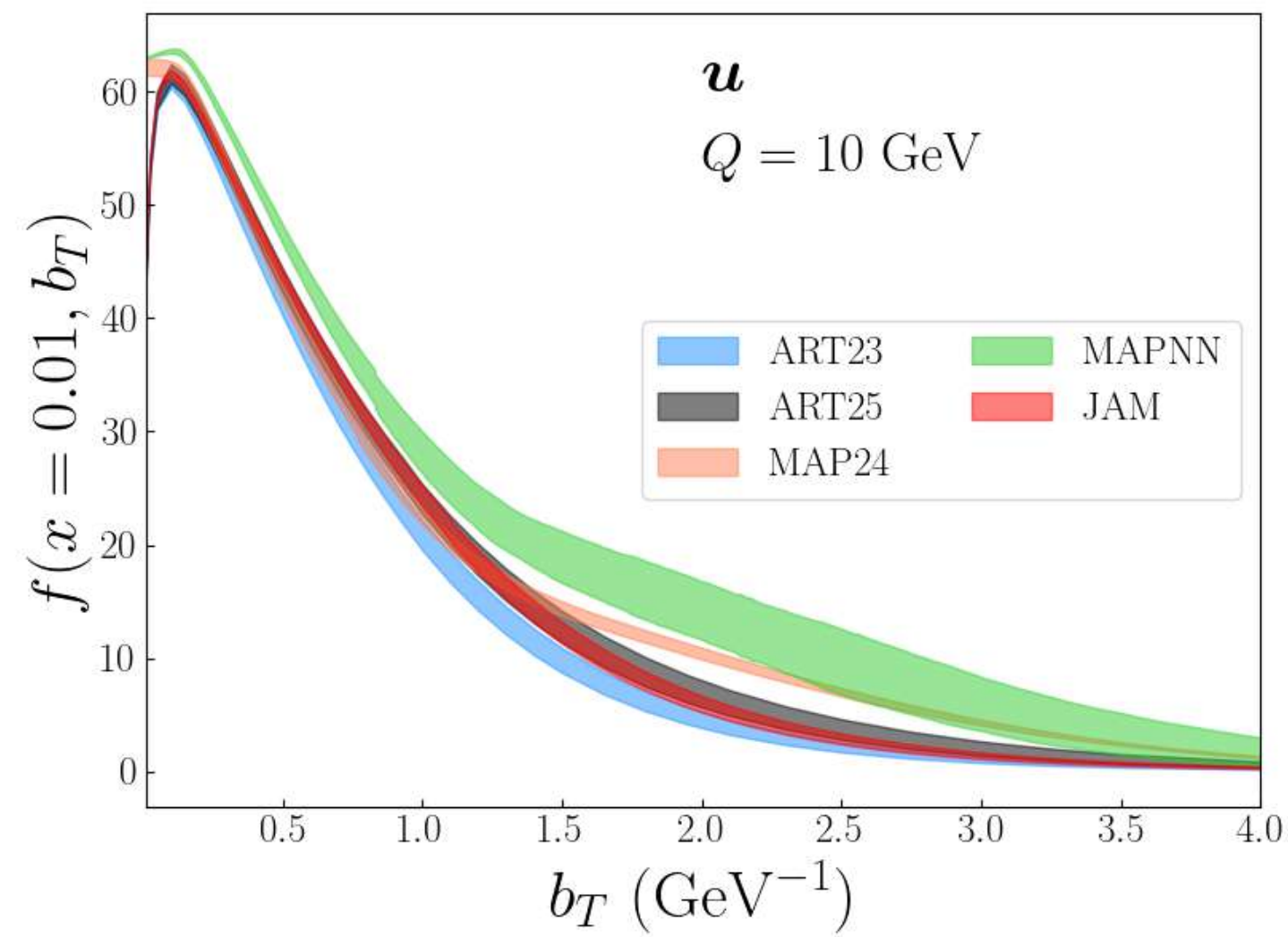
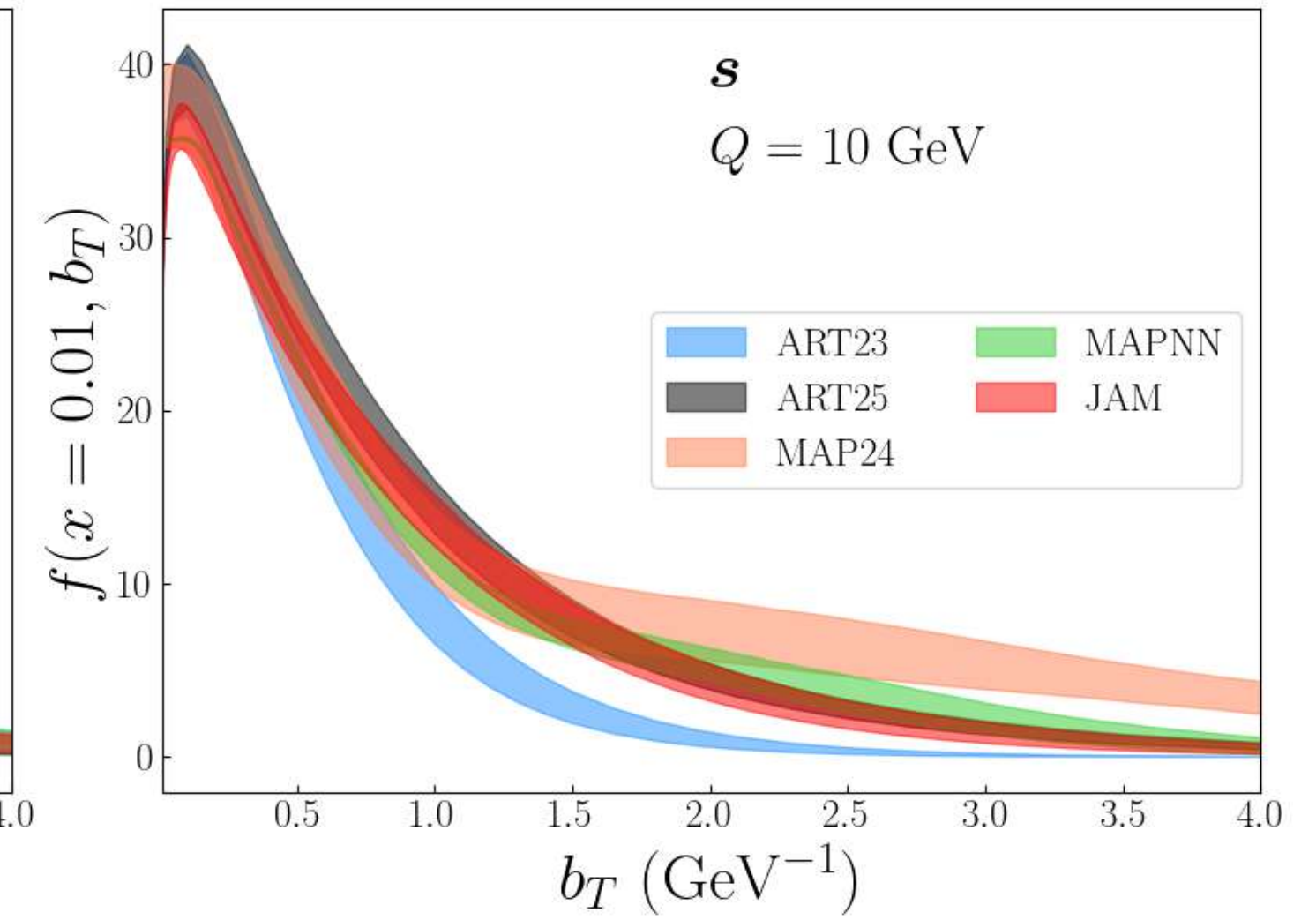
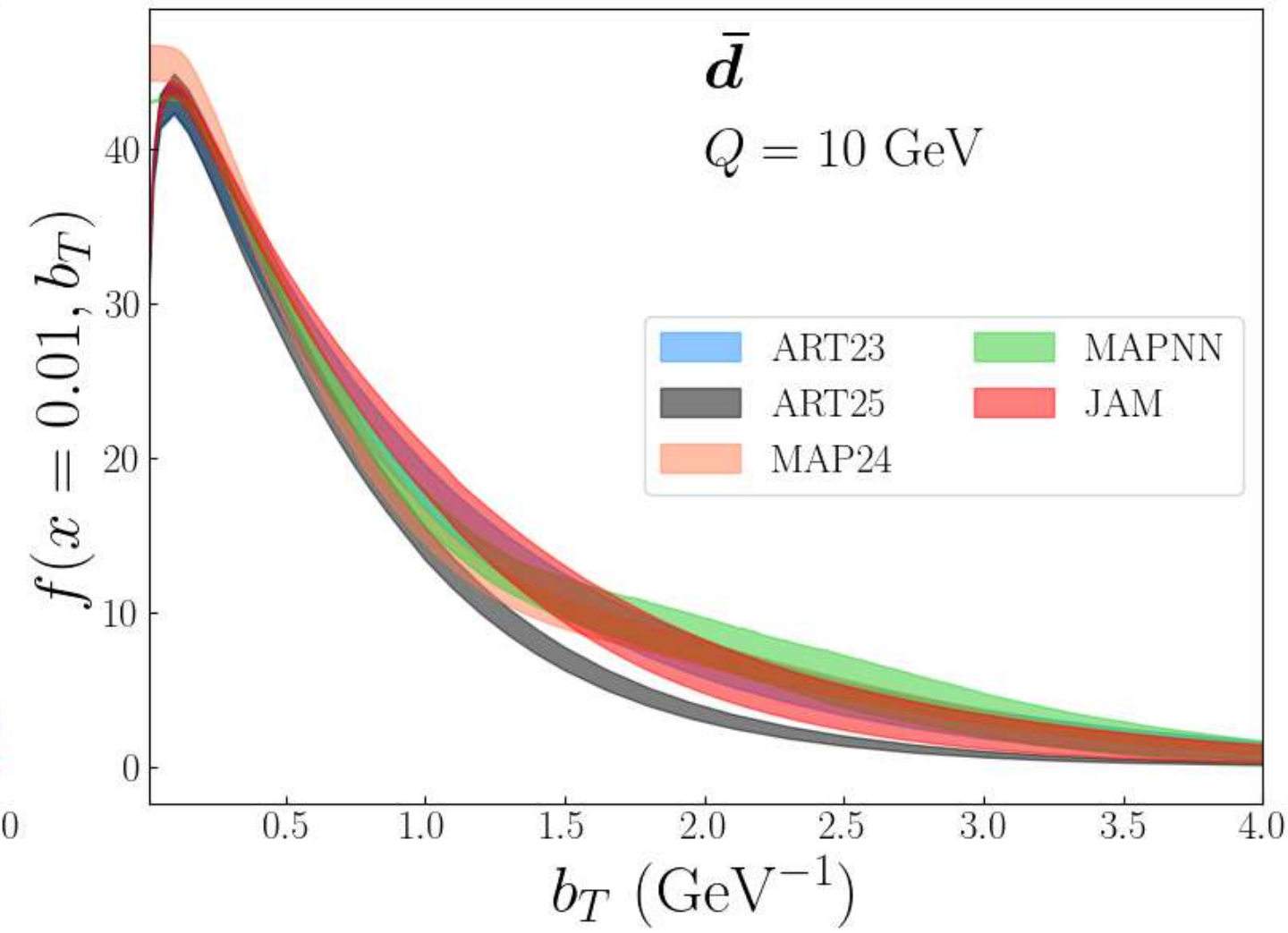
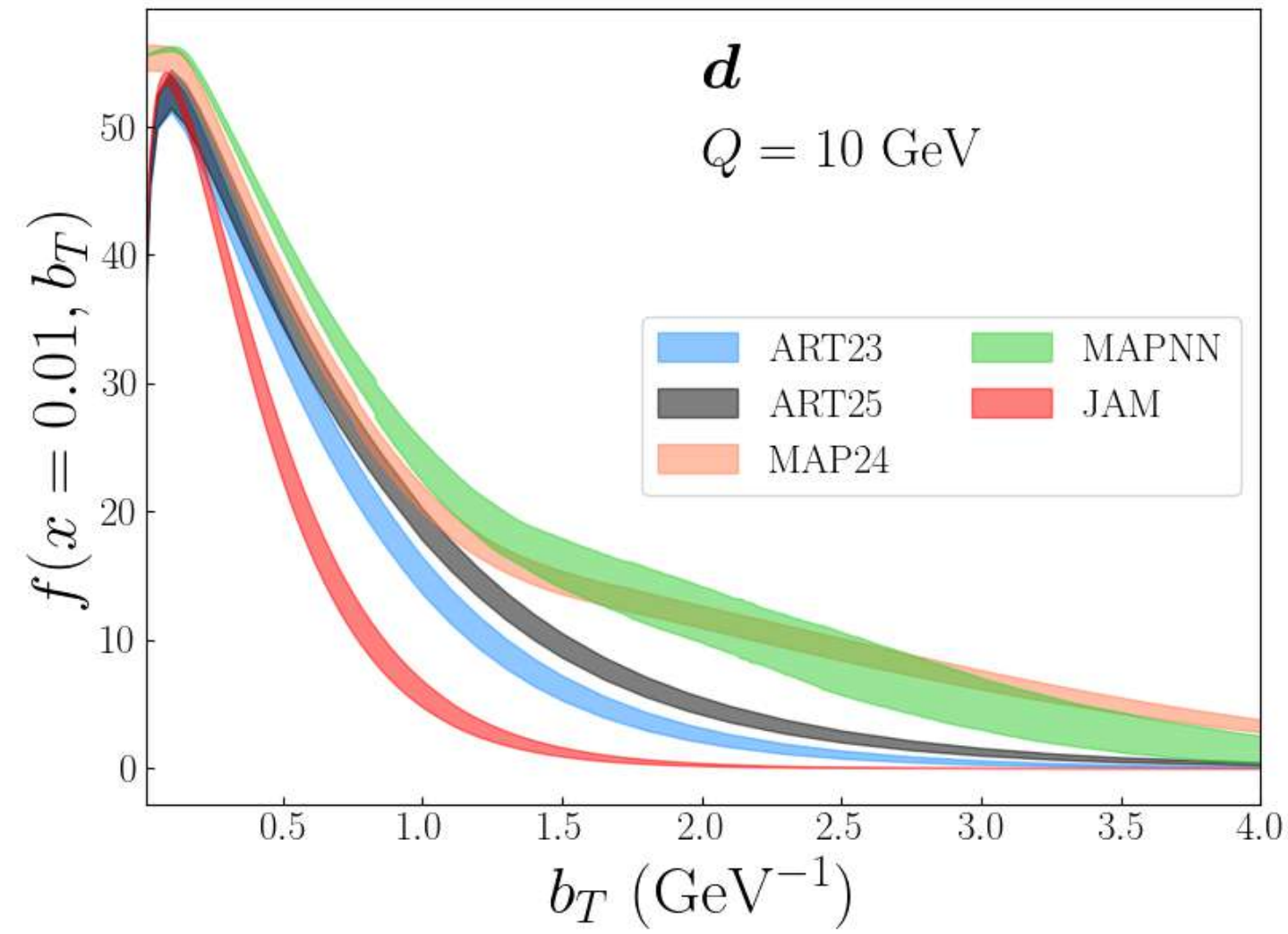
# PDFS FROM JAM25<sub>TMD+PDF</sub>



# TMDS FROM JAM25<sub>TMD+PDF</sub>



# TMDS FROM JAM25<sub>TMD+PDF</sub>



*Phenomenological framework* — The collinear data used in this analysis includes DIS data from BCDMS [64], NMC [65, 66], SLAC [67], Jefferson Lab [68–71] and HERA [72], DY lepton-pair production from the Fermilab E866 NuSea [73] and E906 SeaQuest [74] experiments,  $W^\pm$ -lepton asymmetries from CMS [75–78] and LHCb [79, 80] at the LHC and STAR [81] from RHIC, reconstructed  $Z/\gamma^*$  and  $W^\pm$  asymmetries from the Tevatron [81–85],  $W$ +charm production [86–88] from the LHC, and jet production data from the Tevatron [89, 90] and RHIC [91]. A cut on DIS data of  $W^2 > 3.5 \text{ GeV}^2$  allows sensitivity to quark distributions in the large- $x$  region.

Collins-Soper evolution equations:

$$\mu^2 \frac{d}{d\mu^2} F(x, b; \mu, \zeta) = \frac{\gamma_F(\mu, \zeta)}{2} F(x, b; \mu, \zeta)$$

$$\zeta \frac{d}{d\zeta} F(x, b; \mu, \zeta) = -\mathcal{D}(b, \mu) F(x, b; \mu, \zeta)$$

$$\mathcal{D} = -\tilde{K}/2 \quad \gamma_F(\mu, \zeta) = \Gamma_{cusp}(\mu) \ln\left(\frac{\mu^2}{\zeta}\right) - \gamma_V(\mu)$$

The solution:

$$F(x, b; \mu_f, \zeta_f) = \exp \left[ \int_P \left( \gamma_F(\mu', \zeta') \frac{d\mu'}{\mu'} - \mathcal{D}(b, \mu') \frac{d\zeta'}{\zeta'} \right) \right] F(x, b; \mu_i, \zeta_i)$$

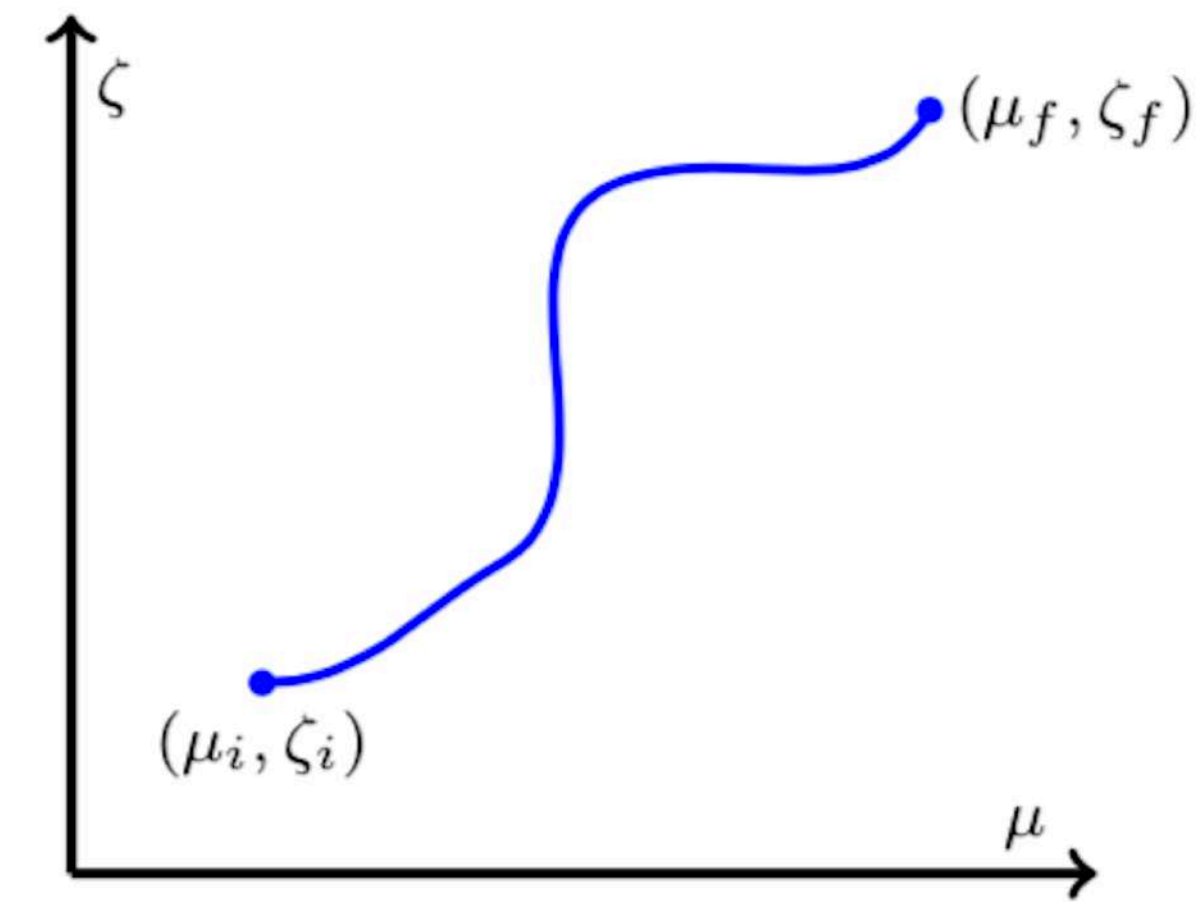
The evolution factor is path independent  
(in principle)

One can derive:

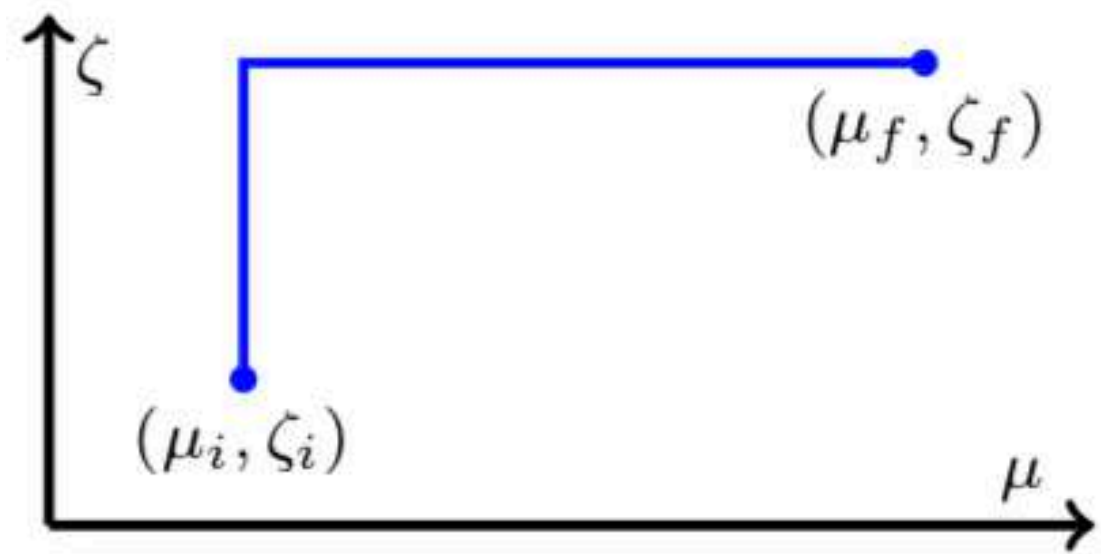
$$\zeta \frac{d}{d\zeta} \gamma_F(\mu, \zeta) = -\Gamma(\mu)$$

$$\mu \frac{d}{d\mu} \mathcal{D}(\mu, b) = \Gamma(\mu)$$

$$\zeta \frac{d}{d\zeta} \gamma_F(\mu, \zeta) = -\mu \frac{d}{d\mu} \mathcal{D}(\mu, b)$$



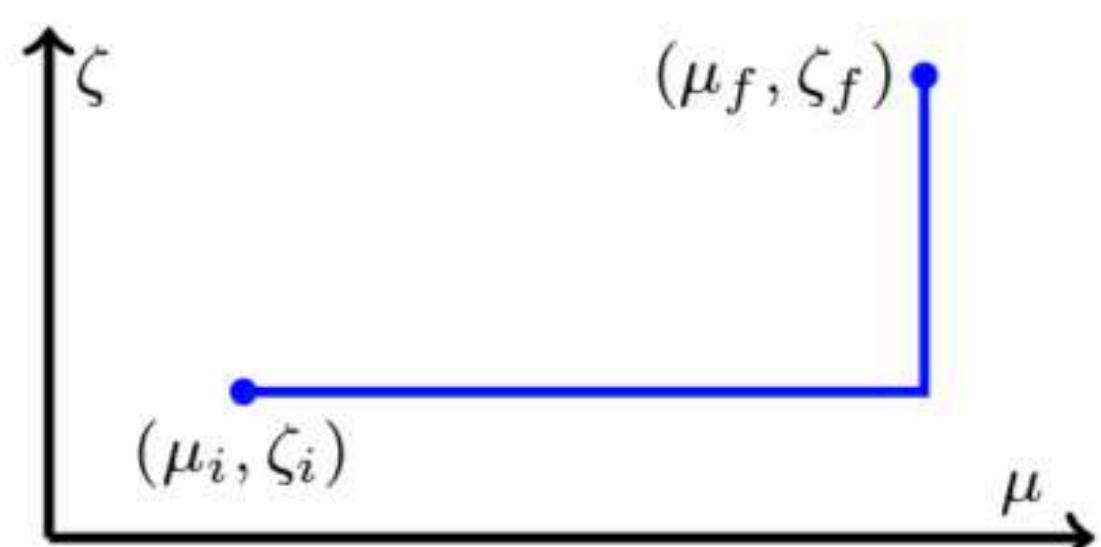
## COUPLED EVOLUTION OF TMD AND TRUNCATION OF THE PERTURBATIVE SERIES



Solution 1

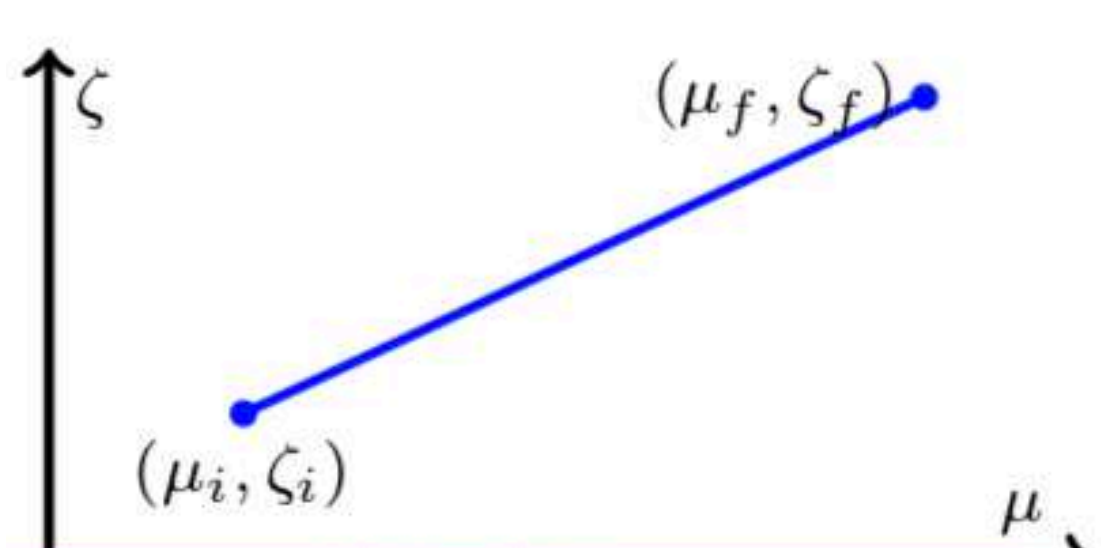
$$\ln R = \int_{\mu_i}^{\mu_f} \frac{d\mu}{\mu} \gamma_F(\mu, \zeta_f) - \mathcal{D}(\mu_i, b) \ln \left( \frac{\zeta_f}{\zeta_i} \right)$$

given in [Collins' textbook]



Solution 2

$$\ln R = \int_{\mu_i}^{\mu_f} \frac{d\mu}{\mu} \gamma_F(\mu, \zeta_i) - \mathcal{D}(\mu_f, b) \ln \left( \frac{\zeta_f}{\zeta_i} \right)$$



Solution 3

$$\ln R = \int_0^1 \left( \gamma_F(\mu(t), \zeta(t)) \frac{\mu_f - \mu_i}{(\mu_f - \mu_i)t + \mu_i} - \mathcal{D}(\mu(t), b) \frac{\zeta_f - \zeta_i}{(\zeta_f - \zeta_i)t + \zeta_i} \right) dt$$

Collins-Soper evolution equations:

$$\mu^2 \frac{d}{d\mu^2} F(x, b; \mu, \zeta) = \frac{\gamma_F(\mu, \zeta)}{2} F(x, b; \mu, \zeta)$$

$$\zeta \frac{d}{d\zeta} F(x, b; \mu, \zeta) = -\mathcal{D}(b, \mu) F(x, b; \mu, \zeta)$$

$$\mathcal{D} = -\tilde{K}/2 \quad \gamma_F(\mu, \zeta) = \Gamma_{cusp}(\mu) \ln\left(\frac{\mu^2}{\zeta}\right) - \gamma_V(\mu)$$

The solution:

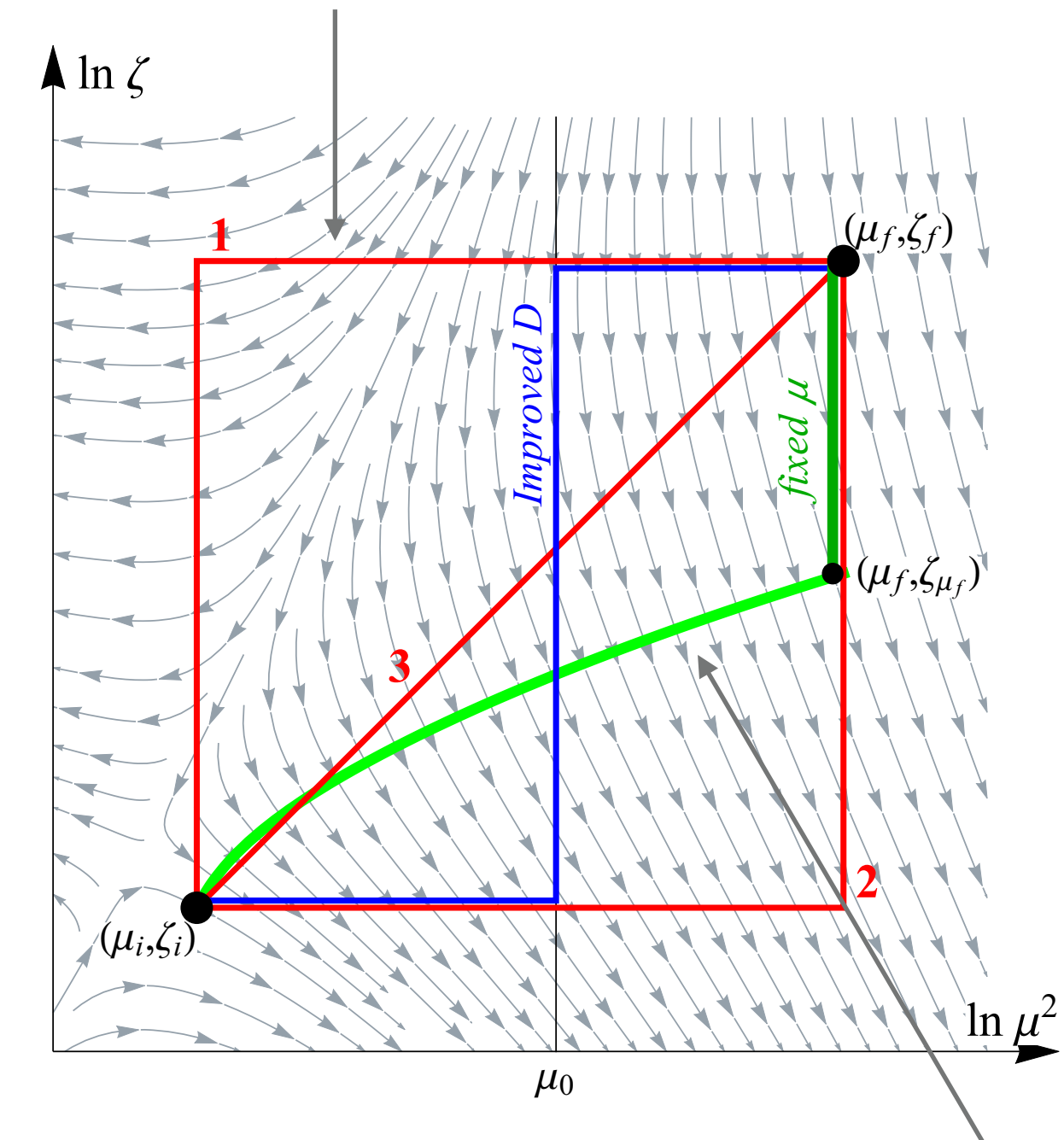
$$F(x, b; \mu_f, \zeta_f) = \exp \left[ \int_P \left( \gamma_F(\mu', \zeta') \frac{d\mu'}{\mu'} - \mathcal{D}(b, \mu') \frac{d\zeta'}{\zeta'} \right) \right] F(x, b; \mu_i, \zeta_i)$$

Equipotential lines exist in the field  $\mathbf{E} \equiv (\gamma_F(\mu, \zeta)/2, -\mathcal{D}(b, \mu))$

$$\Gamma_{cusp}(\mu) \ln\left(\frac{\mu^2}{\zeta_\mu(b)}\right) - \gamma_V(\mu) = 2\mathcal{D}(b, \mu) \frac{d \ln \zeta_\mu(b)}{d \ln \mu^2}$$

TMDs do not evolve along those lines  $\mu^2 \frac{d}{d\mu^2} F(x, b; \mu, \zeta_\mu(b)) = 0$

Standard TMD formalism



$\zeta$ -prescription

Collins-Soper evolution equations:

$$\mu^2 \frac{d}{d\mu^2} \tilde{F}(x, b; \mu, \zeta) = \frac{\gamma_F(\mu, \zeta)}{2} \tilde{F}(x, b; \mu, \zeta)$$

$$\zeta \frac{\partial}{\partial \zeta} \tilde{F}(x, b; \mu, \zeta) = -\mathcal{D}(b, \mu) \tilde{F}(x, b; \mu, \zeta)$$

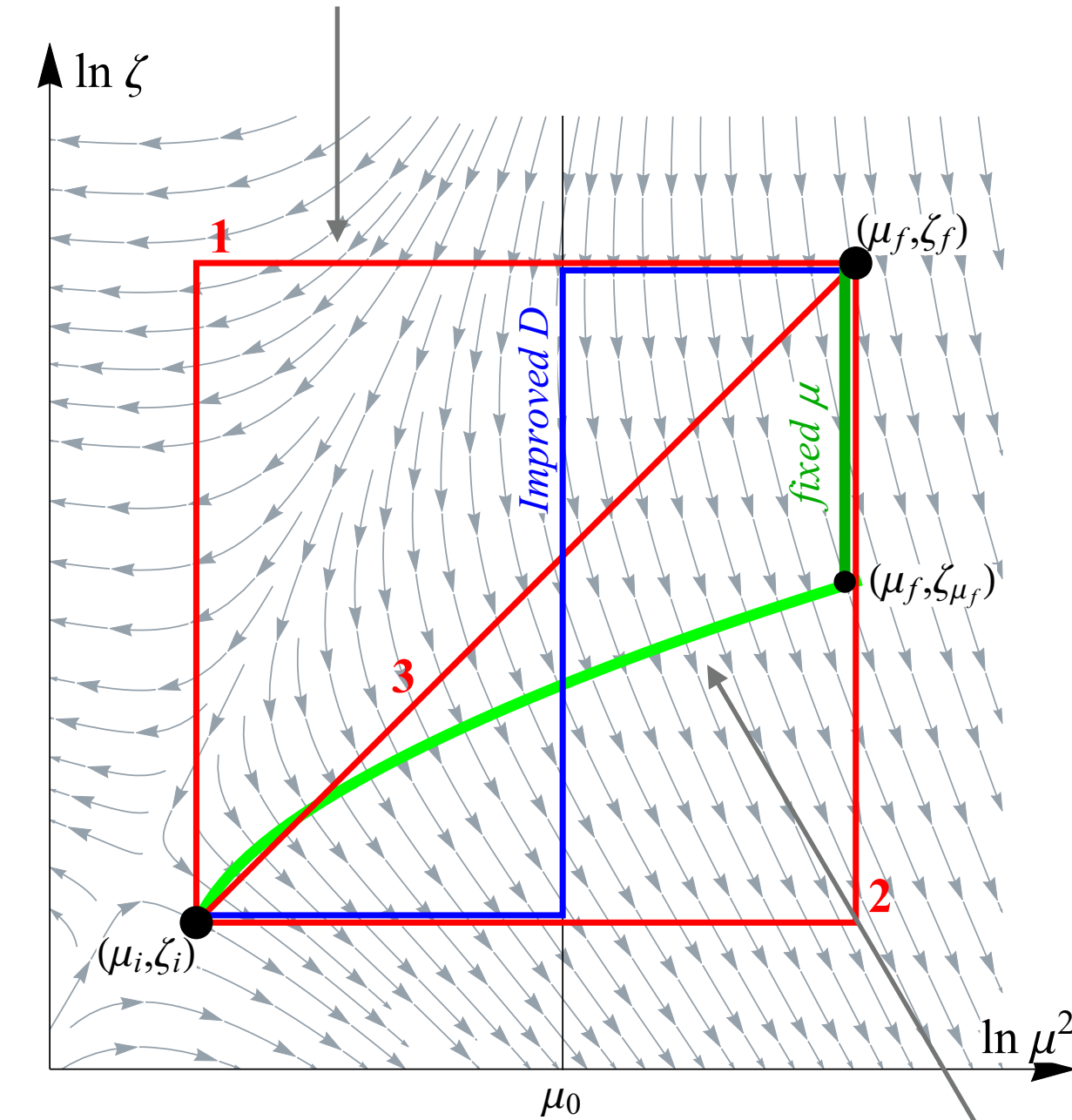
$$\mathcal{D} = -\tilde{K}/2$$

The solution:

$$F(x, b; \mu, \zeta) = \exp \left[ \int_P \left( \gamma_F(\mu', \zeta') \frac{d\mu'}{\mu'} - \mathcal{D}(b, \mu') \frac{d\zeta'}{\zeta'} \right) \right] F(x, b; \mu_0, \zeta_0)$$

Equipotential lines exist in the field  $\mathbf{E} \equiv (\gamma_F(\mu, \zeta)/2, -\mathcal{D}(b, \mu))$   $\zeta$ -prescription

Standard TMD formalism



TMDs do not evolve along those lines

$$\mu^2 \frac{d}{d\mu^2} F(x, b; \mu, \zeta_\mu(b)) = 0$$

One particular line is of interest, the saddle point  $(\mu_0, \zeta_0)$ :

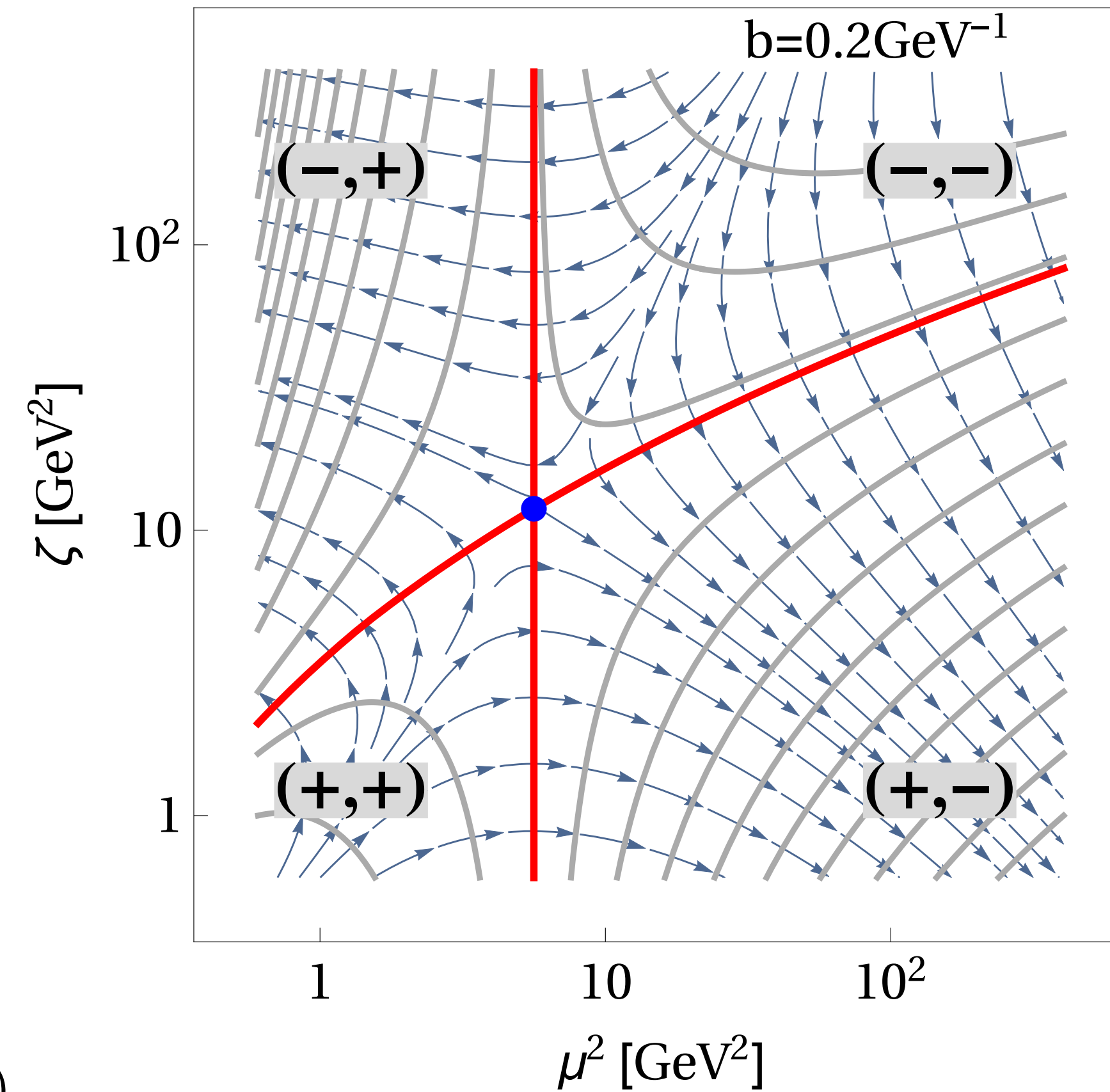
$$\mathcal{D}(b, \mu_0) = 0,$$

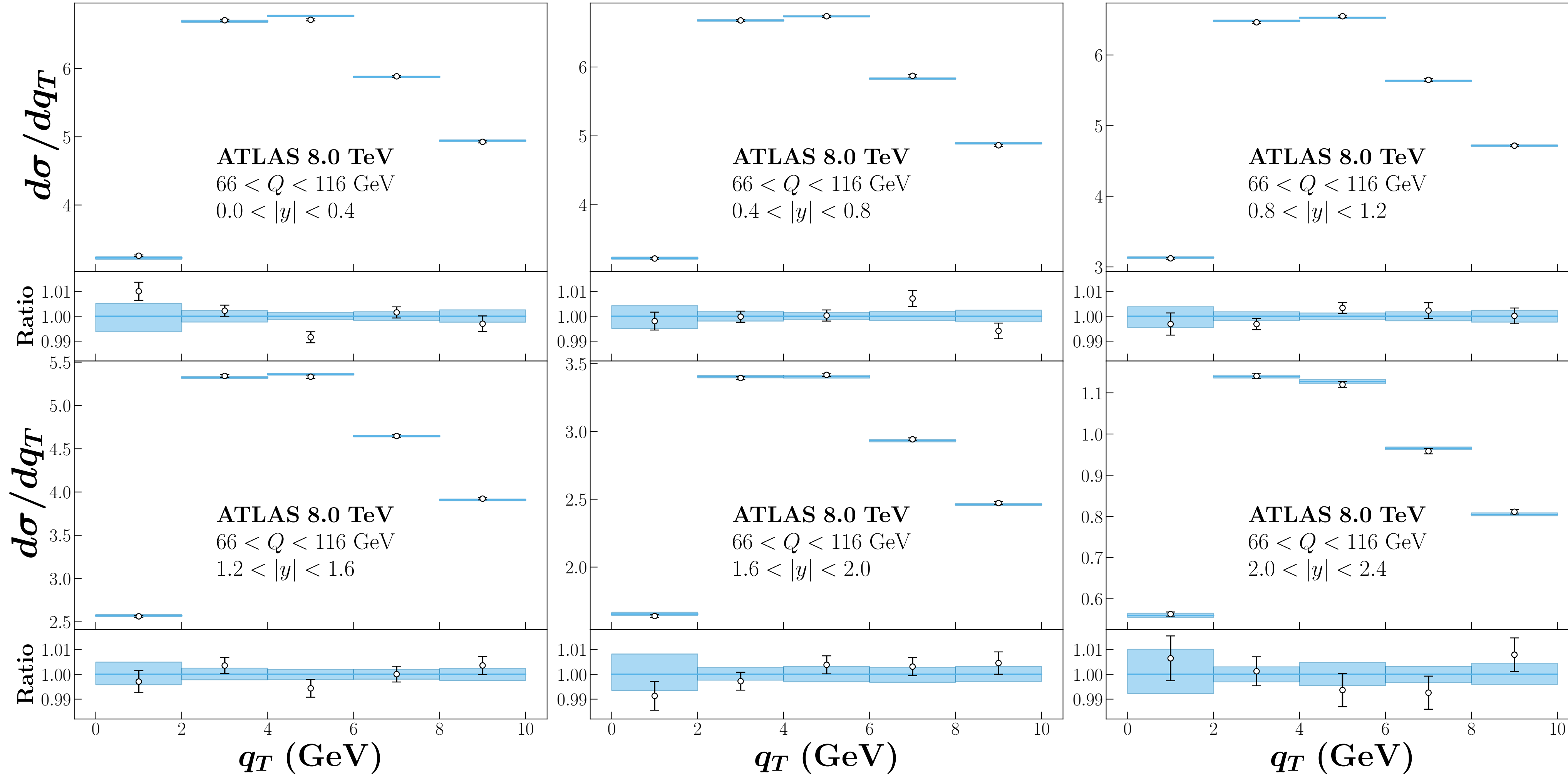
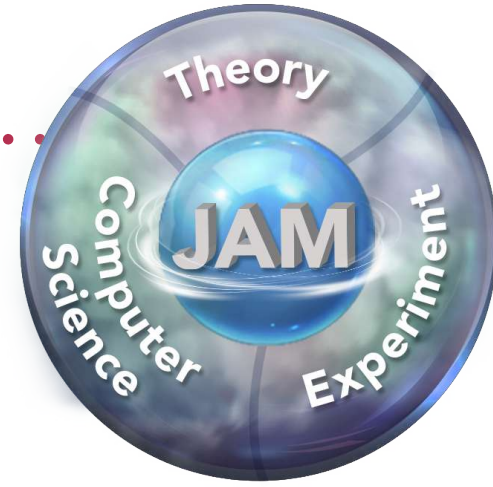
$$\Gamma_{\text{cusp}}(\mu_0) \ln \left( \frac{\mu_0^2}{\zeta_0} \right) - \gamma_V(\mu_0) = 0.$$

TMD on this line is called the optimal TMD and it has no scale dependence  $F(x, b)$

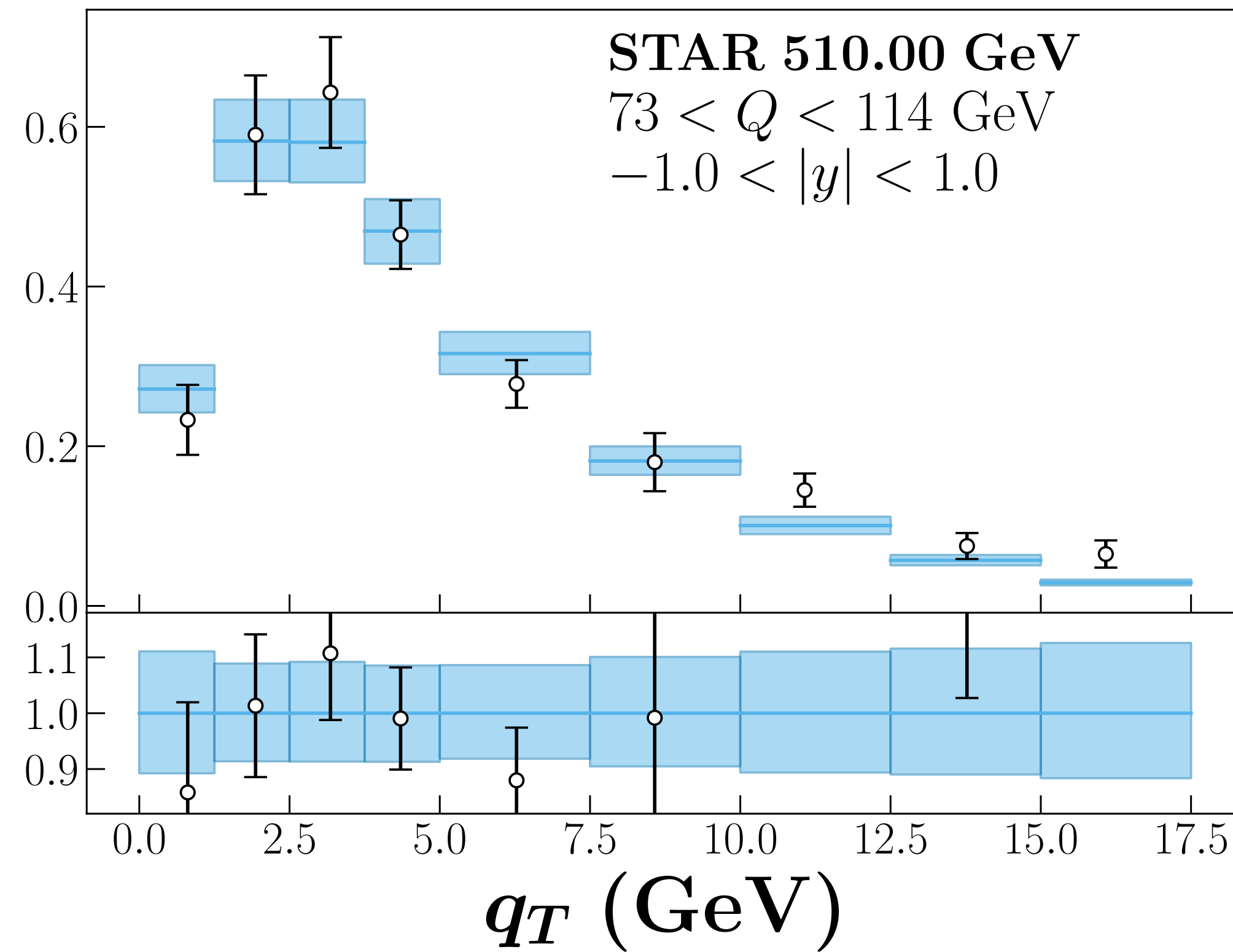
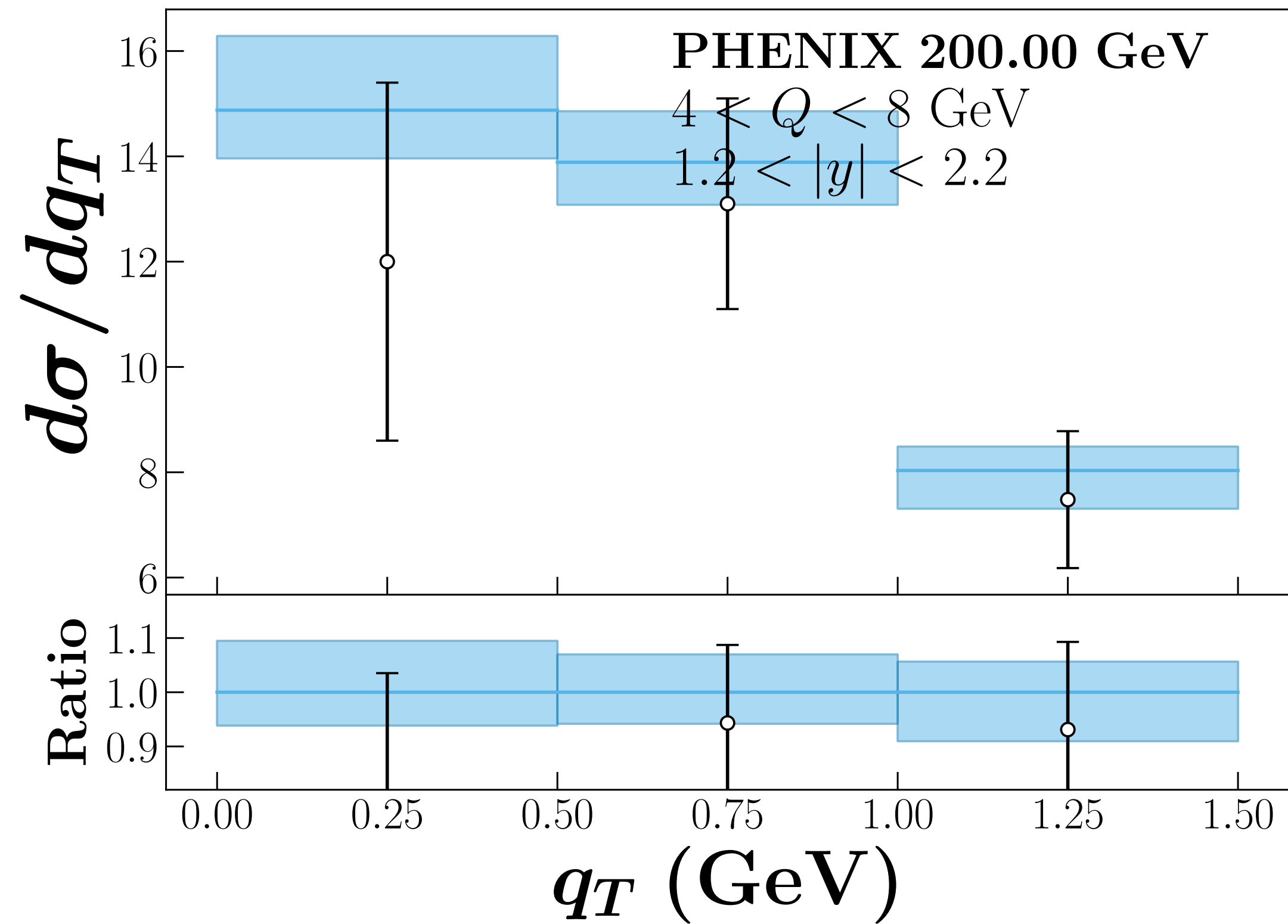
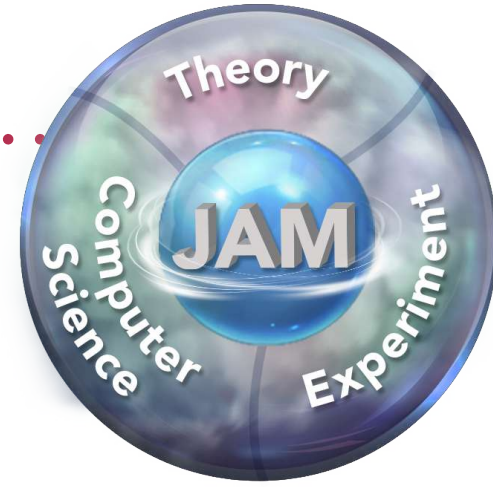
The evolution becomes a multiplicative factor

$$F(x, b; \mu, \zeta) = \left( \frac{\zeta}{\zeta_\mu(b)} \right)^{-\mathcal{D}(b, \mu)} F(x, b)$$



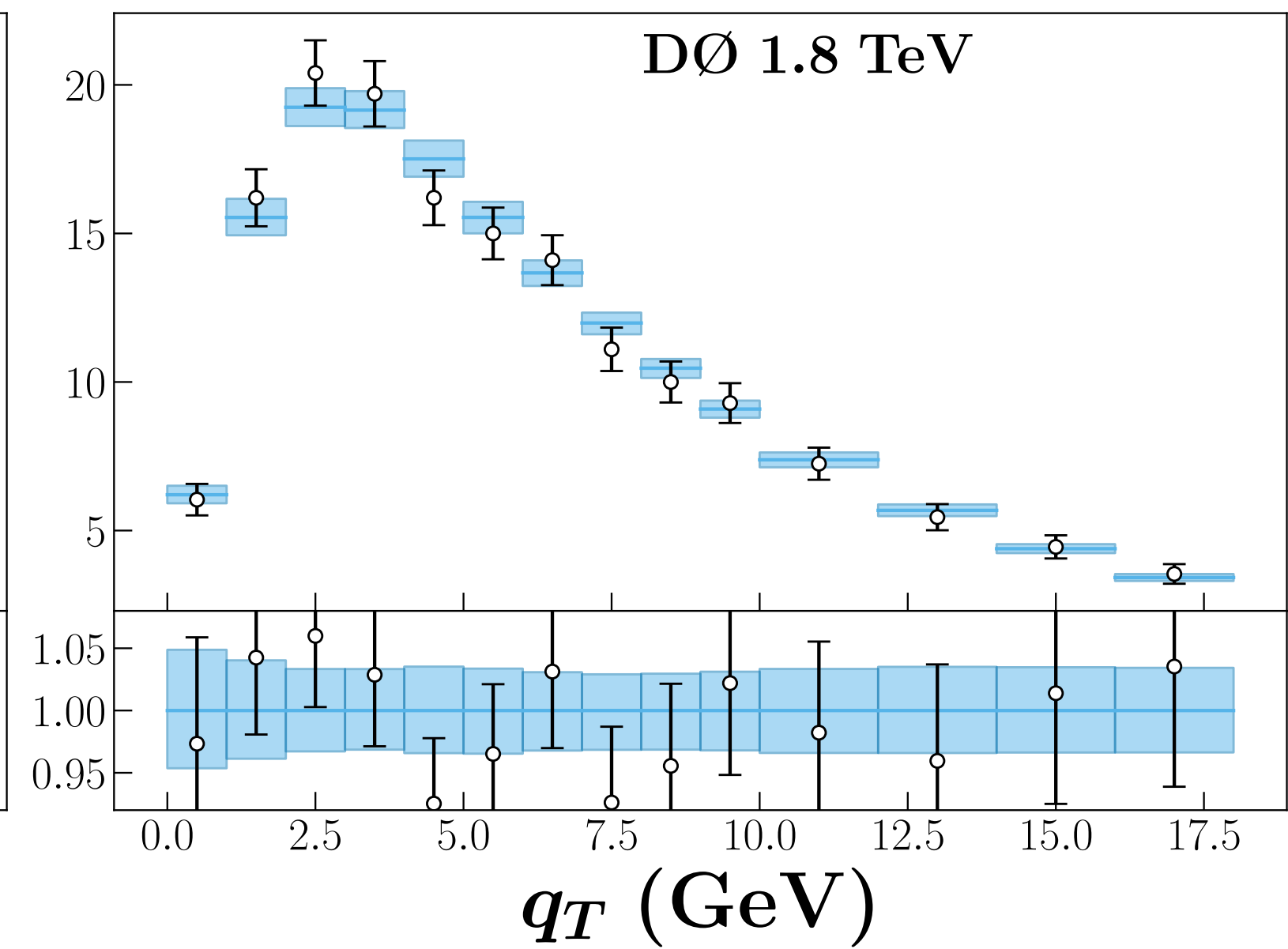
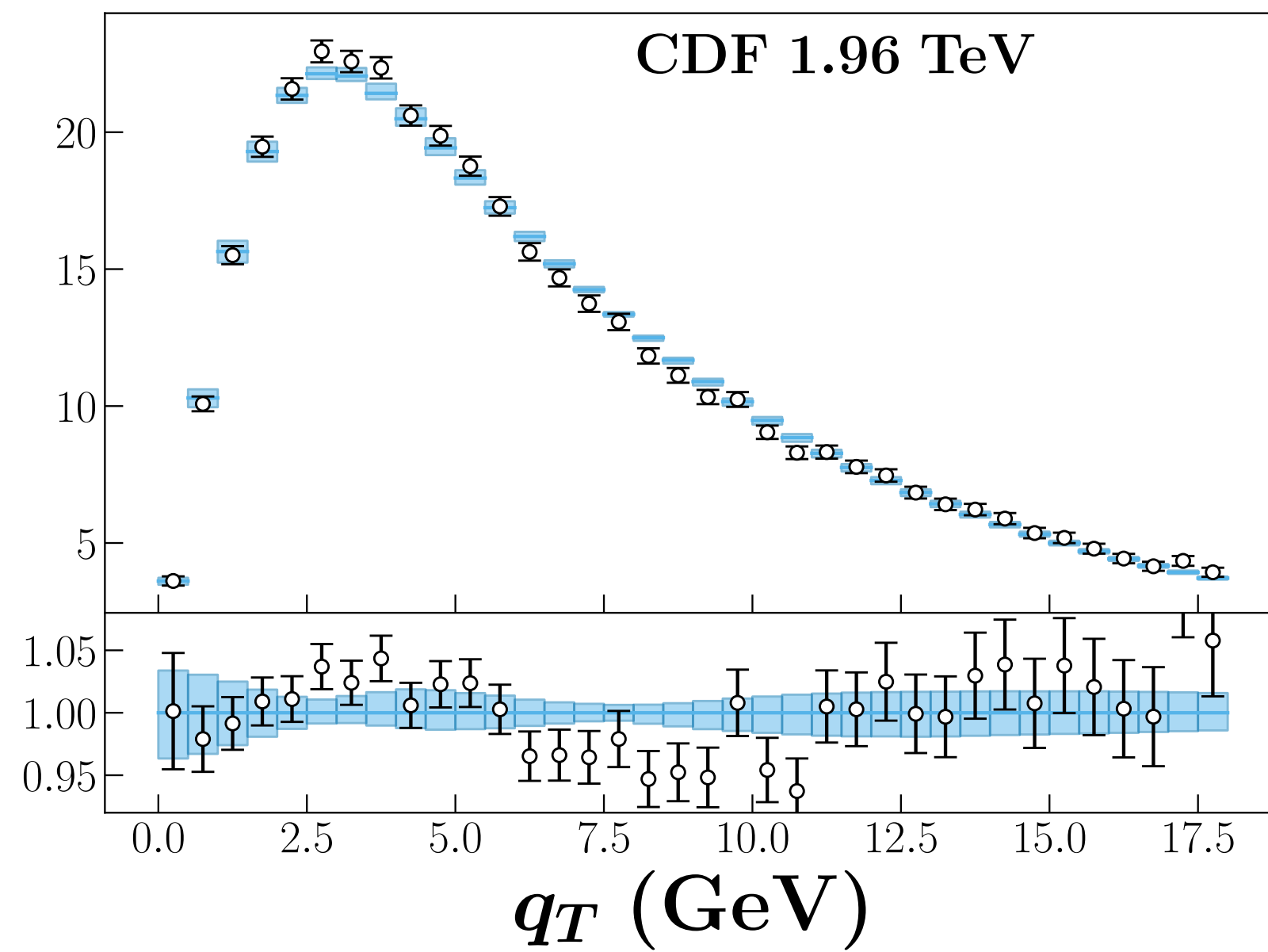
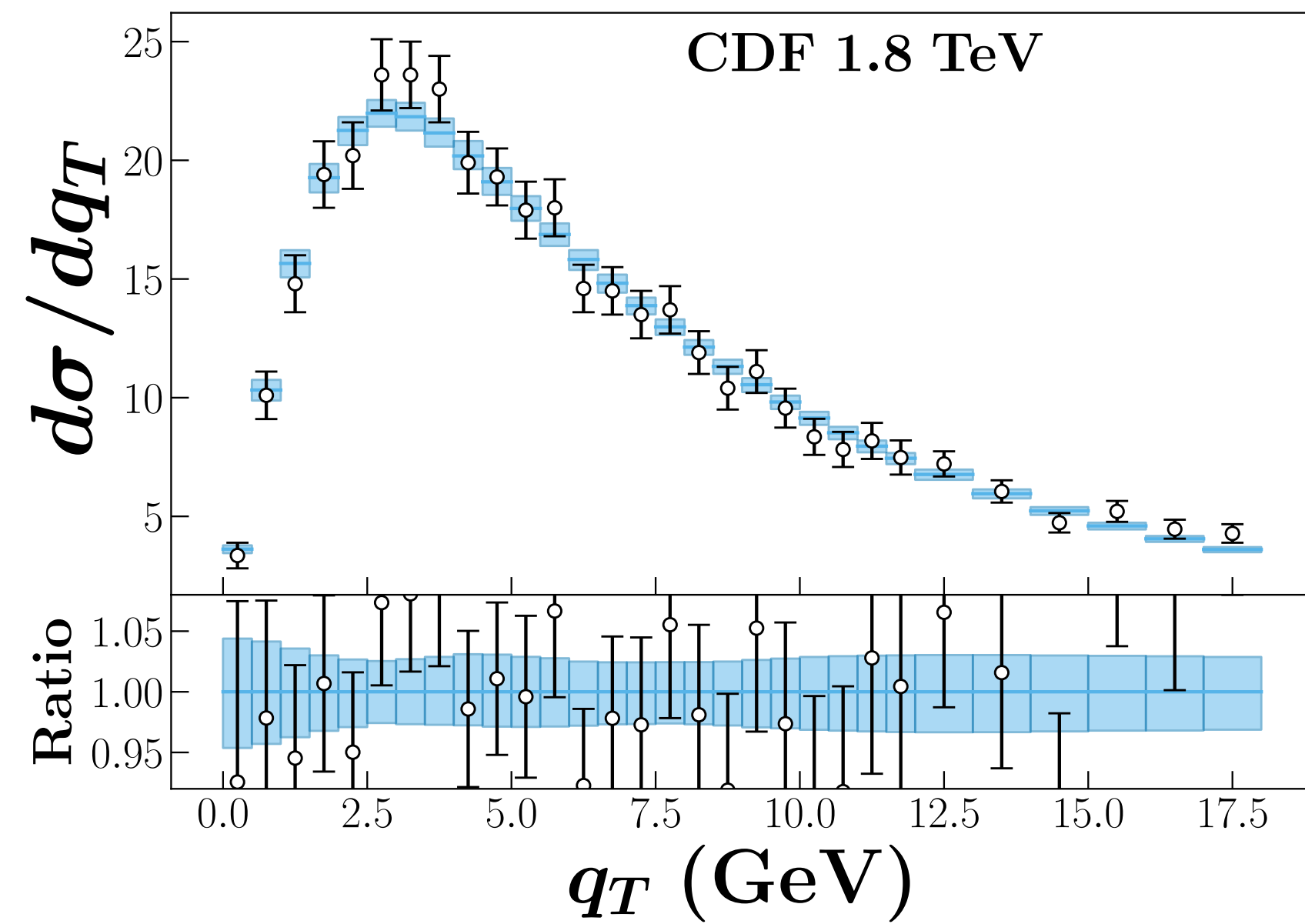
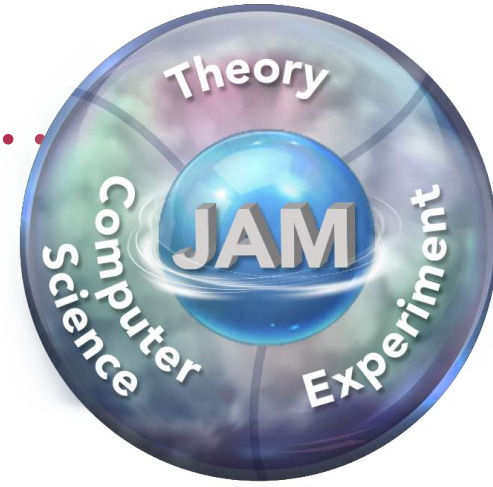


ATLAS data are in bins of rapidity, very sensitive to the shape of PDFs

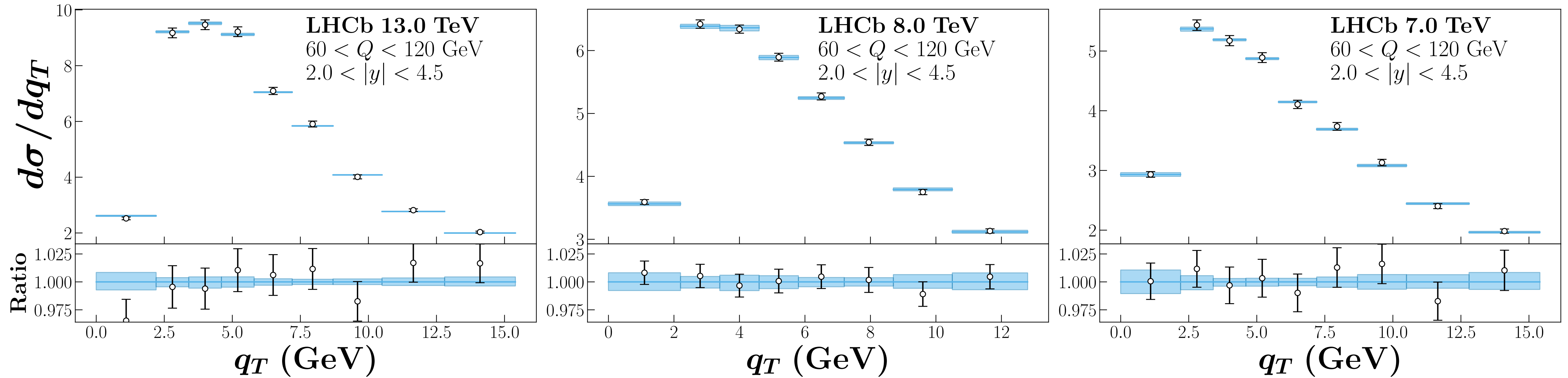
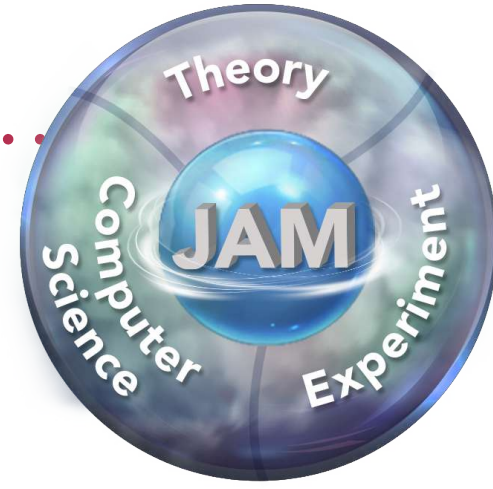


RHIC data are integrated in rapidity

# RESULTS

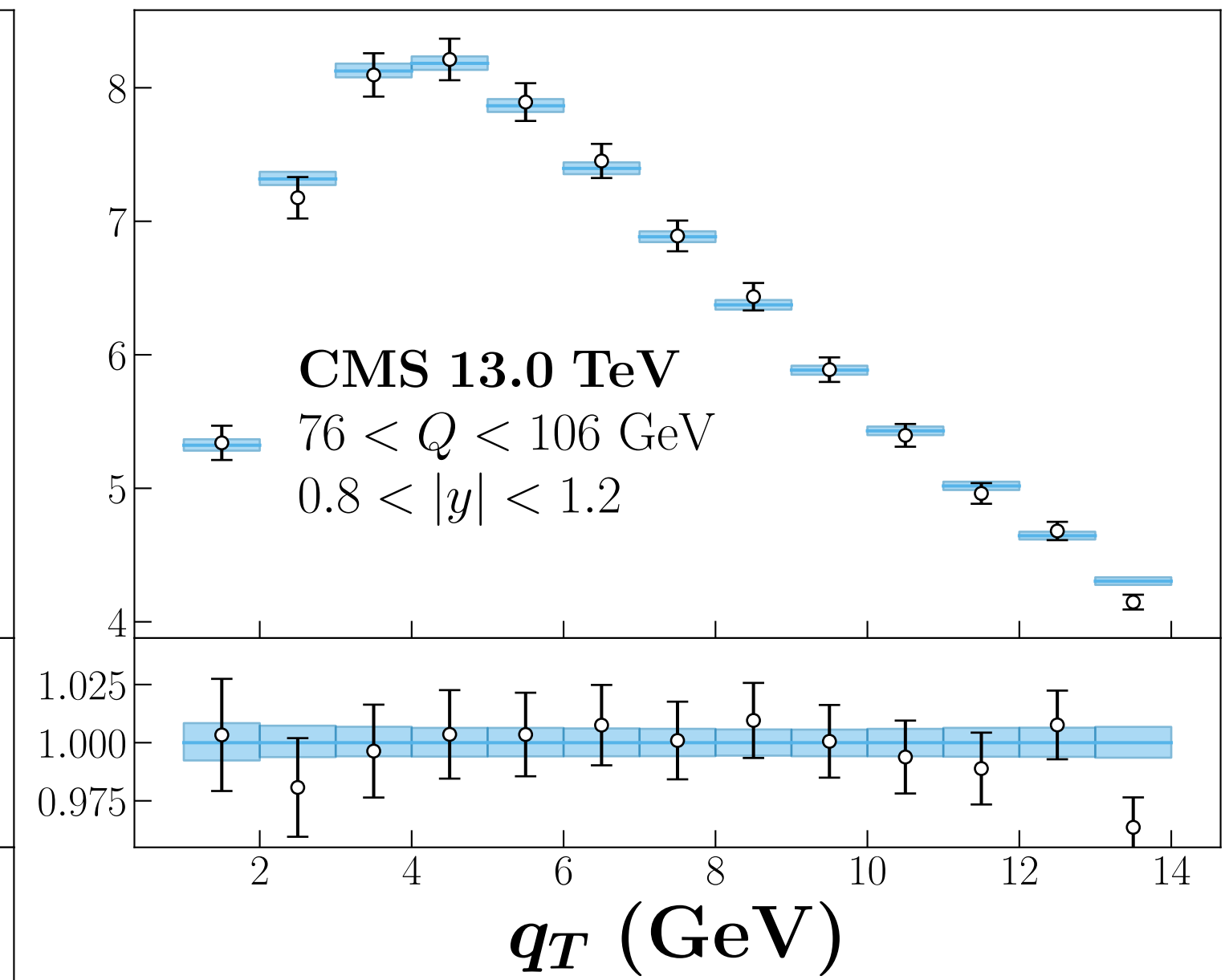
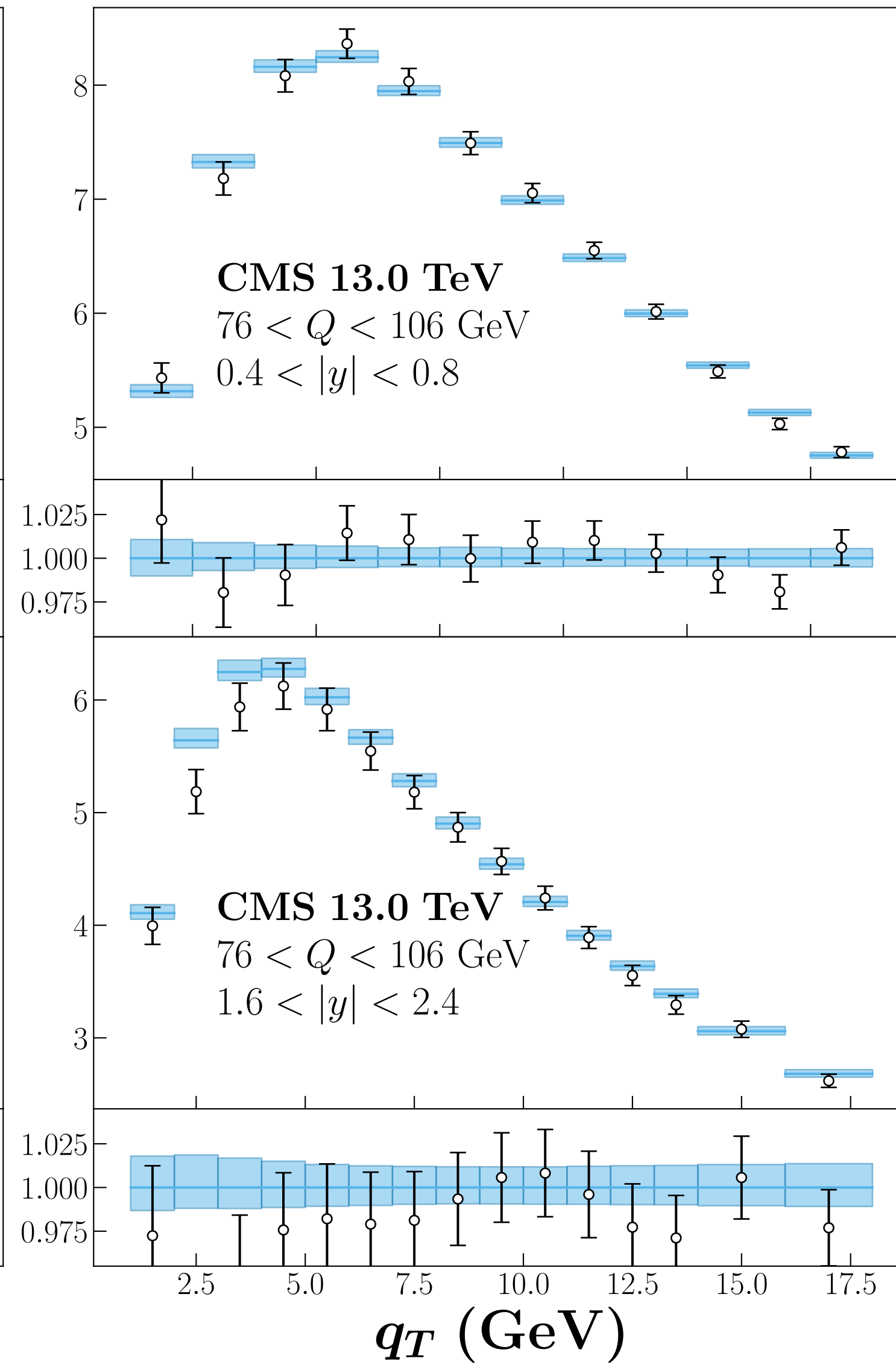
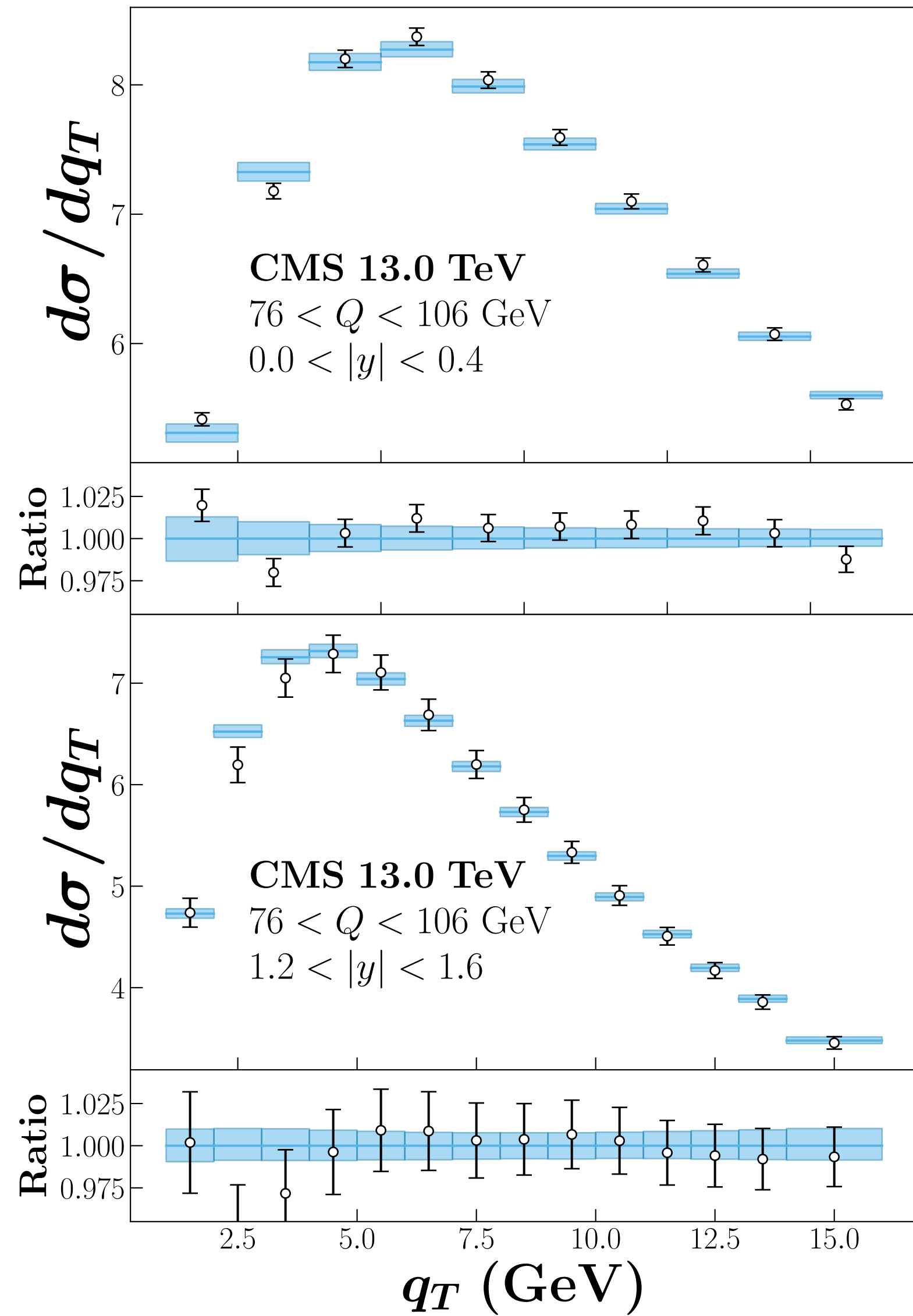
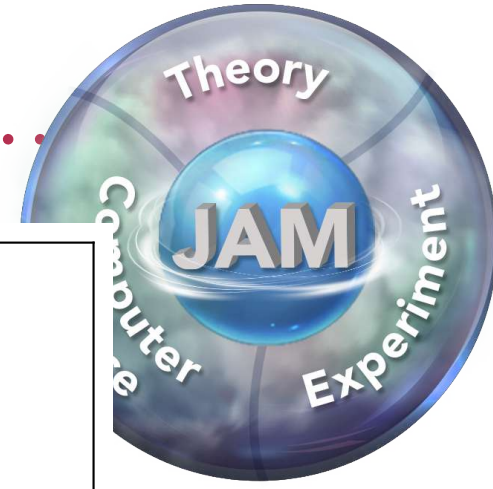


TEVATRON data are integrated in rapidity



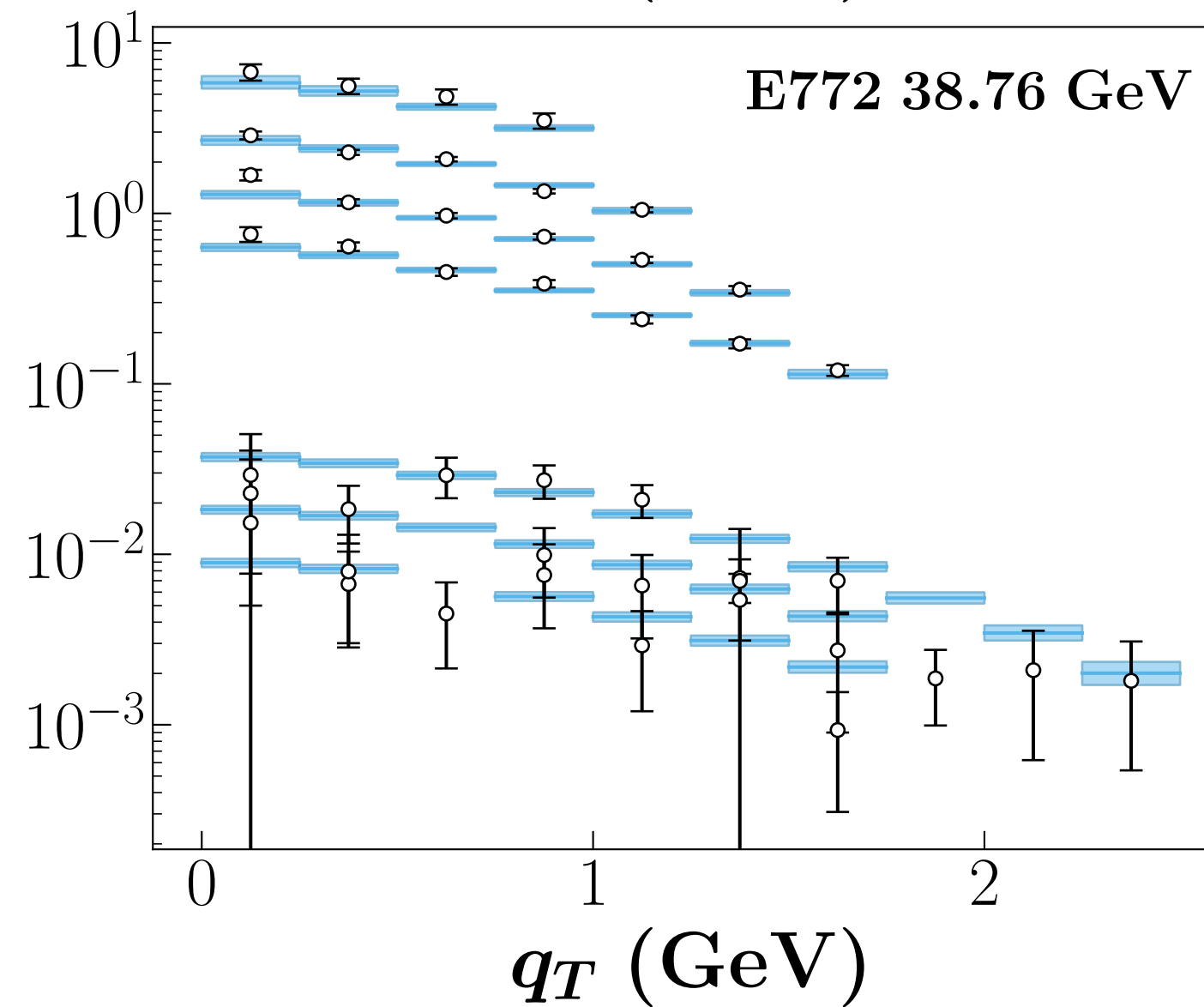
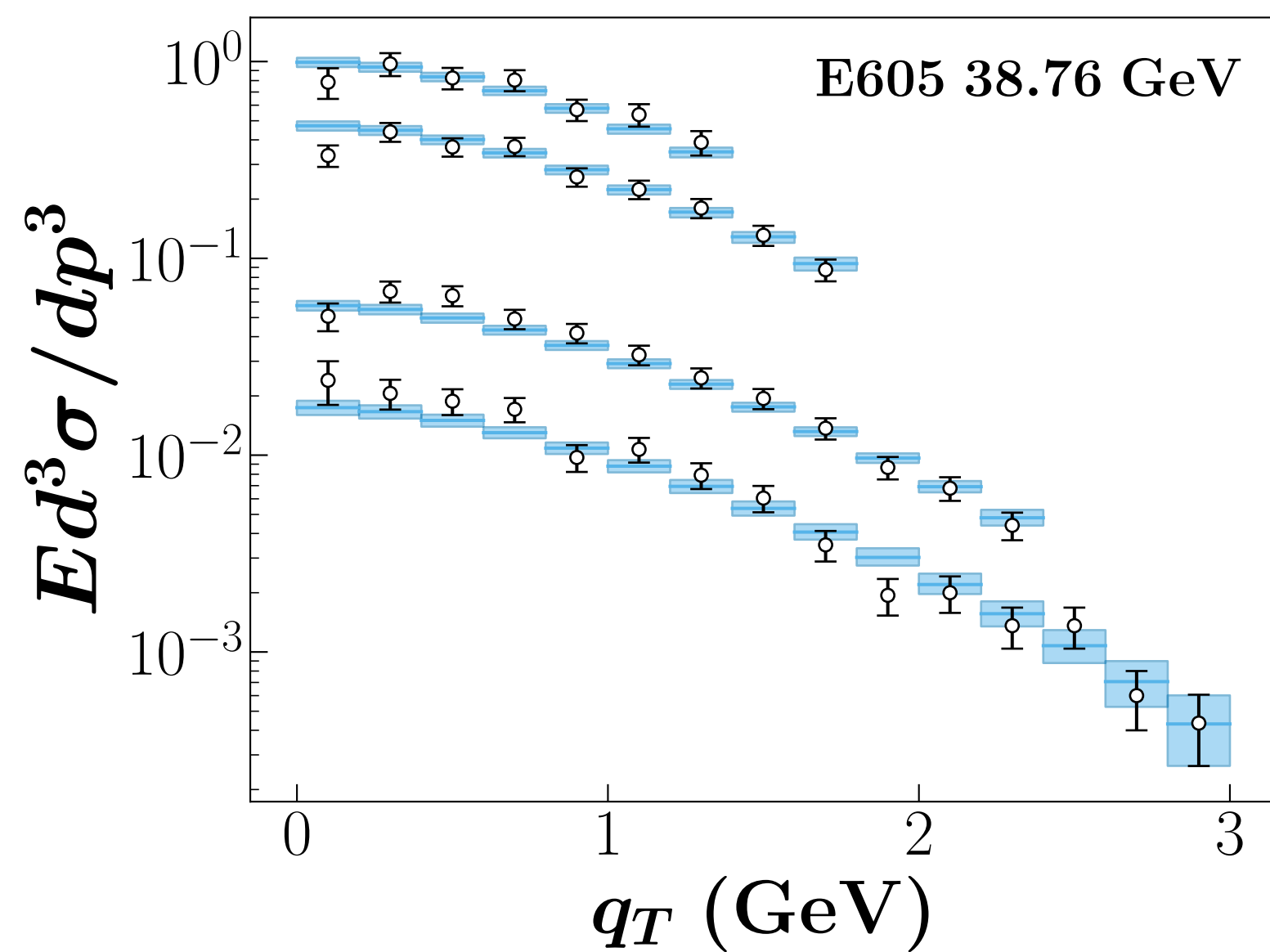
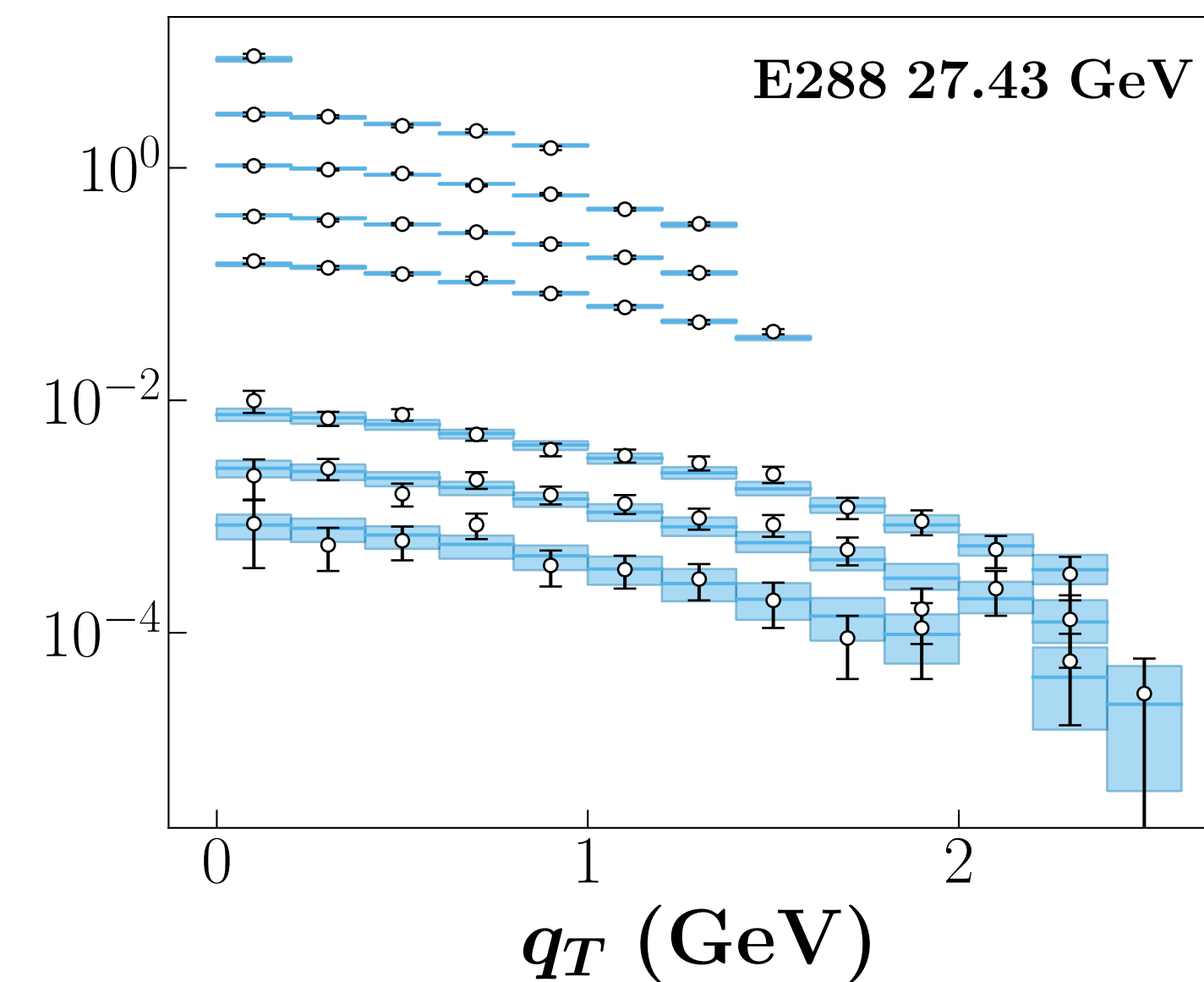
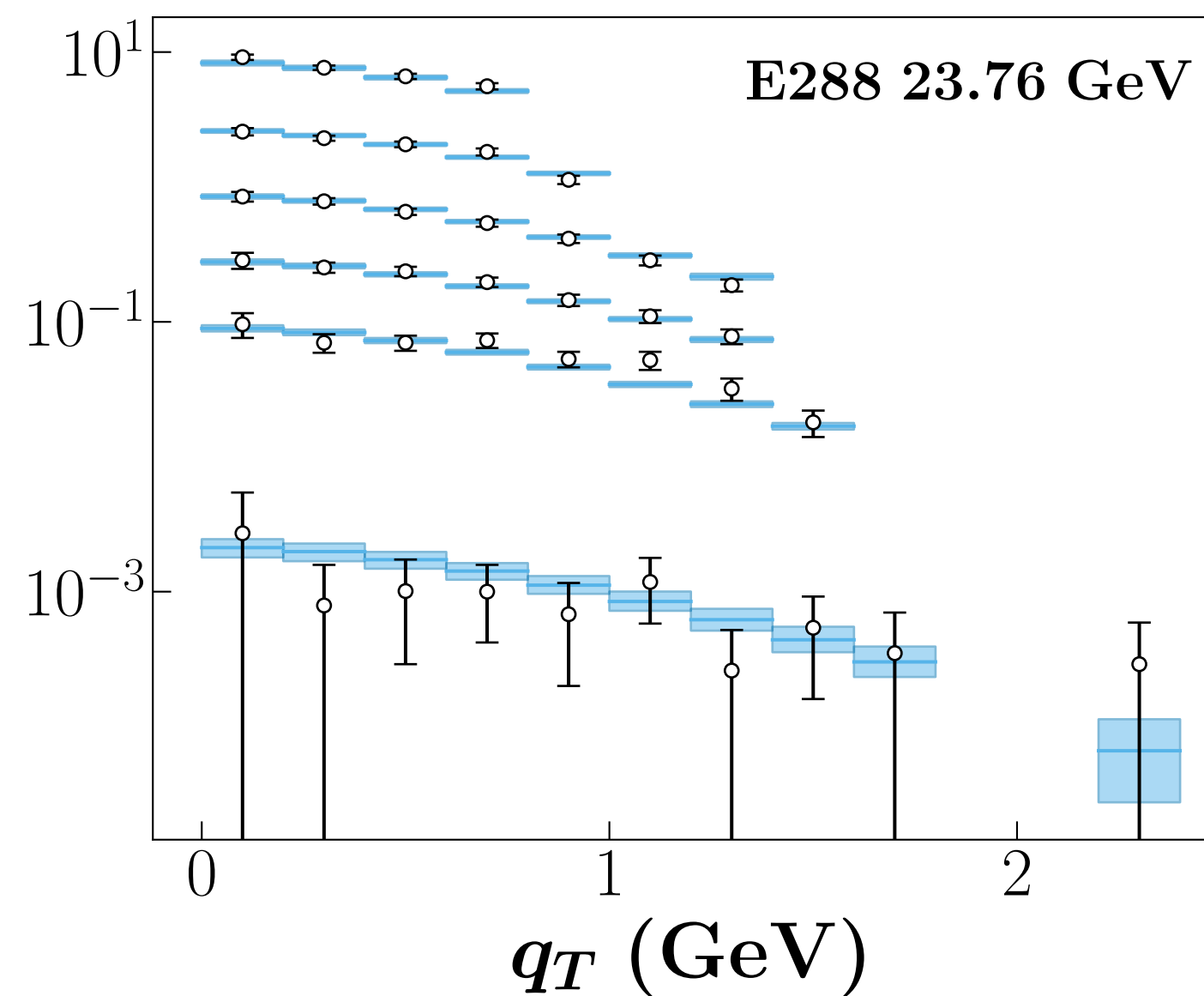
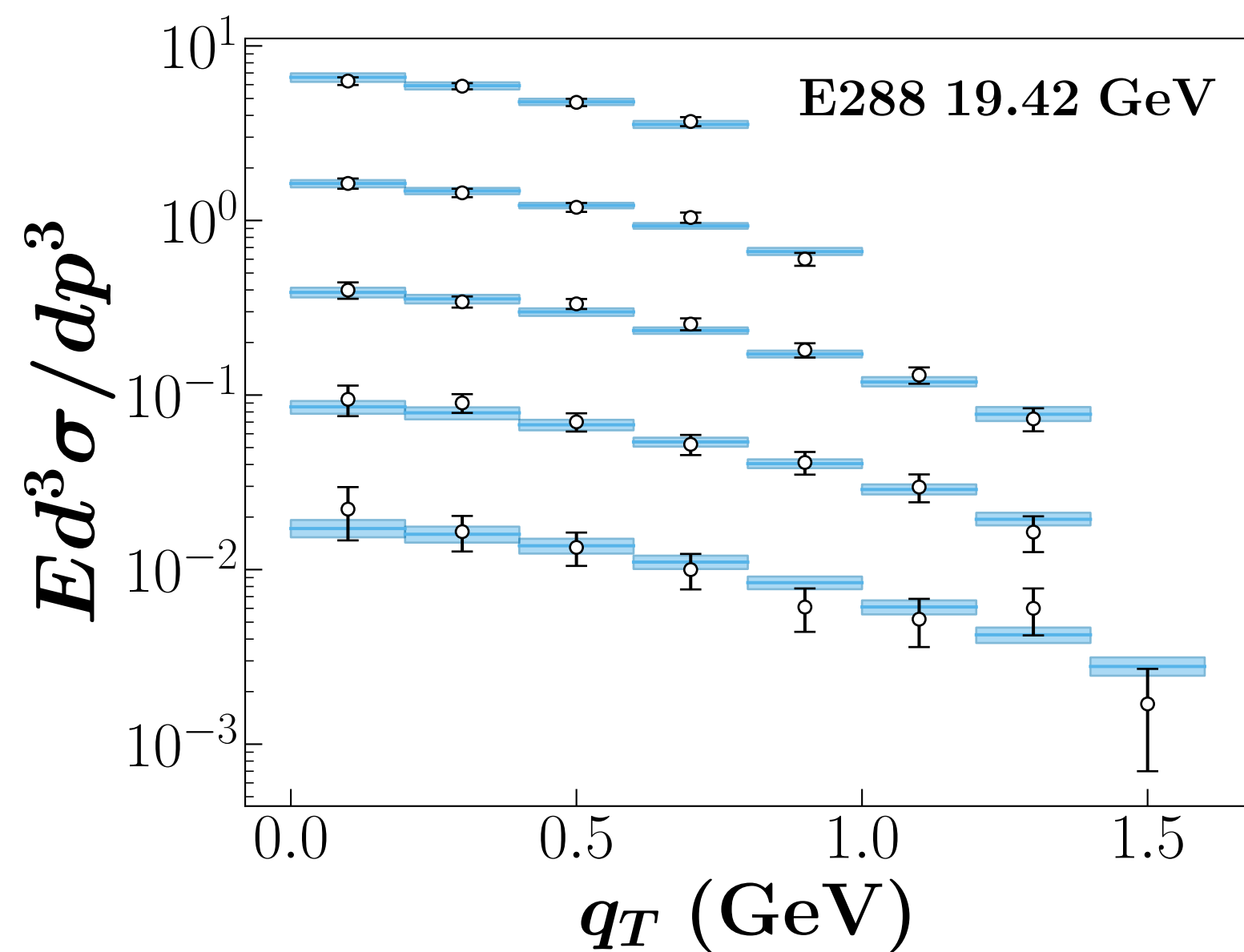
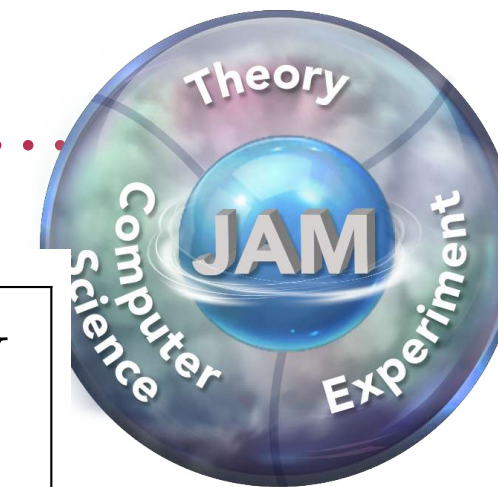
LHCb data are integrated in rapidity

# RESULTS



CMS data are in bins of rapidity

# RESULTS



E288, E605 data are sensitive to high- $x$  behavior of PDFs and TMDs