



Transverse diffeomorphisms

basic ideas &
cosmological consequences

11.12.24

Darío Jaramillo Garrido

Diego Tessainer Bonet

Departamento de
Física Teórica



OUTLINE

DARÍO

1

TDiff: Basic ideas

Refs.

[2307.14861]

[2402.17422]

Motivation | Building the theory | Symmetry breaking |
New interactions | EMT conservation | Constraint



OUTLINE

DARÍO

1

TDiff: Basic ideas

Refs.

[2307.14861]

[2402.17422]

Motivation | Building the theory | Symmetry breaking |
New interactions | EMT conservation | Constraint

DIEGO

2

Cosmological consequences

Refs.

[2409.11991]

Multi-field model | Effective interactions | Dark sector |
Analysis | Conclusions

①

TDiff:
Basic ideas

Motivation

- Newsflash!

Motivation

- Newsflash!
 - The Universe presents an ***accelerated expansion***.
 - To form cosmic structures, we need ***more matter*** than we see.

Motivation

- Newsflash!

- The Universe presents *dark energy* **accelerated expansion**.
- To form cosmic structures, we need **more matter** than we see.

Motivation

- Newsflash!

- The Universe presents **dark energy** *accelerated expansion*.
- To form cosmic structures **dark matter** *more matter* than we see.

Motivation

- Newsflash!

- The Universe presents a *cosmological expansion*.

- To form cosmic structures, we need *more matter* than we see.

DARK SECTOR

dark energy

dark matter

expansion.

re matter than we see.

Motivation

- Newsflash!

- The Universe presents
- To form cosmic structures



expansion.

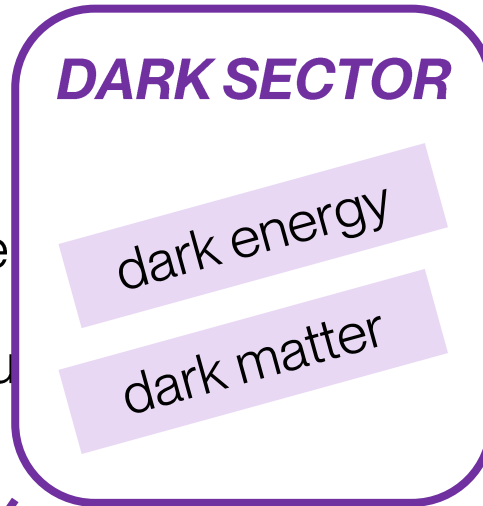
re matter than we see.

What is its nature?

Motivation

- Newsflash!

- The Universe presents
- To form cosmic structures



expansion.

are matter than we see.

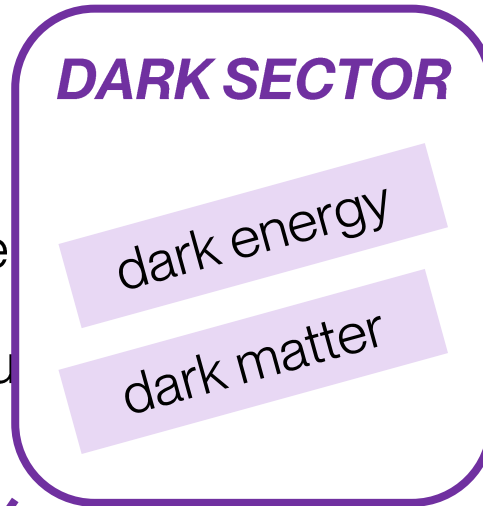
What is its nature?

Is there interaction?

Motivation

- Newsflash!

- The Universe presents **dark energy** to drive **expansion**.
- To form cosmic structures, we need **dark matter** more **matter** than we see.



What is its nature?

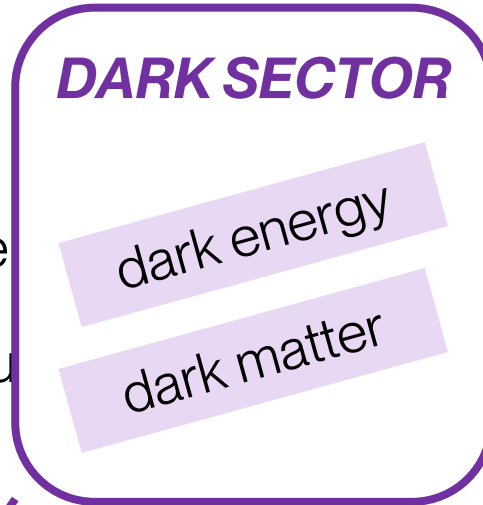
Is there interaction?

Theory modification?

Motivation

- Newsflash!

- The Universe presents a puzzle: *accelerated expansion.*
- To form cosmic structures, we need *more matter* than we see.



What is its nature?

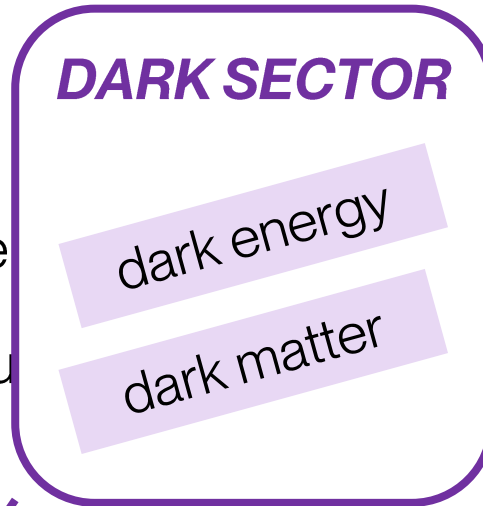
Is there interaction?

Theory modification?

Motivation

- Newsflash!

- The Universe presents **dark energy** for **accelerated expansion**.
- To form cosmic structures, we need **dark matter** **more matter** than we see.



- What are the symmetries of the dark sector of the Universe?

We will consider the possibility of **broken symmetries** in the dark sector.

Building the theory

- In general relativity, the ***diffeomorphism*** symmetry fixes the ***coupling*** between gravity and a matter field:
(... at least in the simplest case of “minimal coupling”)

Building the theory

- In general relativity, the ***diffeomorphism*** symmetry fixes the ***coupling*** between gravity and a matter field:

$$S[g_{\mu\nu}, \text{matter}] = \int_{\mathcal{M}} d^4x \sqrt{g} \mathcal{L}(\text{matter})$$

Building the theory

- In general relativity, the ***diffeomorphism*** symmetry fixes the ***coupling*** between gravity and a matter field:

$$S[g_{\mu\nu}, \text{matter}] = \int_{\mathcal{M}} d^4x \sqrt{g} \mathcal{L}(\text{matter})$$

$g_{\mu\nu}$: metric
("gravitational field")

Building the theory

- In general relativity, the **diffeomorphism** symmetry fixes the **coupling** between gravity and a matter field:

$$S[g_{\mu\nu}, \text{matter}] = \int_{\mathcal{M}} d^4x \sqrt{g} \mathcal{L}(\text{matter})$$

$g_{\mu\nu}$: metric
("gravitational field")

$g = |\det(g_{\mu\nu})|$

Building the theory

- In general relativity, the **diffeomorphism** symmetry fixes the **coupling** between gravity and a matter field:

$$S[g_{\mu\nu}, \text{matter}] = \int_{\mathcal{M}} d^4x \sqrt{g} \mathcal{L}(\text{matter})$$

$g_{\mu\nu}$: metric
("gravitational field")

$g = |\det(g_{\mu\nu})|$

Building the theory

- In general relativity, the **diffeomorphism** symmetry fixes the **coupling** between gravity and a matter field:

$$S[g_{\mu\nu}, \text{matter}] = \int_{\mathcal{M}} d^4x \sqrt{g} \mathcal{L}(\text{matter})$$

$g_{\mu\nu}$: metric
("gravitational field")

$g = |\det(g_{\mu\nu})|$

- Question:

*what happens if we **modify** the coupling between fields and gravity?*

Building the theory

- In general relativity, the **diffeomorphism** symmetry fixes the **coupling** between gravity and a matter field:

$$S[g_{\mu\nu}, \text{matter}] = \int_{\mathcal{M}} d^4x \sqrt{g} \mathcal{L}(\text{matter})$$

$g_{\mu\nu}$: metric
("gravitational field")

$g = |\det(g_{\mu\nu})|$

- Question:

*what happens if we **modify** the coupling between fields and gravity?*

$$\sqrt{g} \longrightarrow f(g)$$

Building the theory

- In general relativity, the symmetry fixes the *coupling* between gravity and the fields:

1. We will modify the *interaction* between gravity and the fields.

2. We will modify the *symmetries* of the theory.

- Question:

*what happens if we **modify** the coupling between fields and gravity?*

$$\sqrt{g} \longrightarrow f(g)$$

Symmetry breaking

2. We will modify the *symmetries* of the theory.

Symmetry breaking

2. We will modify the *symmetries* of the theory.

- Before:
$$S = \int d^4x \sqrt{g} \mathcal{L}$$

Symmetry breaking

2. We will modify the *symmetries* of the theory.

- Before: $S = \int d^4x \sqrt{g} \mathcal{L} \rightarrow$ Diffeomorphisms (*Diff*)

Symmetry breaking

2. We will modify the *symmetries* of the theory.

- Before: $S = \int d^4x \sqrt{g} \mathcal{L} \rightarrow$ Diffeomorphisms (*Diff*)



$$x \rightarrow \hat{x}$$

Symmetry breaking

2. We will modify the *symmetries* of the theory.

• Before: $S = \int d^4x \sqrt{g} \mathcal{L} \rightarrow$ Diffeomorphisms (*Diff*)

• Now: $S = \int d^4x f(g) \mathcal{L}$

Diff

$$x \rightarrow \hat{x}$$

Symmetry breaking

2. We will modify the *symmetries* of the theory.

• Before: $S = \int d^4x \sqrt{g} \mathcal{L} \rightarrow$ Diffeomorphisms (*Diff*)

• Now: $S = \int d^4x f(g) \mathcal{L} \rightarrow$ Transverse Diffeomorphisms (*TDiff*)



$$x \rightarrow \hat{x}$$

Symmetry breaking

2. We will modify the *symmetries* of the theory.

• Before: $S = \int d^4x \sqrt{g} \mathcal{L} \rightarrow$ Diffeomorphisms (*Diff*)

• Now: $S = \int d^4x f(g) \mathcal{L} \rightarrow$ Transverse Diffeomorphisms (*TDiff*)



$$x \rightarrow \hat{x}$$
$$J = \det \left(\frac{\partial \hat{x}}{\partial x} \right) = 1$$
$$\hat{x}^\mu = x^\mu + \xi^\mu, \quad \partial_\mu \xi^\mu = 0$$

New interactions

- 1.** We will modify the *interaction* between gravity and the fields.

New interactions

1. We will modify the *interaction* between gravity and the fields.

- Simplest matter to study: ***scalar fields*** $\psi(x)$

New interactions

1. We will modify the *interaction* between gravity and the fields.

- Simplest matter to study: ***scalar fields*** $\psi(x)$

$$S_{\text{TDiff}} = S_{\text{G}} + S_{\text{M}} =$$

New interactions

1. We will modify the *interaction* between gravity and the fields.

- Simplest matter to study: ***scalar fields*** $\psi(x)$

$$\begin{aligned} S_{\text{TDiff}} &= S_{\text{G}} + S_{\text{M}} = \\ &= \int d^4x \sqrt{g} R + \\ &\quad \text{(unchanged} \\ &\quad \text{Einstein-Hilbert)} \end{aligned}$$

New interactions

1. We will modify the *interaction* between gravity and the fields.

- Simplest matter to study: **scalar fields** $\psi(x)$

$$\begin{aligned} S_{\text{TDiff}} &= S_{\text{G}} + S_{\text{M}} = \\ &= \int d^4x \sqrt{g} R + \int d^4x \left\{ f_k(g) \frac{(\partial\psi)^2}{2} - f_v(g) V(\psi) \right\} \\ &\quad \text{(unchanged Einstein-Hilbert)} \end{aligned}$$

New interactions

1. We will modify the *interaction* between gravity and the fields.

- Simplest matter to study: **scalar fields** $\psi(x)$

$$\begin{aligned} S_{\text{TDiff}} &= S_{\text{G}} + S_{\text{M}} = \\ &= \int d^4x \sqrt{g} R + \int d^4x \left\{ f_k(g) \frac{(\partial\psi)^2}{2} - f_v(g) V(\psi) \right\} \end{aligned}$$

- We obtain modified equations of motion:

$$\text{Einstein:} \quad G_{\mu\nu} = 8\pi G T_{\mu\nu} \quad T_{\mu\nu} = \frac{2}{\sqrt{g}} \frac{\delta S_{\text{M}}}{\delta g^{\mu\nu}}$$

$$\text{Klein-Gordon:} \quad \partial_\mu [f_k(g) \partial^\mu \psi] + f_v(g) V'(\psi) = 0$$

New interactions

1. We will modify the *interaction* between gravity and the fields.

- Simplest matter to study: **scalar fields** $\psi(x)$

$$\begin{aligned} S_{\text{TDiff}} &= S_{\text{G}} + S_{\text{M}} = \\ &= \int d^4x \sqrt{g} R + \int d^4x \left\{ f_k(g) \frac{(\partial\psi)^2}{2} - f_v(g) V(\psi) \right\} \end{aligned}$$

- We obtain modified equations of motion:

Einstein: $G_{\mu\nu} = 8\pi G T_{\mu\nu}$  energy-momentum tensor (EMT)

Klein-Gordon: $\partial_\mu [f_k(g) \partial^\mu \psi] + f_v(g) V'(\psi) = 0$

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)

EMT conservation

(Noether)

- Diff invariance \Rightarrow EMT conservation.
- TDiff: *less symmetry!*

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)
- TDiff: ***less symmetry!***

however, we still have the Einstein equations...

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)
- TDiff: ***less symmetry!***

however, we still have the Einstein equations...

- EMT conservation follows from the ***Bianchi identities:***

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)
- TDiff: ***less symmetry!***

however, we still have the Einstein equations...

- EMT conservation follows from the ***Bianchi identities:***

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\nabla_{\mu} G^{\mu\nu} = 0$$

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)
- TDiff: ***less symmetry!***

however, we still have the Einstein equations...

- EMT conservation follows from the ***Bianchi identities***:

$$\left. \begin{aligned} G_{\mu\nu} &= 8\pi G T_{\mu\nu} \\ \nabla_{\mu} G^{\mu\nu} &= 0 \end{aligned} \right\} \Rightarrow \nabla_{\mu} T^{\mu\nu} = 0$$

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)
- TDiff: ***less symmetry!***

however, we still have the Einstein equations...

- EMT conservation follows from the ***Bianchi identities***:

$$\left. \begin{aligned} G_{\mu\nu} &= 8\pi G T_{\mu\nu} \\ \nabla_{\mu} G^{\mu\nu} &= 0 \end{aligned} \right\} \Rightarrow \nabla_{\mu} T^{\mu\nu} = 0$$

Trivial in GR,
not anymore!

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)
- TDiff: ***less symmetry!***

however, we still have the Einstein equations...

- EMT conservation follows from the ***Bianchi identities***:

$$\left. \begin{aligned} G_{\mu\nu} &= 8\pi G T_{\mu\nu} \\ \nabla_{\mu} G^{\mu\nu} &= 0 \end{aligned} \right\} \Rightarrow \nabla_{\mu} T^{\mu\nu} = 0$$

Trivial in GR,
not anymore!

\Rightarrow ***new information*** in the form of an equation for g (“constraint”)

EMT conservation

- Diff invariance \Rightarrow EMT conservation. (Noether)
- TDiff: **less symmetry!**

however, we still have the Einstein equations...

- EMT conservation follows from the **Bianchi identities**:

$$\left. \begin{aligned} G_{\mu\nu} &= 8\pi G T_{\mu\nu} \\ \nabla_{\mu} G^{\mu\nu} &= 0 \end{aligned} \right\} \Rightarrow \nabla_{\mu} T^{\mu\nu} = 0$$

Trivial in GR,
not anymore!

\Rightarrow **new information** in the form of an equation for g (“constraint”)

! This is true for any TDiff theory where the gravitational action is the Einstein-Hilbert one

Constraint & phenomenology

- The ***general constraint*** in a scalar field theory reads:

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

Constraint & phenomenology

- The ***general constraint*** in a scalar field theory reads:

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

➤ formal expression – need the equations of motion

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

Diff

Theory

EoS

Sound speed

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

	Diff
Theory	$\mathcal{L}_{\psi} = \sqrt{g} \frac{(\partial\psi)^2}{2}$
EoS	$w = 1$
Sound speed	$c_s^2 = 1$

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_\mu T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

	Diff	TDiff
Theory	$\mathcal{L}_\psi = \sqrt{g} \frac{(\partial\psi)^2}{2}$	
EoS	$w = 1$	
Sound speed	$c_s^2 = 1$	

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_\mu T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

	Diff	TDiff
Theory	$\mathcal{L}_\psi = \sqrt{g} \frac{(\partial\psi)^2}{2}$	$\mathcal{L}_\psi = f_k(g) \frac{(\partial\psi)^2}{2}$
EoS	$w = 1$	
Sound speed	$c_s^2 = 1$	

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_\mu T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

	Diff	TDiff
Theory	$\mathcal{L}_\psi = \sqrt{g} \frac{(\partial\psi)^2}{2}$	$\mathcal{L}_\psi = f_k(g) \frac{(\partial\psi)^2}{2}$
EoS	$w = 1$	$w = \frac{gf'_k}{f_k - gf'_k} = w(g)$
Sound speed	$c_s^2 = 1$	

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_\mu T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

	Diff	TDiff
Theory	$\mathcal{L}_\psi = \sqrt{g} \frac{(\partial\psi)^2}{2}$	$\mathcal{L}_\psi = f_k(g) \frac{(\partial\psi)^2}{2}$
EoS	$w = 1$	$w = \frac{gf'_k}{f_k - gf'_k} = w(g)$
Sound speed	$c_s^2 = 1$	$c_s^2 = \frac{-gf_k(f'_k + 2gf''_k)}{(f_k - 2gf'_k)^2 - gf_k(f'_k + 2gf''_k)}$

Constraint & phenomenology

- The **general constraint** in a scalar field theory reads:

$$\nabla_\mu T^{\mu\nu} = 0 \quad \Rightarrow \quad [f_k - 2gf'_k] \frac{(\partial\psi)^2}{2} - [f_v - 2gf'_v] V(\psi) = \text{const}$$

(Diego) \leftarrow \rightarrow formal expression – need the equations of motion

Example: shift symmetry. ($V = 0$)

	Diff	TDiff
Theory	$\mathcal{L}_\psi = \sqrt{g} \frac{(\partial\psi)^2}{2}$	$\mathcal{L}_\psi = f_k(g) \frac{(\partial\psi)^2}{2}$
EoS	$w = 1$	$w = \frac{gf'_k}{f_k - gf'_k} = w(g)$
Sound speed	$c_s^2 = 1$	$c_s^2 = \frac{-gf_k(f'_k + 2gf''_k)}{(f_k - 2gf'_k)^2 - gf_k(f'_k + 2gf''_k)}$

new phenomenology!

②

*Cosmological
consequences*

Multi-field model

- We will consider a *shift-symmetric* model involving two free scalar fields invariant under TDiff transformations:

Multi-field model

- We will consider a *shift-symmetric* model involving two **free scalar fields** invariant under TDiff transformations:

$$S_M = \sum_{i=1}^2 \int d^4x \left(\frac{1}{2} f_{k_i}(g) \partial_\mu \psi_i \partial^\mu \psi_i \right)$$

Multi-field model

- We will consider a *shift-symmetric* model involving two free scalar fields invariant under TDiff transformations:

$$S_M = \sum_{i=1}^2 \int d^4x \left(\frac{1}{2} f_{k_i}(g) \partial_\mu \psi_i \partial^\mu \psi_i \right)$$

Free fields with power-law coupling functions

Multi-field model

- We will consider a *shift-symmetric* model involving two free scalar fields invariant under TDiff transformations:

$$S_M = \sum_{i=1}^2 \int d^4x \left(\frac{1}{2} f_{k_i}(g) \partial_\mu \psi_i \partial^\mu \psi_i \right)$$

Free fields with power-law coupling functions $\longrightarrow f_{k_i}(g) = \lambda_i g^{\alpha_i}$

Multi-field model

- We will consider a *shift-symmetric* model involving two free scalar fields invariant under TDiff transformations:

$$S_M = \sum_{i=1}^2 \int d^4x \left(\frac{1}{2} f_{k_i}(g) \partial_\mu \psi_i \partial^\mu \psi_i \right)$$

Free fields with power-law coupling functions $\longrightarrow f_{k_i}(g) = \lambda_i g^{\alpha_i}$

- Spatially flat Robertson-Walker background:

Multi-field model

- We will consider a *shift-symmetric* model involving two **free scalar fields** invariant under TDiff transformations:

$$S_M = \sum_{i=1}^2 \int d^4x \left(\frac{1}{2} f_{k_i}(g) \partial_\mu \psi_i \partial^\mu \psi_i \right)$$

Free fields with power-law coupling functions $\longrightarrow f_{k_i}(g) = \lambda_i g^{\alpha_i}$

- Spatially flat Robertson-Walker background:

$$ds^2 = b(\tau)^2 d\tau^2 - a(\tau)^2 d\mathbf{x}^2$$

Multi-field model

- We will consider a *shift-symmetric* model involving two free scalar fields invariant under TDiff transformations:

$$S_M = \sum_{i=1}^2 \int d^4x \left(\frac{1}{2} f_{k_i}(g) \partial_\mu \psi_i \partial^\mu \psi_i \right)$$

Free fields with power-law coupling functions $\longrightarrow f_{k_i}(g) = \lambda_i g^{\alpha_i}$

- Spatially flat Robertson-Walker background:

$$ds^2 = b(\tau)^2 d\tau^2 - a(\tau)^2 d\mathbf{x}^2$$

! A new physical degree of freedom appears in the metric as a consequence of having less symmetry

Multi-field model

- We will consider a *shift-symmetric* model involving two free scalar fields invariant under TDiff transformations:

$$S_M = \sum_{i=1}^2 \int d^4x \left(\frac{1}{2} f_{k_i}(g) \partial_\mu \psi_i \partial^\mu \psi_i \right)$$

Free fields with power-law coupling functions $\longrightarrow f_{k_i}(g) = \lambda_i g^{\alpha_i}$

- Spatially flat Robertson-Walker background:

$$ds^2 = b(\tau)^2 d\tau^2 - a(\tau)^2 d\mathbf{x}^2$$

! A new physical degree of freedom appears in the metric as a consequence of having less symmetry

$$\partial_\mu \xi^\mu = 0$$

Multi-field model

- We will regard the perfect fluid approach:

Multi-field model

- We will regard the perfect fluid approach:

$$T_{\mu\nu} = (\rho + p)u_{\mu}u_{\nu} - pg_{\mu\nu}$$

Multi-field model

- We will regard the perfect fluid approach:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu - pg_{\mu\nu}$$

$$A = 1, 2$$

$$\rho_A = C_A(1 - \alpha_A) \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

$$p_A = C_A \alpha_A \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

Multi-field model

- We will regard the perfect fluid approach:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu - pg_{\mu\nu}$$

$$A = 1, 2$$

$$\rho_A = C_A(1 - \alpha_A) \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$
$$p_A = C_A\alpha_A \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

➤ The equation of state parameters read \Rightarrow

$$w_A = \frac{p_A}{\rho_A} = \frac{\alpha_A}{1 - \alpha_A}$$

Multi-field model

- We will regard the perfect fluid approach:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu - pg_{\mu\nu}$$

$$A = 1, 2$$

$$\rho_A = C_A(1 - \alpha_A) \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$
$$p_A = C_A\alpha_A \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

➤ The equation of state parameters read

$$w_A = \frac{p_A}{\rho_A} = \frac{\alpha_A}{1 - \alpha_A}$$

- The dynamics are given by Einstein Equations and the conservation law for the total EMT:

Multi-field model

- We will regard the perfect fluid approach:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu - pg_{\mu\nu}$$

$$A = 1, 2$$

$$\rho_A = C_A(1 - \alpha_A) \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$
$$p_A = C_A\alpha_A \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

➤ The equation of state parameters read \Rightarrow

$$w_A = \frac{p_A}{\rho_A} = \frac{\alpha_A}{1 - \alpha_A}$$

- The dynamics are given by Einstein Equations and the conservation law for the total EMT:

$$\rho_{\text{tot}} = \rho_1 + \rho_2$$

$$\left(\frac{a'}{a}\right)^2 = \frac{8\pi G}{3}\rho_{\text{tot}}b(\tau)^2$$

Multi-field model

- We will regard the perfect fluid approach:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu - pg_{\mu\nu}$$

$$A = 1, 2$$

$$\rho_A = C_A(1 - \alpha_A) \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$
$$p_A = C_A\alpha_A \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

➤ The equation of state parameters read \Rightarrow

$$w_A = \frac{p_A}{\rho_A} = \frac{\alpha_A}{1 - \alpha_A}$$

- The dynamics are given by Einstein Equations and the conservation law for the total EMT:

$$\rho_{\text{tot}} = \rho_1 + \rho_2$$

$$\left(\frac{a'}{a}\right)^2 = \frac{8\pi G}{3}\rho_{\text{tot}}b(\tau)^2$$

+

$$\rho'_{\text{tot}} + 3\frac{a'}{a}(\rho_{\text{tot}} + p_{\text{tot}}) = 0$$

Multi-field model

- We will regard the perfect fluid approach:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu - pg_{\mu\nu}$$

$$A = 1, 2$$

$$\rho_A = C_A(1 - \alpha_A) \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

$$p_A = C_A\alpha_A \frac{b(\tau)^{1-2\alpha_A}}{a(\tau)^{6\alpha_A+3}}$$

➤ The equation of state parameters read \Rightarrow

$$w_A = \frac{p_A}{\rho_A} = \frac{\alpha_A}{1 - \alpha_A}$$

- The dynamics are given by Einstein Equations and the conservation law for the total EM

$$\left(\frac{a'}{a}\right)^2 = \frac{8\pi G}{3} \rho_{\text{tot}} b(\tau)^2$$

+

$$\rho'_{\text{tot}} + 3\frac{a'}{a}(\rho_{\text{tot}} + p_{\text{tot}}) = 0$$

This will provide the geometrical constraint that will be responsible for the interactions

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

$$C_1 g^{1-\alpha_1} |2\alpha_1 - 1| + C_2 g^{1-\alpha_2} |2\alpha_2 - 1| = C_g a^6 \quad \Rightarrow \quad b(\tau) \rightarrow b(a(\tau))$$

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

$$C_1 g^{1-\alpha_1} |2\alpha_1 - 1| + C_2 g^{1-\alpha_2} |2\alpha_2 - 1| = C_g a^6 \quad \Rightarrow \quad b(\tau) \rightarrow b(a(\tau))$$



Contributions from both components, the geometrical constraint will be different than in the single-field case

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

$$C_1 g^{1-\alpha_1} |2\alpha_1 - 1| + C_2 g^{1-\alpha_2} |2\alpha_2 - 1| = C_g a^6 \quad \Rightarrow \quad b(\tau) \rightarrow b(a(\tau))$$

$$\begin{aligned} \nabla_\mu T_A^{\mu\nu} &\neq 0 \\ \nabla_\mu (T_1^{\mu\nu} + T_2^{\mu\nu}) &= 0 \end{aligned}$$



Contributions from both components, the geometrical constraint will be different than in the single-field case

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

$$C_1 g^{1-\alpha_1} |2\alpha_1 - 1| + C_2 g^{1-\alpha_2} |2\alpha_2 - 1| = C_g a^6 \quad \Rightarrow \quad b(\tau) \rightarrow b(a(\tau))$$

$$\begin{aligned} \nabla_\mu T_A^{\mu\nu} &\neq 0 \\ \nabla_\mu (T_1^{\mu\nu} + T_2^{\mu\nu}) &= 0 \end{aligned}$$



Contributions from both components, the geometrical constraint will be different than in the single-field case

- The single-field domination regimes allow us to study the phenomenological possibilities of the model

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

$$C_1 g^{1-\alpha_1} |2\alpha_1 - 1| + C_2 g^{1-\alpha_2} |2\alpha_2 - 1| = C_g a^6 \quad \Rightarrow \quad b(\tau) \rightarrow b(a(\tau))$$

$$\begin{aligned} \nabla_\mu T_A^{\mu\nu} &\neq 0 \\ \nabla_\mu (T_1^{\mu\nu} + T_2^{\mu\nu}) &= 0 \end{aligned}$$



Contributions from both components, the geometrical constraint will be different than in the single-field case

- The single-field domination regimes allow us to study the phenomenological possibilities of the model

➤ ψ_1 is dominant \Rightarrow

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

$$C_1 g^{1-\alpha_1} |2\alpha_1 - 1| + C_2 g^{1-\alpha_2} |2\alpha_2 - 1| = C_g a^6 \quad \Rightarrow \quad b(\tau) \rightarrow b(a(\tau))$$

$$\begin{aligned} \nabla_\mu T_A^{\mu\nu} &\neq 0 \\ \nabla_\mu (T_1^{\mu\nu} + T_2^{\mu\nu}) &= 0 \end{aligned}$$



Contributions from both components, the geometrical constraint will be different than in the single-field case

- The single-field domination regimes allow us to study the phenomenological possibilities of the model

➤ ψ_1 is dominant \Rightarrow

$$\begin{aligned} \rho_1(a) &\propto a^{-3(1+w_1)} \\ \rho_2(a) &\propto a^{-3(1+w_{\text{eff}})} \end{aligned}$$

Effective interactions

- The conservation law for the total EMT yields the following geometrical constraint:

$$C_1 g^{1-\alpha_1} |2\alpha_1 - 1| + C_2 g^{1-\alpha_2} |2\alpha_2 - 1| = C_g a^6 \quad \Rightarrow \quad b(\tau) \rightarrow b(a(\tau))$$

$$\begin{aligned} \nabla_\mu T_A^{\mu\nu} &\neq 0 \\ \nabla_\mu (T_1^{\mu\nu} + T_2^{\mu\nu}) &= 0 \end{aligned}$$



Contributions from both components, the geometrical constraint will be different than in the single-field case

- The single-field domination regimes allow us to study the phenomenological possibilities of the model

➤ ψ_1 is dominant \Rightarrow

$$\begin{aligned} \rho_1(a) &\propto a^{-3(1+w_1)} \\ \rho_2(a) &\propto a^{-3(1+w_{\text{eff}})} \end{aligned}$$

$$w_{\text{eff}} = \frac{2w_2 - w_1 + w_1 w_2}{1 + w_2}$$

Effective interactions

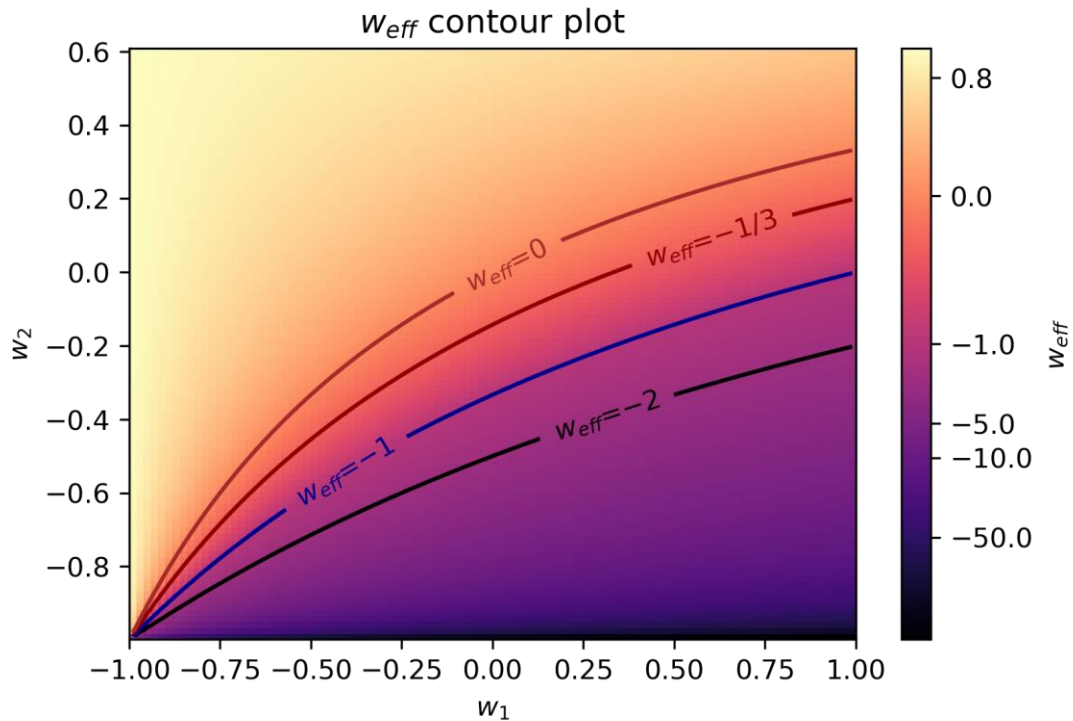
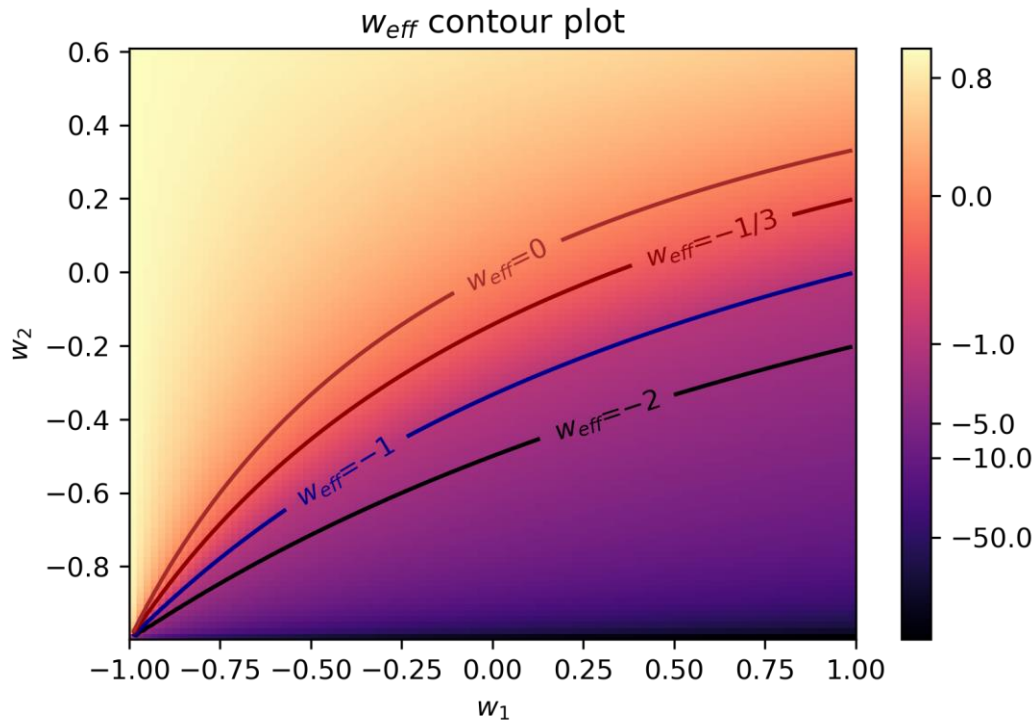


Fig.1. Effective equation of state parameter for the subdominant field in terms of w_1 and w_2 .

Effective interactions



- Wide range of possible phenomenology.
- Possible to have a cosmological constant in subdominant regimes.
- Possibility of phantom-crossing

Fig.1. Effective equation of state parameter for the subdominant field in terms of w_1 and w_2 .

Dark sector

- We consider a generic dark sector model:

Dark sector

- We consider a generic dark sector model:

$$\alpha_1 = 0 \implies w_1 = 0$$

Dark sector

- We consider a generic dark sector model:

$$\boxed{\alpha_1 = 0 \implies w_1 = 0} \implies \text{our parameters are } \left\{ \begin{array}{l} \alpha_2 \quad (w_2) \\ \gamma = \frac{\Omega_{\text{DM}}}{\Omega_{\text{DE}}} \end{array} \right.$$

Dark sector

- We consider a generic dark sector model:

$$\boxed{\alpha_1 = 0 \implies w_1 = 0} \implies \text{our parameters are } \left\{ \begin{array}{l} \alpha_2 \quad (w_2) \\ \gamma = \frac{\Omega_{\text{DM}}}{\Omega_{\text{DE}}} \end{array} \right.$$

$< -1/3$

Dark sector

- We consider a generic dark sector model:

$$\boxed{\alpha_1 = 0 \implies w_1 = 0} \implies \text{our parameters are } \left\{ \begin{array}{l} \alpha_2 \quad (w_2) \\ \gamma = \frac{\Omega_{\text{DM}}}{\Omega_{\text{DE}}} \end{array} \right. \quad \begin{array}{l} < -1/3 \end{array}$$

- The Friedmann equation allows to obtain the Hubble rate:

Dark sector

- We consider a generic dark sector model:

$$\alpha_1 = 0 \implies w_1 = 0 \implies \text{our parameters are } \left\{ \begin{array}{l} \alpha_2 \quad (w_2) \\ \gamma = \frac{\Omega_{\text{DM}}}{\Omega_{\text{DE}}} \end{array} \right. \quad \triangleleft -1/3$$

- The Friedmann equation allows to obtain the Hubble rate:

$$H^2(z) = H_0^2 \left[\Omega_B(1+z)^3 + (1 - \Omega_B) \left(1 + \frac{1}{\gamma}\right)^{-1} b(z)(1+z)^3 + \frac{1}{\gamma}(1 - \Omega_B) \left(1 + \frac{1}{\gamma}\right)^{-1} b(z)^{1-2\alpha_2} (1+z)^{6\alpha_2+3} \right]$$

Dark sector

- We consider a generic dark sector model:

$$\boxed{\alpha_1 = 0 \implies w_1 = 0} \implies \text{our parameters are } \left\{ \begin{array}{l} \alpha_2 \quad (w_2) \\ \gamma = \frac{\Omega_{\text{DM}}}{\Omega_{\text{DE}}} \end{array} \right. \quad \triangleleft -1/3$$

- The Friedmann equation allows to obtain the Hubble rate:

$$H^2(z) = H_0^2 \left[\Omega_B(1+z)^3 + (1 - \Omega_B) \left(1 + \frac{1}{\gamma}\right)^{-1} b(z)(1+z)^3 + \frac{1}{\gamma}(1 - \Omega_B) \left(1 + \frac{1}{\gamma}\right)^{-1} b(z)^{1-2\alpha_2}(1+z)^{6\alpha_2+3} \right]$$

- We will use the matter abundance at high redshift

Dark sector

- We consider a generic dark sector model:

$$\alpha_1 = 0 \implies w_1 = 0 \implies \text{our parameters are } \left\{ \begin{array}{l} \alpha_2 \quad (w_2) \\ \gamma = \frac{\Omega_{\text{DM}}}{\Omega_{\text{DE}}} \end{array} \right. \quad \triangleleft -1/3$$

- The Friedmann equation allows to obtain the Hubble rate:

$$H^2(z) = H_0^2 \left[\Omega_B(1+z)^3 + (1 - \Omega_B) \left(1 + \frac{1}{\gamma}\right)^{-1} b(z)(1+z)^3 + \frac{1}{\gamma}(1 - \Omega_B) \left(1 + \frac{1}{\gamma}\right)^{-1} b(z)^{1-2\alpha_2}(1+z)^{6\alpha_2+3} \right]$$

➤ We will use the matter abundance at high redshift

$$\Omega_M^{\text{eff}} = \Omega_B + (1 - \Omega_B) \left(1 + \frac{1}{\gamma}\right)^{-1} b_{\text{early}}$$

Analysis

- We perform an observational analysis of the model

Analysis

- We perform an **observational analysis** of the model



SN Ia

Analysis

- We perform an **observational analysis** of the model



SNIa

- Union2 database
- Distance moduli

Analysis

- We perform an observational analysis of the model



SNIa

- Union2 database
- Distance moduli

$$\mu(z) = 5 \log_{10}(H_0 d_L(z)) + M$$

$$\chi_{\text{SNIa}}^2 = \sum_i \frac{(\mu^{\text{obs}}(z_i) - \mu^{\text{th}}(z_i))^2}{E_i^2}$$

Analysis

- We perform an observational analysis of the model

SN Ia

CMB

- Union2 database
- Distance moduli

$$\mu(z) = 5 \log_{10}(H_0 d_L(z)) + M$$

$$\chi_{\text{SN Ia}}^2 = \sum_i \frac{(\mu^{\text{obs}}(z_i) - \mu^{\text{th}}(z_i))^2}{E_i^2}$$

Analysis

- We perform an observational analysis of the model

SN Ia

- Union2 database
- Distance moduli

CMB

- Distance priors

$$\mu(z) = 5 \log_{10}(H_0 d_L(z)) + M$$

$$\chi_{\text{SN Ia}}^2 = \sum_i \frac{(\mu^{\text{obs}}(z_i) - \mu^{\text{th}}(z_i))^2}{E_i^2}$$

Analysis

- We perform an observational analysis of the model

SN Ia

- Union2 database
- Distance moduli

$$\mu(z) = 5 \log_{10}(H_0 d_L(z)) + M$$

$$\chi_{\text{SN Ia}}^2 = \sum_i \frac{(\mu^{\text{obs}}(z_i) - \mu^{\text{th}}(z_i))^2}{E_i^2}$$

CMB

- Distance priors

$$R = \sqrt{\Omega_M^{\text{eff}} H_0^2 (1 + z_*)} d_A(z_*)$$

$$\ell_a = \pi (1 + z_*) \frac{d_A(z_*)}{r_s}$$

$$\chi_{\text{CMB}}^2 = \Delta^T \cdot \mathbf{Cov}^{-1} \cdot \Delta$$

Analysis

- We perform an **observational analysis** of the model

SN Ia

- Union2 database
- Distance moduli

$$\mu(z) = 5 \log_{10}(H_0 d_L(z)) + M$$
$$\chi_{\text{SN Ia}}^2 = \sum_i \frac{(\mu^{\text{obs}}(z_i) - \mu^{\text{th}}(z_i))^2}{E_i^2}$$

CMB

- Distance priors

$$R = \sqrt{\Omega_M^{\text{eff}} H_0^2} (1 + z_*) d_A(z_*)$$
$$\ell_a = \pi (1 + z_*) \frac{d_A(z_*)}{r_s}$$
$$\chi_{\text{CMB}}^2 = \Delta^T \cdot \mathbf{Cov}^{-1} \cdot \Delta$$



We fixed $\omega_B = \Omega_b h^2$ to the value obtained from the abundance of light elements and marginalized H_0 .

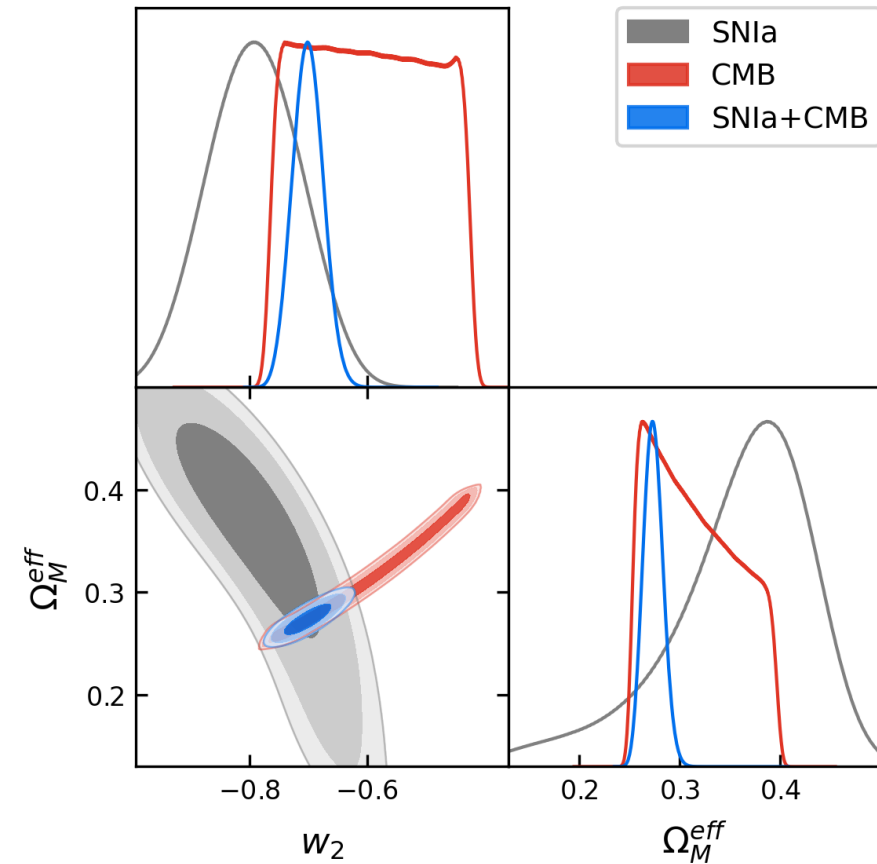
$$\Omega_B h^2 = 0.02240 \pm 0.00069$$

Analysis

- Results

Analysis

- Results

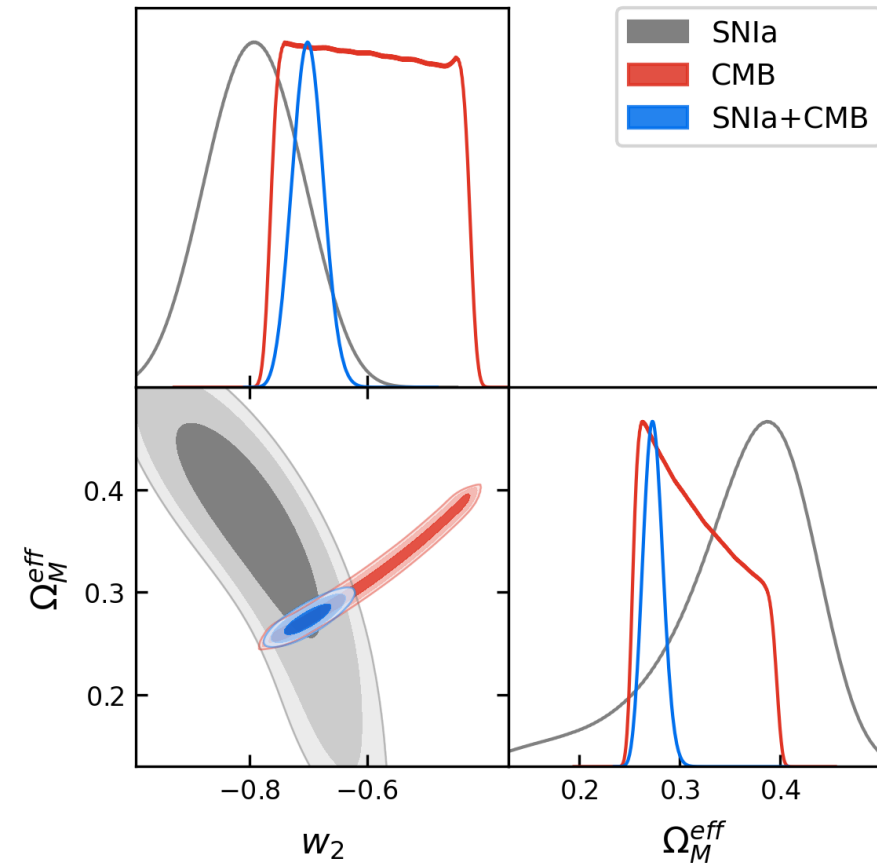


Analysis

- Results

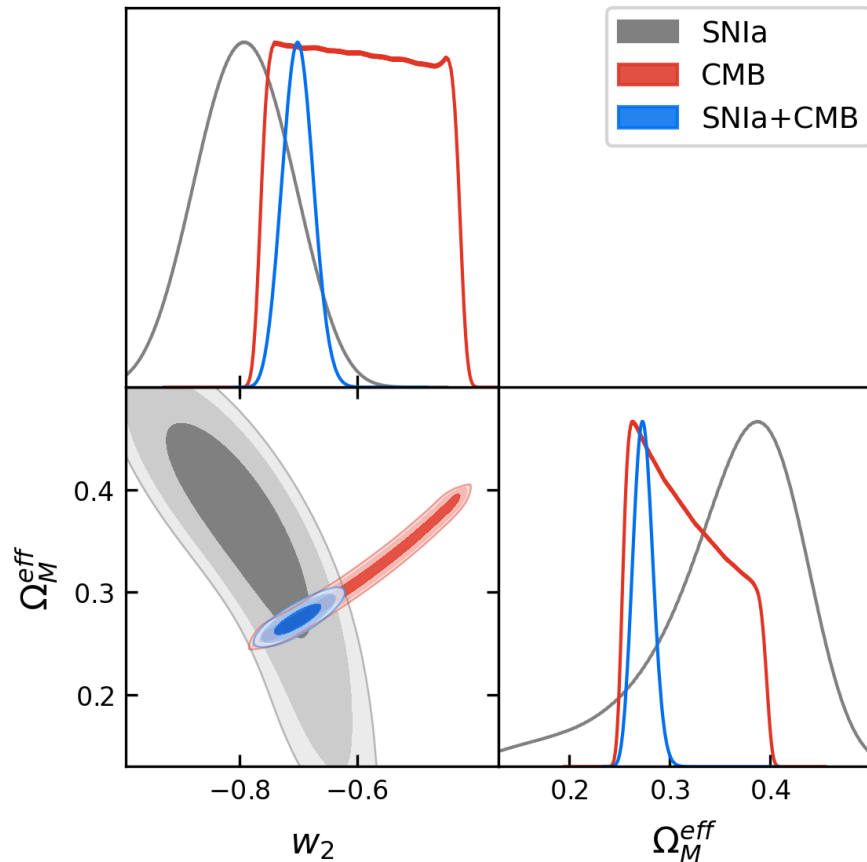
$$\Delta\text{DIC} = \text{DIC}_{\Lambda\text{CDM}} - \text{DIC}_{\text{TDiff}} = 0.76$$

$$\text{DIC} = 2\bar{\chi}^2(\mathbf{x}) - \chi^2(\bar{\mathbf{x}})$$



Analysis

- Results



$$\Delta\text{DIC} = \text{DIC}_{\Lambda\text{CDM}} - \text{DIC}_{\text{TDiff}} = 0.76$$

$$\text{DIC} = 2\bar{\chi}^2(\mathbf{x}) - \chi^2(\bar{\mathbf{x}})$$

SNIa	Best fit	χ^2_{\min}
TDiff	$w_2 = -0.813^{+0.102}_{-0.060}, \Omega_M^{\text{eff}} = 0.387^{+0.056}_{-0.078}$	542.16
wCDM	$w = -1.142^{+0.145}_{-0.184}, \Omega_M = 0.346^{+0.082}_{-0.083}$	542.64

CMB	Best fit	χ^2_{\min}
TDiff	$w_2 = -0.722^{+0.191}_{-0.058}, \Omega_M^{\text{eff}} = 0.263^{+0.077}_{-0.016}$	-
wCDM	$w = -1.278^{+0.378}_{-0.042}, \Omega_M = 0.236^{+0.097}_{-0.002}$	-

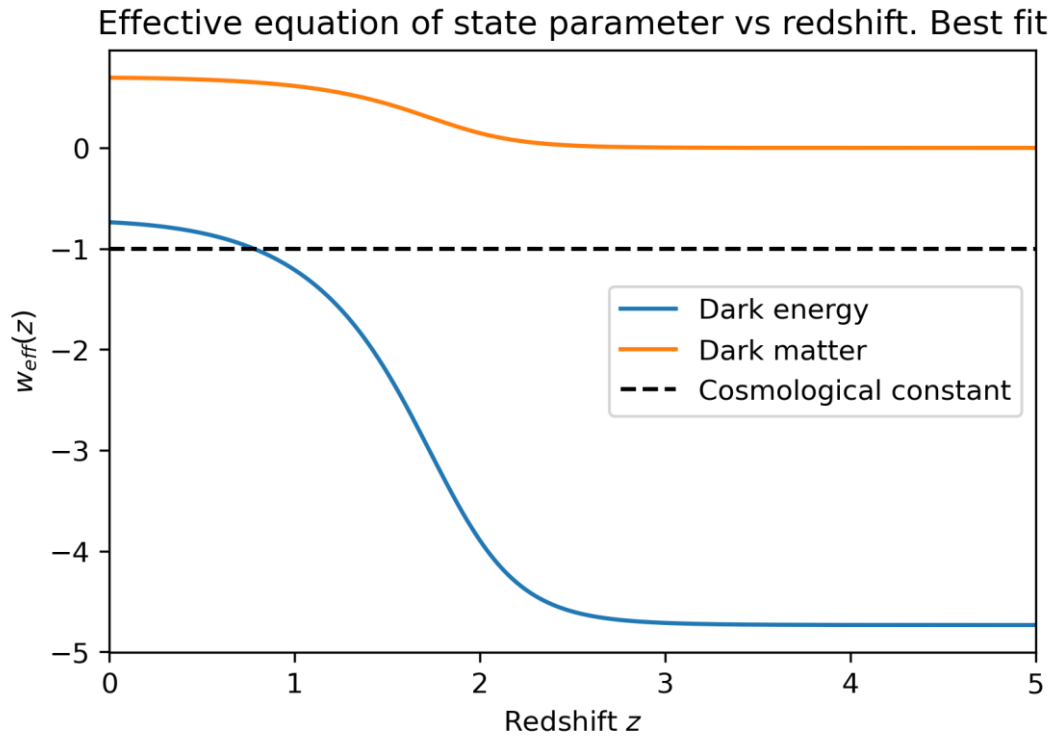
SNIa+CMB	Best fit	χ^2_{\min}
TDiff	$w_2 = -0.703^{+0.026}_{-0.026}, \Omega_M^{\text{eff}} = 0.273^{+0.010}_{-0.010}$	557.97
wCDM	$w = -1.092^{+0.034}_{-0.034}, \Omega_M = 0.292^{+0.010}_{-0.010}$	556.63

Analysis

- Equation of state evolution for the best fit

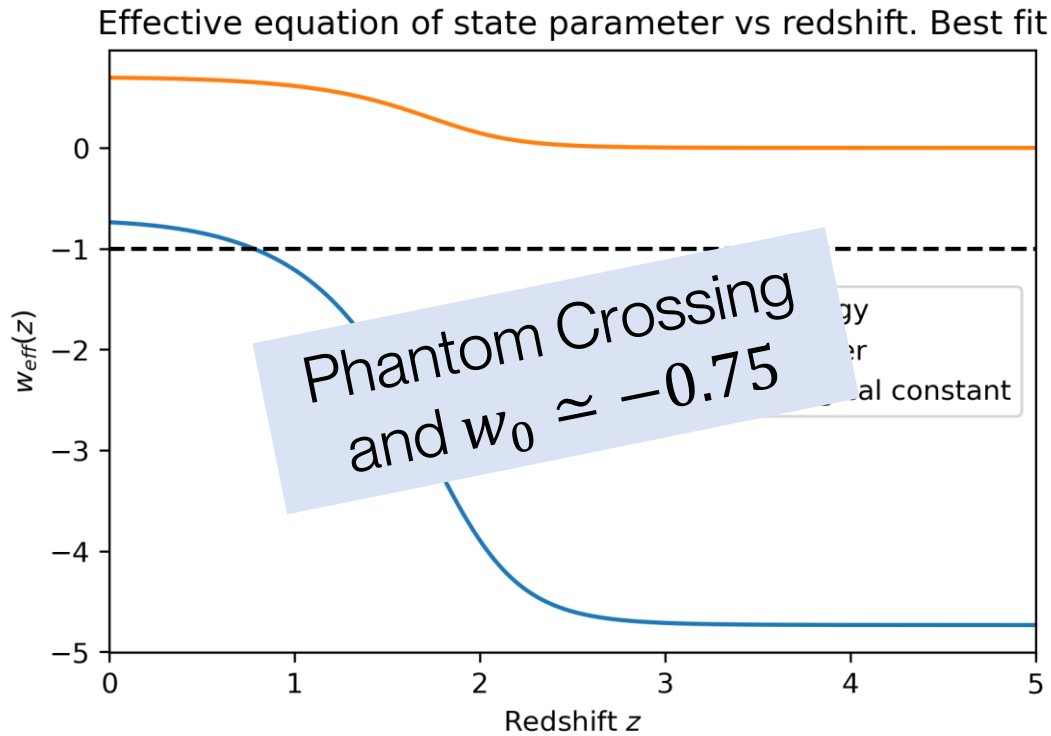
Analysis

- Equation of state evolution for the best fit



Analysis

- Equation of state evolution for the best fit





Conclusions

Conclusions

Refs.

[2307.14861]
[2402.17422]
[2409.11991]

- We have explored the consequences of breaking the symmetry from Diff \rightarrow TDiff in the dark sector.

Conclusions

Refs.

[2307.14861]
[2402.17422]
[2409.11991]

- We have explored the consequences of breaking the symmetry from Diff \rightarrow TDiff in the dark sector.

scalar field(s) with arbitrary couplings $f(g)$.

Conclusions

Refs.

[2307.14861]
[2402.17422]
[2409.11991]

- We have explored the consequences of breaking the symmetry from Diff \rightarrow TDiff in the dark sector.

scalar field(s) with arbitrary couplings $f(g)$.

- In general, there arises a constraint on the metric in the form of an equation for the determinant g .

Conclusions

Refs.

[2307.14861]
[2402.17422]
[2409.11991]

- We have explored the consequences of breaking the symmetry from Diff \rightarrow TDiff in the dark sector.

scalar field(s) with arbitrary couplings $f(g)$.

- In general, there arises a constraint on the metric in the form of an equation for the determinant g .

due to the total EMT conservation

Conclusions

Refs.

[2307.14861]
[2402.17422]
[2409.11991]

- We have explored the consequences of breaking the symmetry from Diff \rightarrow TDiff in the dark sector.

scalar field(s) with arbitrary couplings $f(g)$.

- In general, there arises a constraint on the metric in the form of an equation for the determinant g .

due to the total EMT conservation \longrightarrow effective interaction

Conclusions

Refs.

[2307.14861]
[2402.17422]
[2409.11991]

- We have explored the consequences of breaking the symmetry from Diff \rightarrow TDiff in the dark sector.

scalar field(s) with arbitrary couplings $f(g)$.

- In general, there arises a constraint on the metric in the form of an equation for the determinant g .

due to the total EMT conservation \longrightarrow effective interaction

- Possibility to describe an interacting dark sector with dynamical dark energy without relying on a specific choice for the potential.

Conclusions

Refs.

[2307.14861]
[2402.17422]
[2409.11991]

- We have explored the consequences of breaking the symmetry from Diff \rightarrow TDiff in the dark sector.

scalar field(s) with arbitrary couplings $f(g)$.

- In general, there arises a constraint on the metric in the form of an equation for the determinant g .

due to the total EMT conservation \longrightarrow effective interaction

- Possibility to describe an interacting dark sector with dynamical dark energy without relying on a specific choice for the potential.

Phantom crossing \longrightarrow H_0 tension?