



JYVÄSKYLÄN YLIOPISTO
UNIVERSITY OF JYVÄSKYLÄ



Centre of Excellence
in Quark Matter

Diffractive Structure Functions with JIMWLK

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2nd European School on the Physics of the EIC and Related Topics

June 28, 2025 @ Benicàssim



European Research Council

Established by the European Commission

Introduction



Maaninka

Population (2014-11-30)^[2]

- **Total** 3,747
- **Density** 6.5/km²

Jyväskylä

- **Urban** 117,974
- **Urban density** 1,188.7/km²

Benicàssim

Population (2018)^[1]

- **Total** 18,055
- **Density** 500/km²

3,144.44 km

Measure distance
Click on the map to add to your path
Total distance: 3,144.44 km (1,953.86 mi)

- Masters thesis 12/2024: Pure glue lattice QCD



Kvarkittoman QCD:n faasitransitio hilalla

by **Runko**, **Pyry**

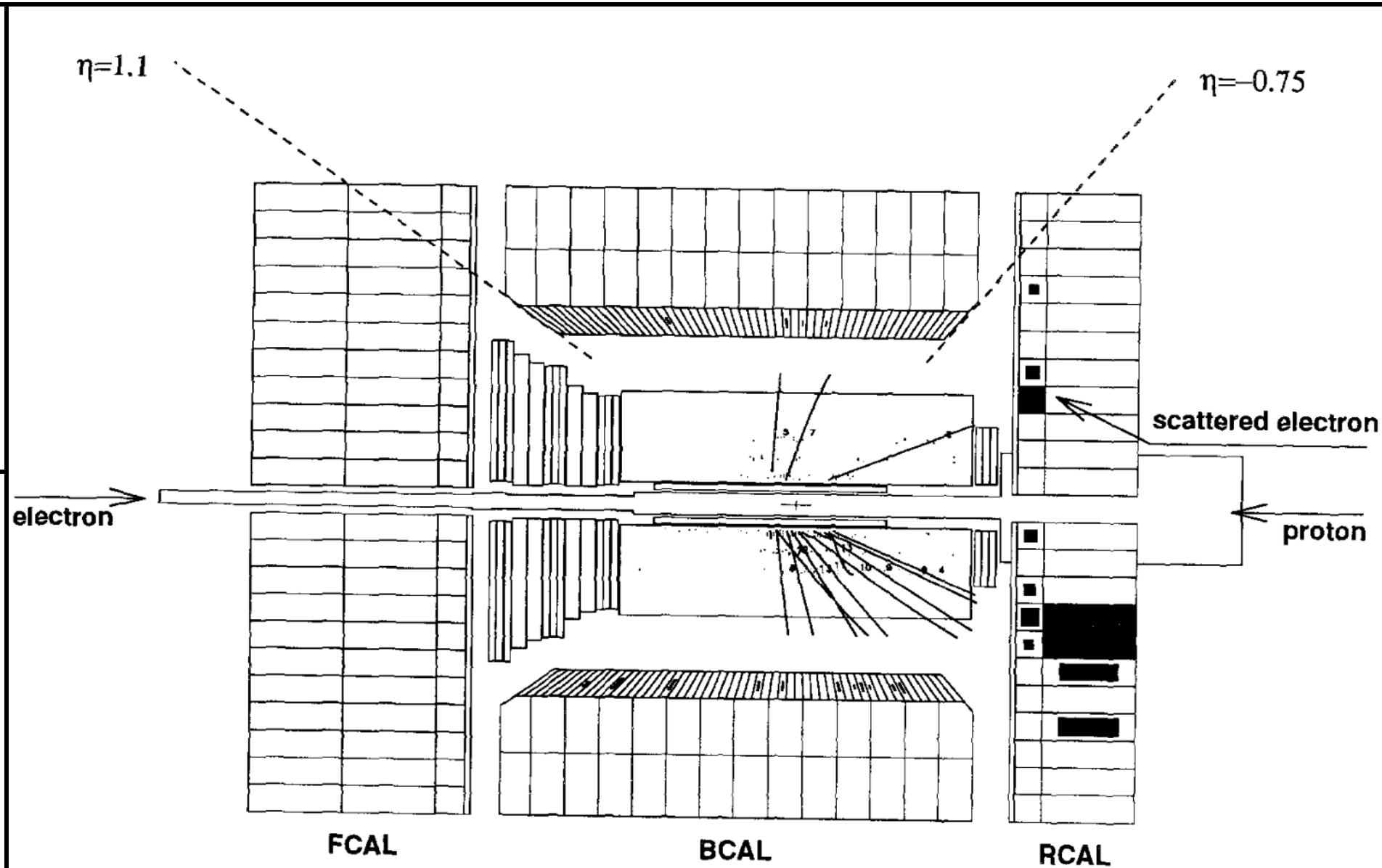
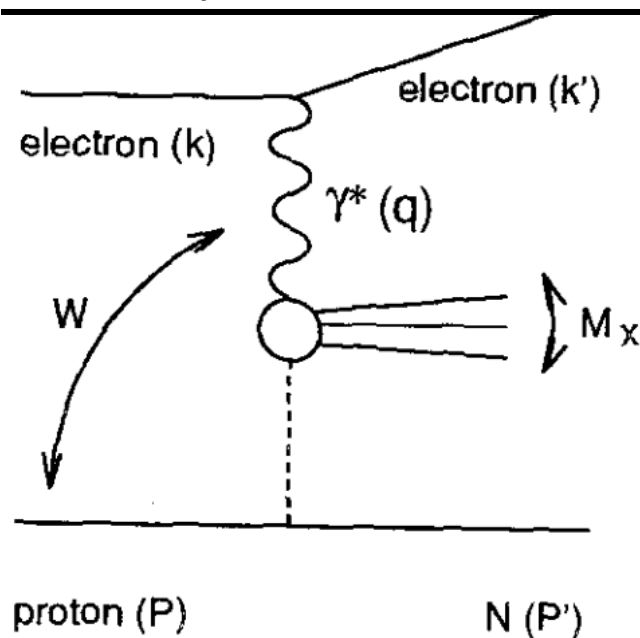
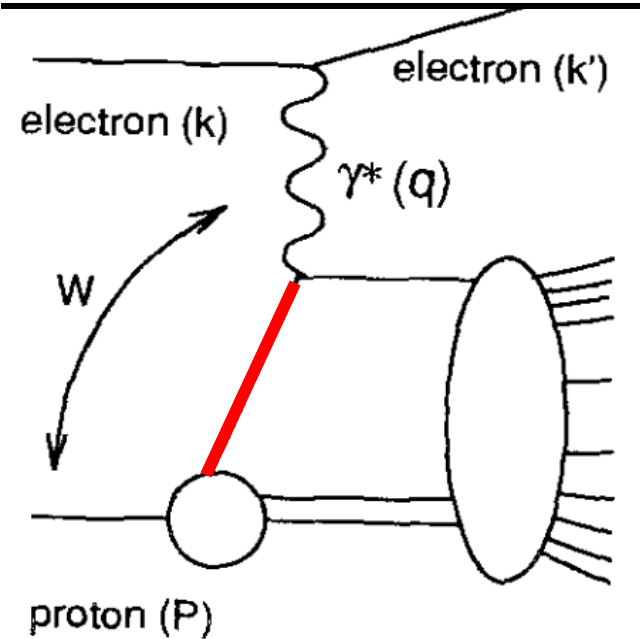
Published 2024

Theses **Master thesis**

- Summer student at Cern in JYU ALICE group in 2023,
Flow analysis



Diffraction Structure Functions with JIMWLK



Diffractive DIS

- Characteristic "Rapidity gap" in the detector between outgoing target and remnants of the virtual photon (X)
- Rapidity gap implies colorless interaction via the pomeron ($\mathbb{P}$)
- $\sim 15\%$ of DIS events were diffractive at HERA ('92 - '07)
- Note: in proton rest frame the virtual photon has energy $\sim \text{TeV}$!

**Hadron-Elektron-
Ringanlage**

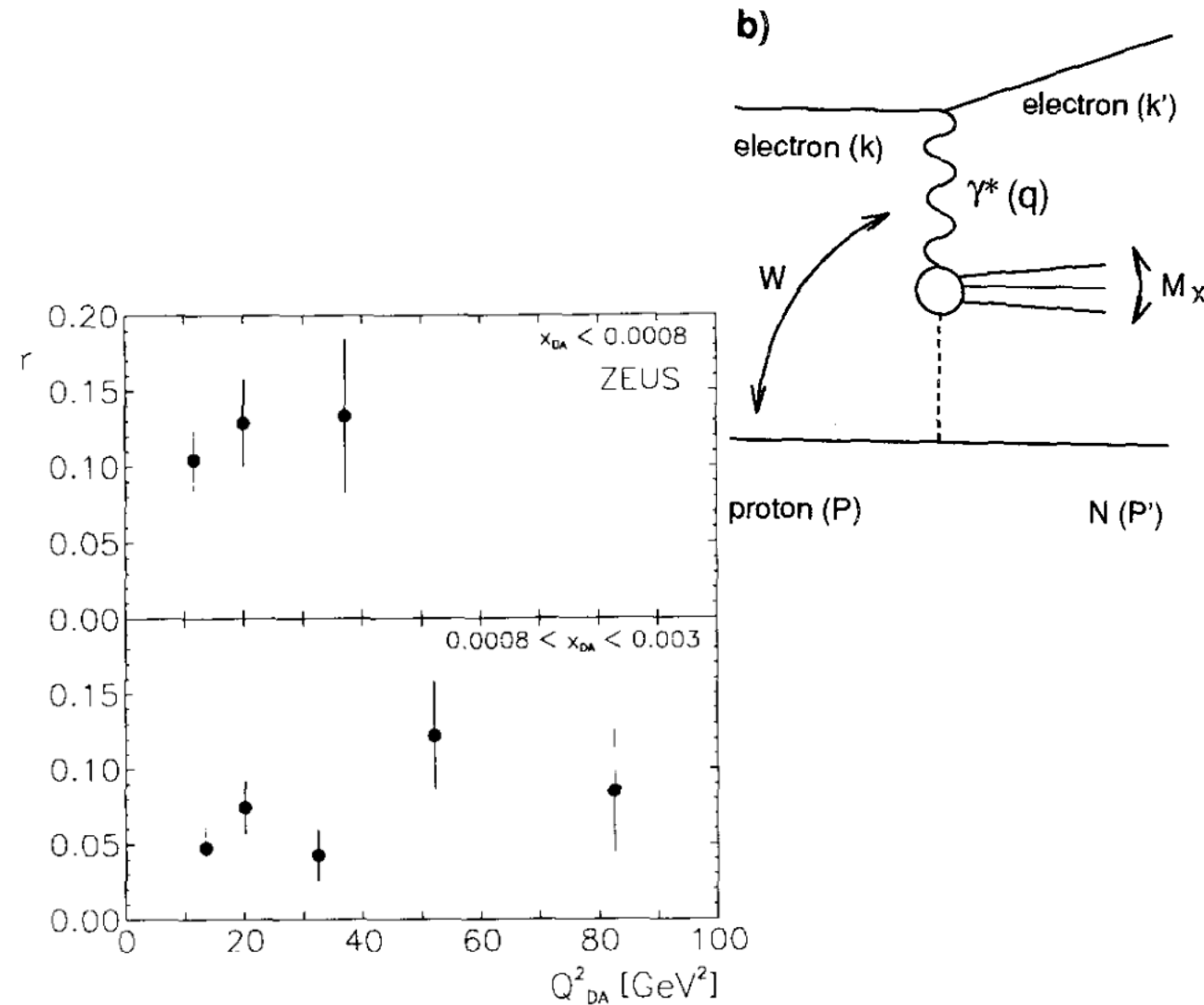
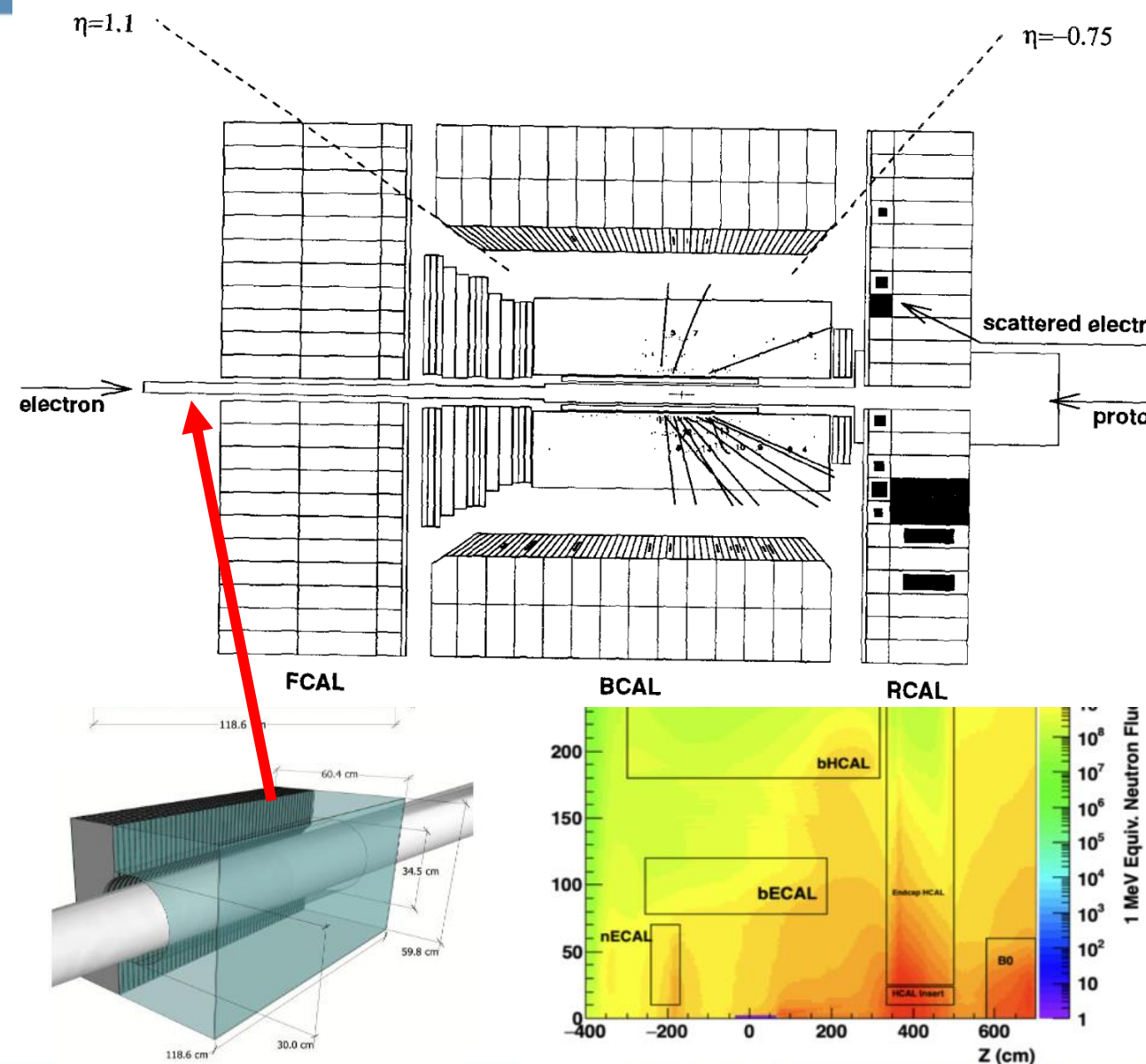


Fig. 4. Fraction r of events with a large rapidity gap, $\eta_{\max} < 1.5$, as a function of Q^2_{DA} for two ranges of x_{DA} . No acceptance corrections have been applied.

- Insert to extend central detector acceptance as much as possible in the forward direction
 - Crucial for containment of hadronic final state for inclusive kinematic reconstruction and rapidity gap diffractive measurements
 - $3.2 < \eta < 4$
 - Needs to accommodate complex beam pipe geometry
- Small hexagonal scintillator tiles for better spatial uniformity & increased light yield
- High radiation load
 - SiPMs will be serviceable
 - $\sim 10^{12}$ neq/cm²/year



Diffraction Structure Functions with JIMWLK

<https://arxiv.org/abs/1806.06783> Mäntysaari, Schenke (2018)

Structure functions

$$\frac{d^4\sigma_{ep}^{diff}}{d\beta dQ^2 dx_{IP} dt} = \frac{2\pi\alpha^2}{\beta Q^4} (1 + (1 - y)^2) F_2^{D(4)}(\beta, Q^2, x_{IP}, t)$$

- Most precise HERA data* is given as a reduced cross section (2012)
- Idea: Use JIMWLK to compute the energy ($x_{\mathbb{P}}$) dependence for diffractive structure functions (new!) for the proton and heavy nuclei $\rightarrow$ saturation
- Goals: Compare with HERA & predict for the EIC, constrain parameters with additional cross sections

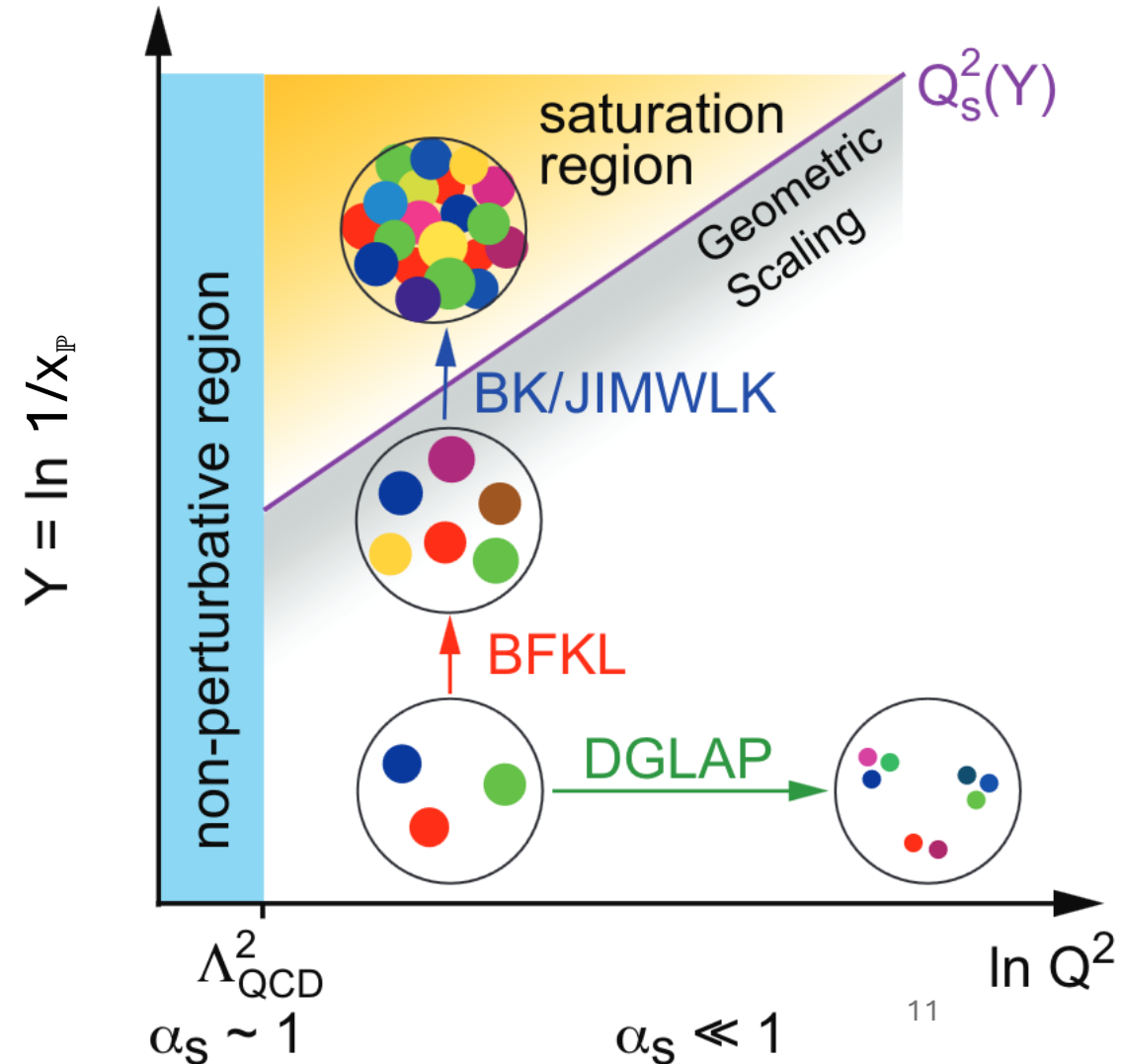
* <https://arxiv.org/abs/1207.4864> H1, ZEUS (2012)

Diffraction Structure Functions with **JIMWLK**

JIMWLK equation (Jalilian-Marian, Iancu, McLerran, Weigert, Leonidov, Kovner)

- A renormalization group equation at low x (generalises BK)
 - Color Glass Condensate: quarks are static sources of dynamic color fields
 - Classical color fields and scattering off them is described using Wilson lines $V \sim P \exp(-iA)$, **stochastic variables**
- JIMWLK: evolve the **probability distributions** of low x Wilson lines to even lower x
- Langevin formulation of LO JIMWLK:

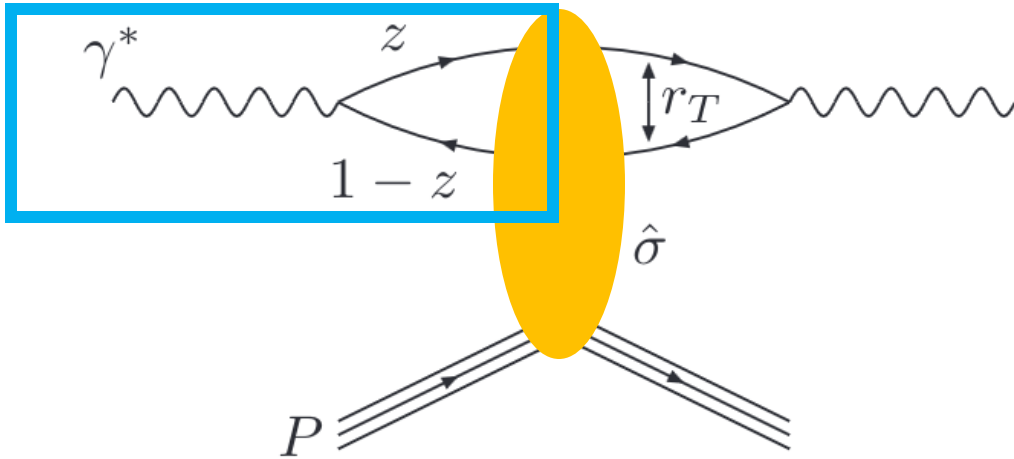
$$\frac{d}{dy} V_{\mathbf{x}} = V_{\mathbf{x}}(it^a) \left[\int d^2\mathbf{z} \varepsilon_{\mathbf{x},\mathbf{z}}^{ab,i} \xi_{\mathbf{z}}(y)_i^b + \sigma_{\mathbf{x}}^a \right]$$



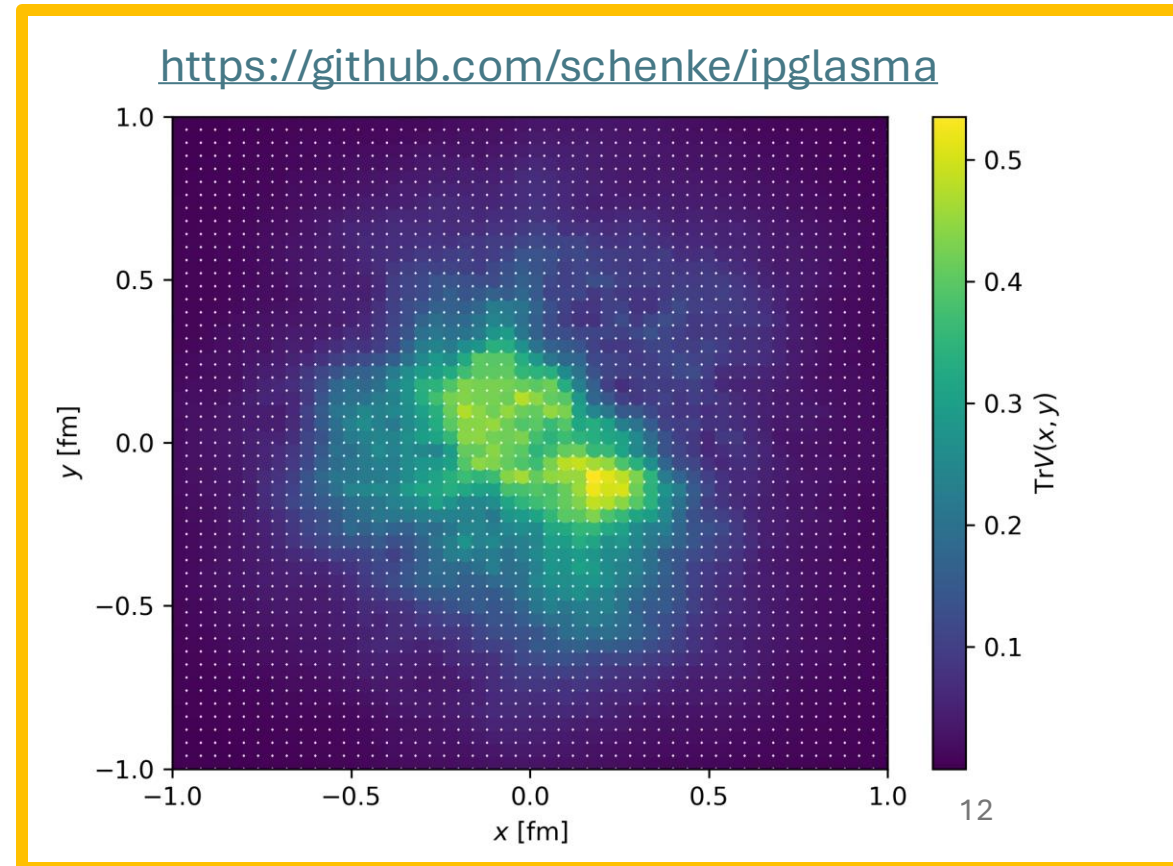
Longitudinal diffractive cross section

$$\frac{d\sigma_{L,q\bar{q}}^D}{dM_X^2 d|t|} \propto \left| \left\langle \int d^2x_0 d^2x_1 \psi_{\gamma_L^* \rightarrow q\bar{q}}(\mathbf{r}_T, z) \left(1 - \frac{1}{N_c} \text{Tr} \left[V(\mathbf{x}_0) V^\dagger(\mathbf{x}_1) \right] \right) \times \text{phase} \right\rangle \right|^2$$

- Coherent cross section = proton stays intact
- Factorizes at high energy, schematically above
- Average over Wilson line configurations



Dipole picture @ LO: photon to quark-antiquark pair



What's new?

$$\frac{d\sigma_{L,q\bar{q}}^D}{dM_X^2 d|t|} \propto \left| \left\langle \int d^2x_0 d^2x_1 \psi_{\gamma_L^* \rightarrow q\bar{q}}(\mathbf{r}_T, z) \left(1 - \frac{1}{N_c} \text{Tr} \left[V(\mathbf{x}_0) V^\dagger(\mathbf{x}_1) \right] \right) \times \text{phase} \right\rangle \right|^2$$

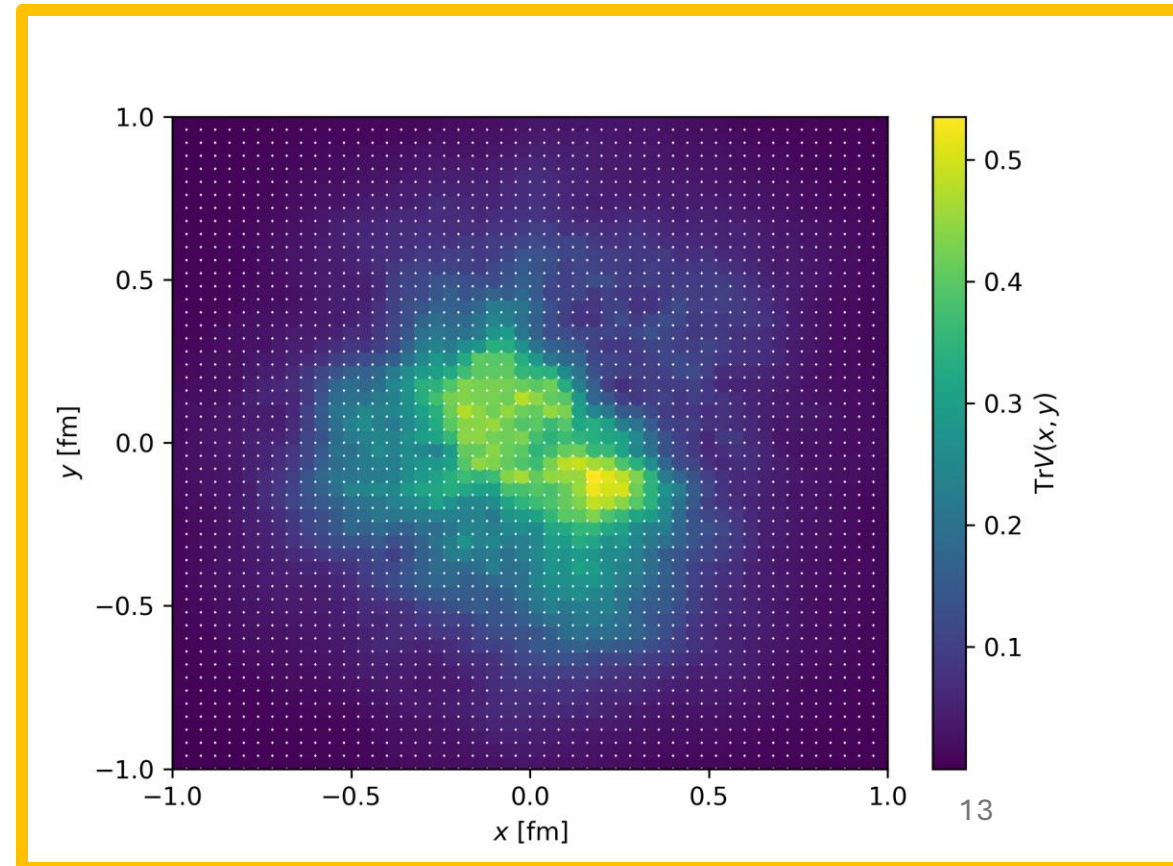
- Energy dependence is solved using JIMWLK *
 - Wilson lines already used to describe other HERA data ^
 - Evolution of the proton geometry is naturally accounted for

* <https://github.com/hejajama/jimwlk>

Schenke, Mäntysaari

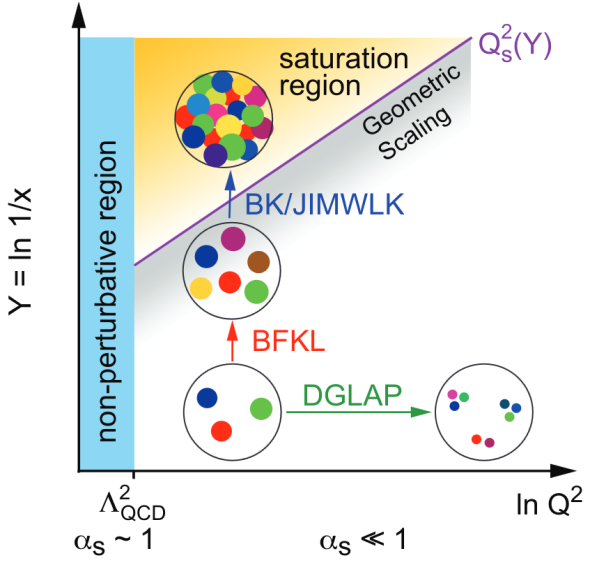
^ <https://arxiv.org/abs/1806.06783>

Mäntysaari, Schenke



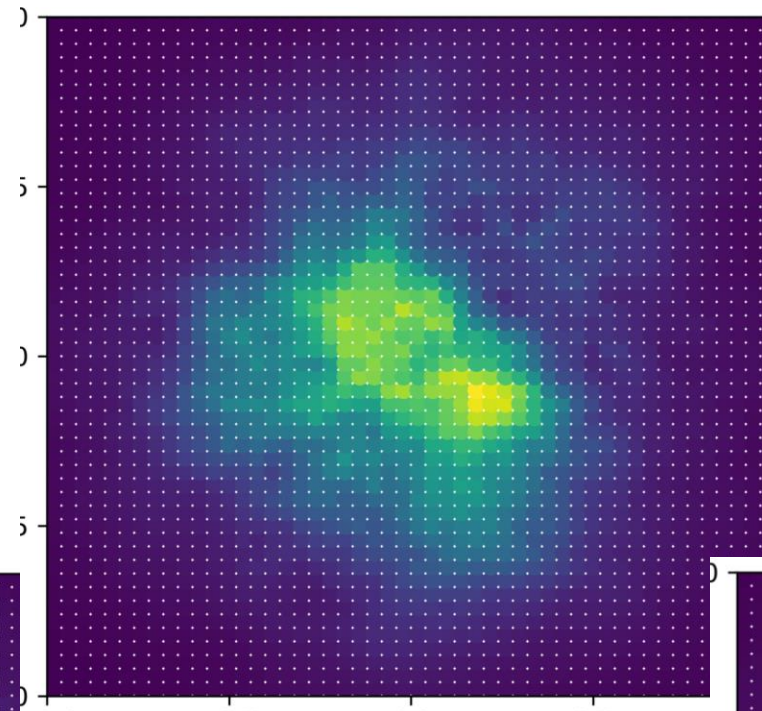
Plans for the future

- Investigate the Gaussian approximation for JIMWLK with TMDs
 - Often used to reduce the computational load, but validity has been checked only for (two) specific coordinate configurations [3]
- Work towards a Langevin formulation of NLO JIMWLK

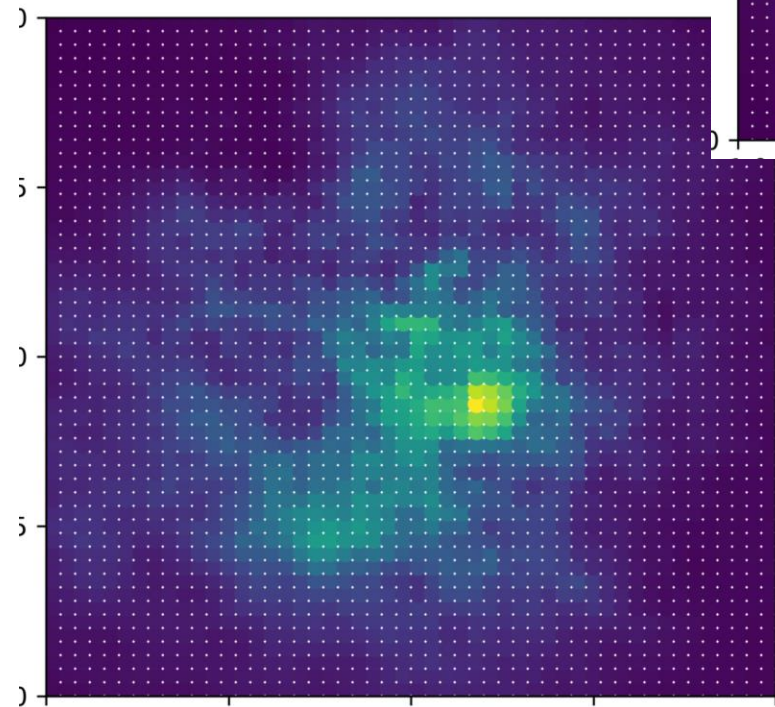


$Y = 0$

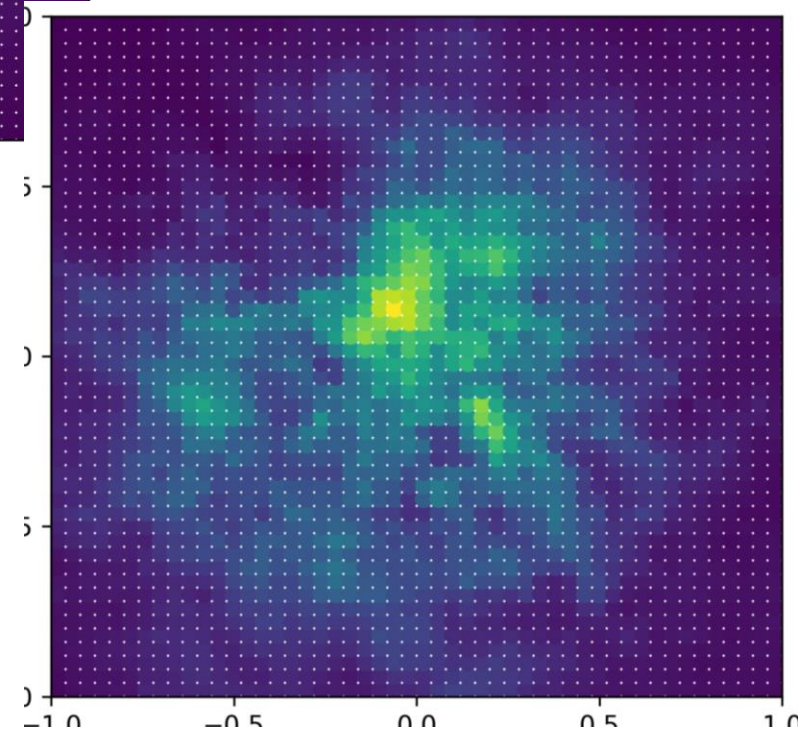
$X_0 = 0.01$



$Y = 4$



$Y = 8$



$$\frac{d}{dy} V_{\mathbf{x}} = V_{\mathbf{x}}(it^a) \left[\int d^2 \mathbf{z} \varepsilon_{\mathbf{x}, \mathbf{z}}^{ab,i} \xi_{\mathbf{z}}(y)_i^b + \sigma_{\mathbf{x}}^a \right]$$

<https://arxiv.org/abs/1806.06783>

Mäntysaari, Schenke (2018)

The random noise is Gaussian and local in coordinates, color, and rapidity with expectation value zero and

$$\langle \xi_{\mathbf{x},i}^a(y) \xi_{\mathbf{y},j}^b(y') \rangle = \delta^{ab} \delta^{ij} \delta_{\mathbf{xy}}^{(2)} \delta(y - y'). \quad (4)$$

The free parameters to be determined by the HERA data are the following:

- $g^4 \mu^2$: controls the saturation scale of the proton at $x = x_0$
- B_p : controls the proton size at the initial condition
- m : models confinement effects in the JIMWLK evolution limiting the growth of the proton at long distances
- constant α_s or Λ_{QCD} which sets the scale for the running of α_s

The deterministic drift term is

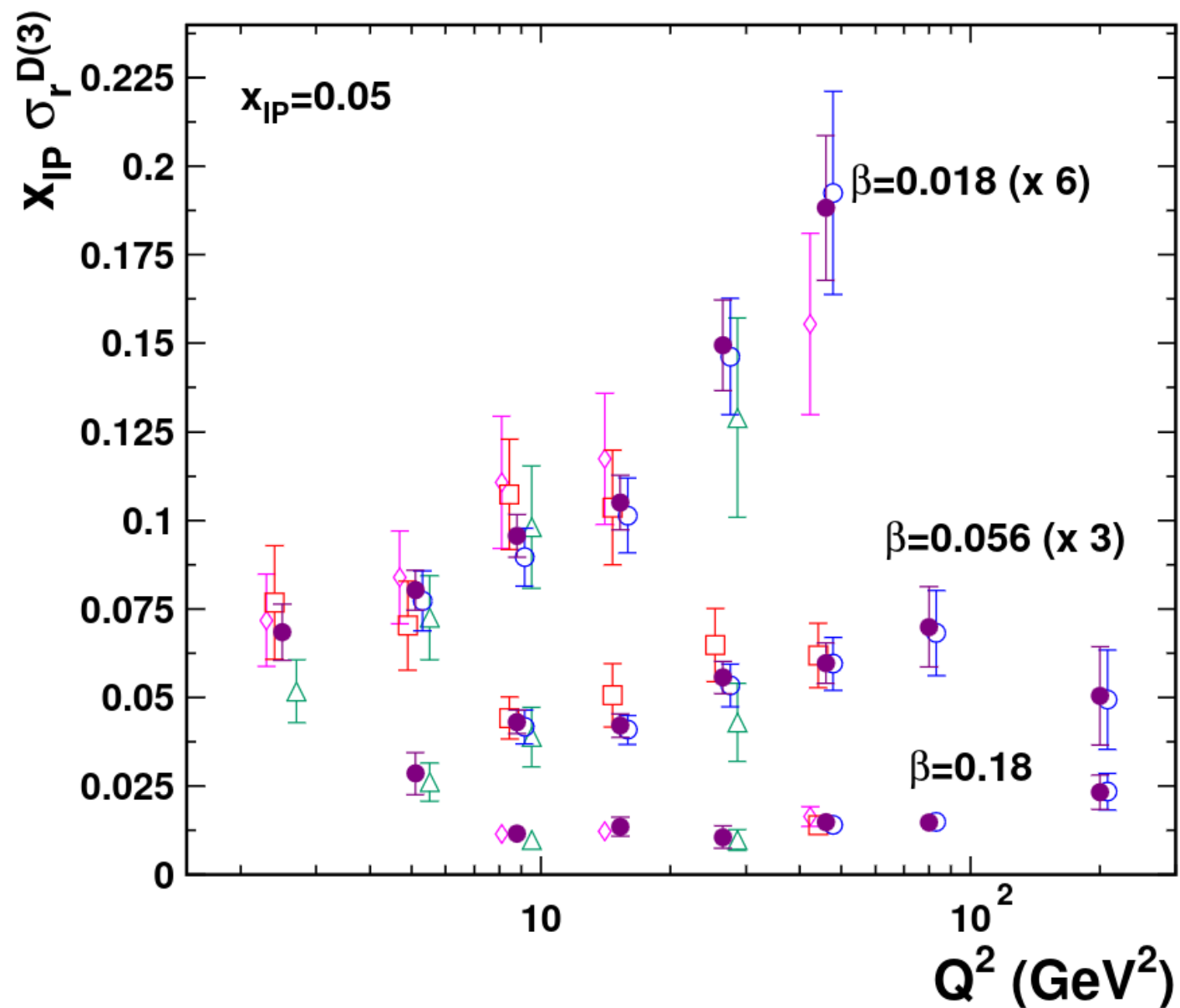
$$\sigma_{\mathbf{x}}^a = -i \frac{\alpha_s}{2\pi^2} \int d^2 \mathbf{z} S_{\mathbf{x}-\mathbf{z}} \text{Tr} [T^a U_{\mathbf{x}}^\dagger U_{\mathbf{z}}], \quad (3)$$

with $S_{\mathbf{x}} = 1/\mathbf{x}^2$ and T^a the generators of the adjoint representation. $U_{\mathbf{x}}$ are Wilson lines in the adjoint representation.

$$K_{\mathbf{x}}^i \rightarrow m|\mathbf{x}| K_1(m|\mathbf{x}|) \frac{x^i}{\mathbf{x}^2}.$$

H1 and ZEUS

○ H1 FPS HERA II △ H1 FPS HERA I ● HERA
□ ZEUS LPS 2 ◇ ZEUS LPS 1 $0.09 < |t| < 0.55 \text{ GeV}^2$



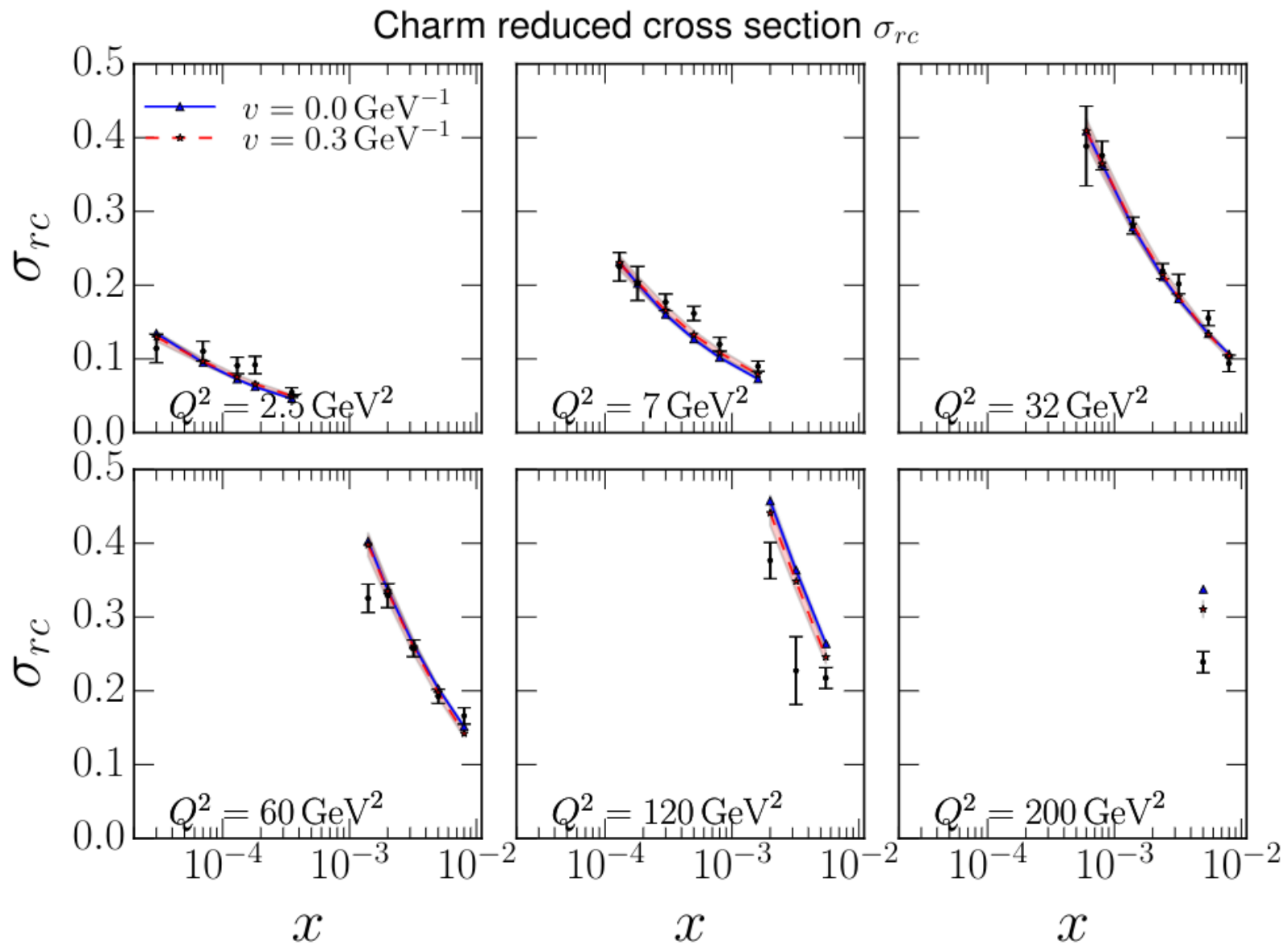
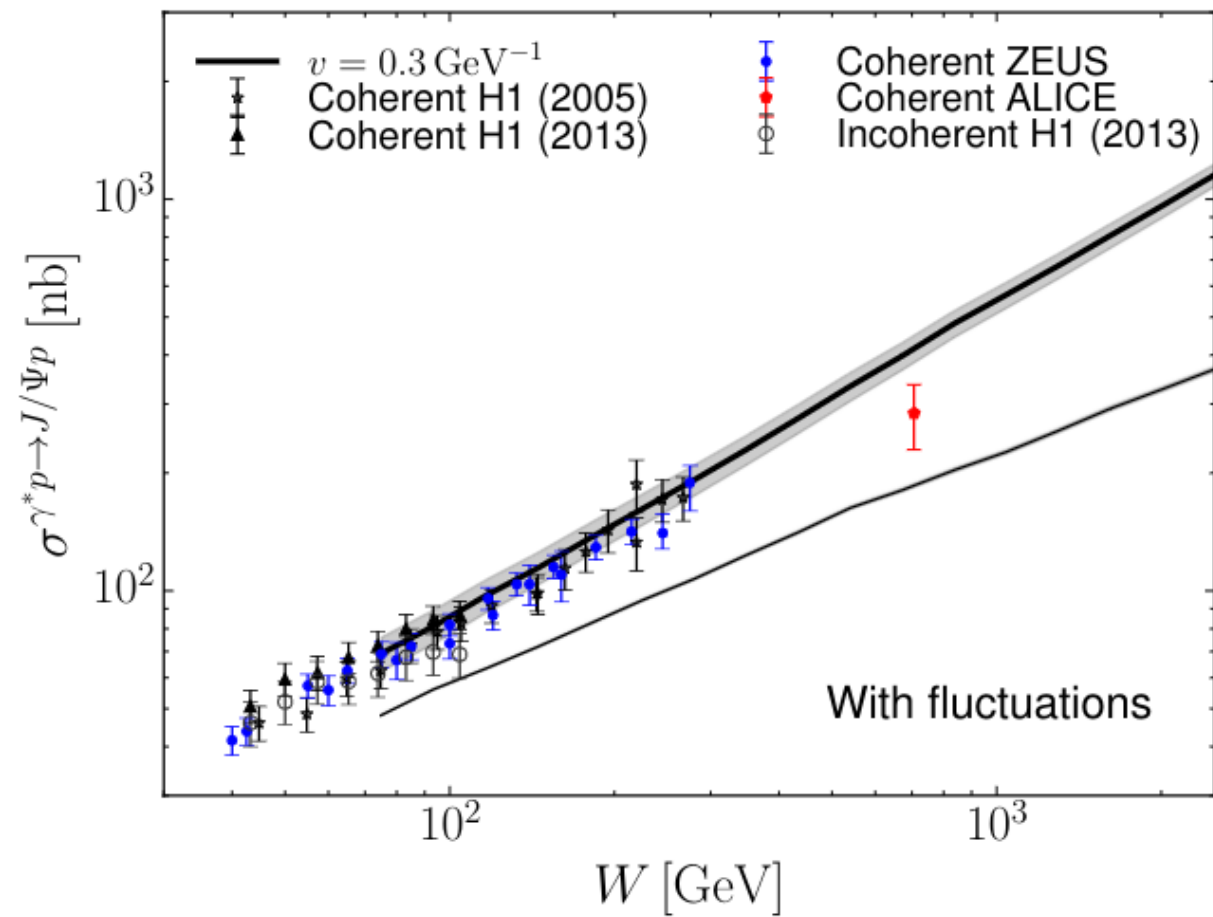
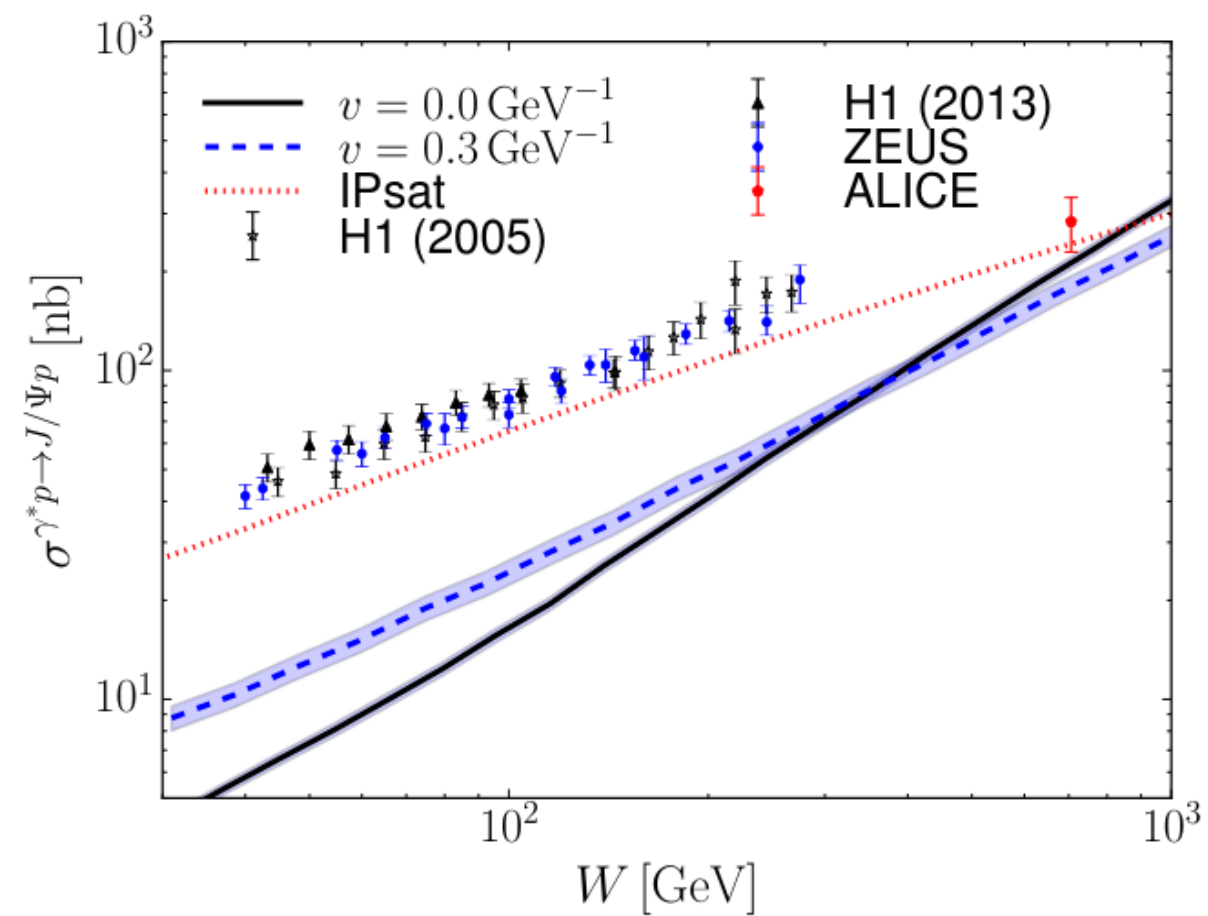


FIG. 2: Description of the charm reduced cross section using the unmodified MV model ($v = 0, \chi^2/N = 4.3$) and with an ultraviolet damping in the initial condition ($v = 0.3 \text{ GeV}^{-1}, \chi^2/N = 2.7$).

<https://arxiv.org/abs/1806.06783>

Mäntysaari, Schenke



<https://arxiv.org/abs/1806.06783>

Mäntysaari, Schenke

$$\begin{aligned}
\frac{d\sigma_{L,q\bar{q}}^D}{dM_X^2 d|t|} &= \frac{N_c}{4\pi} \frac{1}{4(2\pi)^3} \sum_{z_i} dz \sum_{\theta_{l\Delta,i}} d\theta_{l\Delta} \left| \frac{1}{N_V} \sum_V^{N_V} \sum_{\substack{m,n, \\ v,w \in V}}^N a^4 \tilde{\psi}_L^{\gamma^* \rightarrow u\bar{u}} \right. \\
&\times \left[S_{m,n,v,w} - 1 \right] e^{-i(|\mathbf{l}|\hat{\mathbf{i}} \cdot (\mathbf{x}_{m,n} - \mathbf{x}_{v,w}) + |\mathbf{\Delta}| \cos \theta_{l\Delta,i} \cdot (z_i \mathbf{x}_{m,n} + (1-z_i) \mathbf{x}_{v,w}))} \Bigg|^2 \Bigg|_{\substack{|\mathbf{l}| = \sqrt{z_i(1-z_i)} M_X \\ |\mathbf{\Delta}| = \sqrt{|t|}}}
\end{aligned}$$

$$\frac{d^4\sigma_{ep}^{diff}}{d\beta dQ^2 dx_{IP} dt} = \frac{2\pi\alpha^2}{\beta Q^4} (1 + (1 - y)^2) F_2^{D(4)}(\beta, Q^2, x_{IP}, t)$$

In addition to the usual DIS variables, $Q^2 = -q^2 = -(k - k')^2$, $y = (P \cdot q)/(P \cdot k)$, $x = Q^2/(2P \cdot q)$, other quantities are used to describe the kinematics of a diffractive DIS event:

$$x_{IP} = \frac{q(P - P')}{qp} \simeq \frac{Q^2 + M_X^2}{Q^2 + W^2}, \quad t = (P - P')^2, \quad \beta = \frac{Q^2}{2q \cdot (P - P')} \simeq \frac{Q^2}{Q^2 + M_X^2}; \quad (1)$$